## 6

## Bipartite entanglement

Entanglement is a fundamental concept in quantum information theory, considered by many to be a quintessential characteristic that distinguishes quantum systems from their classical counterparts. Informally speaking, a state of a collection of registers $\mathrm{X}_{1}, \ldots, \mathrm{X}_{n}$ is said to be entangled when it is not possible to specify the correlations that exist among the registers in classical terms. When it is possible to describe these correlations in classical terms, the registers are said to be in a separable state. Entanglement among two or more registers is therefore synonymous with a lack of separability.

This chapter introduces notions associated with bipartite entanglement, in which correlations between precisely two registers (or two collections of registers) are considered. Topics to be discussed include the property of separability, which is applicable not only to states but also to channels and measurements; aspects of entanglement manipulation and quantification; and a discussion of operational phenomena associated with entanglement, including teleportation, dense coding, and non-classical correlations among measurements on separated systems.

### 6.1 Separability

This section introduces the notion of separability, which is applicable to states, channels, and measurements on bipartite systems. It is possible to define a multipartite variant of this concept, but only bipartite separability is considered in this book.

### 6.1.1 Separable operators and states

The property of separability for operators acting on bipartite tensor product spaces is defined as follows.

Definition 6.1 For any choice of complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, the set $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is defined as the set containing all positive semidefinite operators $R \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ for which there exists an alphabet $\Sigma$ and two collections of positive semidefinite operators,

$$
\begin{equation*}
\left\{P_{a}: a \in \Sigma\right\} \subset \operatorname{Pos}(\mathcal{X}) \quad \text { and } \quad\left\{Q_{a}: a \in \Sigma\right\} \subset \operatorname{Pos}(\mathcal{Y}) \tag{6.1}
\end{equation*}
$$

such that

$$
\begin{equation*}
R=\sum_{a \in \Sigma} P_{a} \otimes Q_{a} \tag{6.2}
\end{equation*}
$$

Elements of the set $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ are called separable operators.
Remark It must be stressed that separability is defined with respect to a particular tensor product structure of the underlying complex Euclidean space of a given operator, as the previous definition reflects. When the term separable operator is used, one must therefore make this tensor product structure known (if it is not implicit). An operator $R \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y} \otimes \mathcal{Z})$ may, for instance, be an element of $\operatorname{Sep}(\mathcal{X}: \mathcal{Y} \otimes \mathcal{Z})$ but not $\operatorname{Sep}(\mathcal{X} \otimes \mathcal{Y}: \mathcal{Z})$.

By restricting the definition above to density operators, one obtains a definition of separable states.

Definition 6.2 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces. One defines

$$
\begin{equation*}
\operatorname{Sep} \mathrm{D}(\mathcal{X}: \mathcal{Y})=\operatorname{Sep}(\mathcal{X}: \mathcal{Y}) \cap \mathrm{D}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.3}
\end{equation*}
$$

Elements of the set $\operatorname{Sep} \mathrm{D}(\mathcal{X}: \mathcal{Y})$ are called separable states (or separable density operators).

## Convex properties of separable operators and states

The sets $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ and $\operatorname{Sep} \mathrm{D}(\mathcal{X}: \mathcal{Y})$ possess various properties relating to convexity, a few of which will now be observed.

Proposition 6.3 For every choice of complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, the set $\operatorname{SepD}(\mathcal{X}: \mathcal{Y})$ is convex, and the set $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is a convex cone.

Proof It will first be proved that $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is a convex cone. It suffices to prove that $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is closed under addition as well as multiplication by any nonnegative real number. To this end, assume that $R_{0}, R_{1} \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ are separable operators and $\lambda \geq 0$ is a nonnegative real number. One may write

$$
\begin{equation*}
R_{0}=\sum_{a \in \Sigma_{0}} P_{a} \otimes Q_{a} \quad \text { and } \quad R_{1}=\sum_{a \in \Sigma_{1}} P_{a} \otimes Q_{a} \tag{6.4}
\end{equation*}
$$

for disjoint alphabets, $\Sigma_{0}$ and $\Sigma_{1}$, and two collections of positive semidefinite operators,

$$
\begin{align*}
& \left\{P_{a}: a \in \Sigma_{0} \cup \Sigma_{1}\right\} \subset \operatorname{Pos}(\mathcal{X}) \\
& \left\{Q_{a}: a \in \Sigma_{0} \cup \Sigma_{1}\right\} \subset \operatorname{Pos}(\mathcal{Y}) \tag{6.5}
\end{align*}
$$

It holds that

$$
\begin{equation*}
R_{0}+R_{1}=\sum_{a \in \Sigma_{0} \cup \Sigma_{1}} P_{a} \otimes Q_{a} \tag{6.6}
\end{equation*}
$$

and therefore $R_{0}+R_{1} \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$. Moreover, it holds that

$$
\begin{equation*}
\lambda R_{0}=\sum_{a \in \Sigma_{0}}\left(\lambda P_{a}\right) \otimes Q_{a} \tag{6.7}
\end{equation*}
$$

As $\lambda P \in \operatorname{Pos}(\mathcal{X})$ for every positive semidefinite operator $P \in \operatorname{Pos}(\mathcal{X})$, it follows that $\lambda R_{0} \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$.

The fact that $\operatorname{Sep} \mathrm{D}(\mathcal{X}: \mathcal{Y})$ is convex follows from the fact that it is equal to the intersection of two convex sets, $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ and $\mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$.

The next proposition, when combined with the previous one, implies that $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is equal to the cone generated by $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$.

Proposition 6.4 Let $\mathcal{Z}$ be a complex Euclidean space, let $\mathcal{A} \subseteq \operatorname{Pos}(\mathcal{Z})$ be a cone, and assume that $\mathcal{B}=\mathcal{A} \cap \mathrm{D}(\mathcal{Z})$ is nonempty. It holds that

$$
\begin{equation*}
\mathcal{A}=\operatorname{cone}(\mathcal{B}) \tag{6.8}
\end{equation*}
$$

Proof Suppose first that $\rho \in \mathcal{B}$ and $\lambda \geq 0$. It follows that $\lambda \rho \in \mathcal{A}$ by virtue of the fact that $\mathcal{B} \subseteq \mathcal{A}$ and $\mathcal{A}$ is a cone, and therefore

$$
\begin{equation*}
\operatorname{cone}(\mathcal{B}) \subseteq \mathcal{A} \tag{6.9}
\end{equation*}
$$

Now suppose that $P \in \mathcal{A}$. If $P=0$, then one has that $P=\lambda \rho$ for $\lambda=0$ and $\rho \in \mathcal{B}$ being chosen arbitrarily. If $P \neq 0$, then consider the density operator $\rho=P / \operatorname{Tr}(P)$. It holds that $\rho \in \mathcal{A}$ because $1 / \operatorname{Tr}(P)>0$ and $\mathcal{A}$ is a cone, and therefore $\rho \in \mathcal{B}$. As $P=\lambda \rho$ for $\lambda=\operatorname{Tr}(P)>0$, it follows that $P \in \operatorname{cone}(\mathcal{B})$. Therefore,

$$
\begin{equation*}
\mathcal{A} \subseteq \operatorname{cone}(\mathcal{B}) \tag{6.10}
\end{equation*}
$$

which completes the proof.
Two equivalent ways of specifying separable states are provided by the next proposition, which is a straightforward consequence of the spectral theorem.

Proposition 6.5 Let $\xi \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ be a density operator, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$. The following statements are equivalent:

1. $\xi \in \operatorname{SepD}(\mathcal{X}: \mathcal{Y})$.
2. There exists an alphabet $\Sigma$, collections of states $\left\{\rho_{a}: a \in \Sigma\right\} \subseteq \mathrm{D}(\mathcal{X})$ and $\left\{\sigma_{a}: a \in \Sigma\right\} \subseteq \mathrm{D}(\mathcal{Y})$, and a probability vector $p \in \mathcal{P}(\Sigma)$, such that

$$
\begin{equation*}
\xi=\sum_{a \in \Sigma} p(a) \rho_{a} \otimes \sigma_{a} \tag{6.11}
\end{equation*}
$$

3. There exists an alphabet $\Sigma$, collections of unit vectors $\left\{x_{a}: a \in \Sigma\right\} \subset \mathcal{X}$ and $\left\{y_{a}: a \in \Sigma\right\} \subset \mathcal{Y}$, and a probability vector $p \in \mathcal{P}(\Sigma)$, such that

$$
\begin{equation*}
\xi=\sum_{a \in \Sigma} p(a) x_{a} x_{a}^{*} \otimes y_{a} y_{a}^{*} \tag{6.12}
\end{equation*}
$$

Proof The third statement trivially implies the second, and it is immediate that the second statement implies the first, as $\operatorname{SepD}(\mathcal{X}: \mathcal{Y})$ is convex and $\rho_{a} \otimes \sigma_{a} \in \operatorname{SepD}(\mathcal{X}: \mathcal{Y})$ for each $a \in \Sigma$. It remains to prove that the first statement implies the third.

Let $\xi \in \operatorname{SepD}(\mathcal{X}: \mathcal{Y})$. As $\xi \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$, one may write

$$
\begin{equation*}
\xi=\sum_{b \in \Gamma} P_{b} \otimes Q_{b} \tag{6.13}
\end{equation*}
$$

for some choice of an alphabet $\Gamma$ and collections $\left\{P_{b}: b \in \Gamma\right\} \subset \operatorname{Pos}(\mathcal{X})$ and $\left\{Q_{b}: b \in \Gamma\right\} \subset \operatorname{Pos}(\mathcal{Y})$ of positive semidefinite operators. Let $n=\operatorname{dim}(\mathcal{X})$, let $m=\operatorname{dim}(\mathcal{Y})$, and consider spectral decompositions of these operators as follows:

$$
\begin{equation*}
P_{b}=\sum_{j=1}^{n} \lambda_{j}\left(P_{b}\right) u_{b, j} u_{b, j}^{*} \quad \text { and } \quad Q_{b}=\sum_{k=1}^{m} \lambda_{k}\left(Q_{b}\right) v_{b, k} v_{b, k}^{*} \tag{6.14}
\end{equation*}
$$

for each $b \in \Gamma$. Define $\Sigma=\Gamma \times\{1, \ldots, n\} \times\{1, \ldots, m\}$, and define

$$
\begin{align*}
p((b, j, k)) & =\lambda_{j}\left(P_{b}\right) \lambda_{k}\left(Q_{b}\right), \\
x_{(b, j, k)} & =u_{b, j}  \tag{6.15}\\
y_{(b, j, k)} & =v_{b, k}
\end{align*}
$$

for every $(b, j, k) \in \Sigma$. A straightforward computation reveals that

$$
\begin{equation*}
\sum_{a \in \Sigma} p(a) x_{a} x_{a}^{*} \otimes y_{a} y_{a}^{*}=\sum_{b \in \Gamma} P_{b} \otimes Q_{b}=\xi \tag{6.16}
\end{equation*}
$$

Moreover, each value $p(a)$ is nonnegative, and because

$$
\begin{equation*}
\sum_{a \in \Sigma} p(a)=\operatorname{Tr}(\xi)=1 \tag{6.17}
\end{equation*}
$$

it follows that $p$ is a probability vector. It has therefore been proved that statement 1 implies statement 3 .

By the equivalence of the first and second statements in the previous proposition, it holds that a given separable state $\xi \in \operatorname{Sep} \mathrm{D}(\mathcal{X}: \mathcal{Y})$ represents a classical probability distribution over independent quantum states of a pair of registers $(X, Y)$; and in this sense the possible states of the registers X and Y , when considered in isolation, are classically correlated.

For a separable state $\xi \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$, the expression (6.12) is generally not unique - there may be many inequivalent ways that $\xi$ can be expressed in this form. It is important to observe that an expression of this form cannot necessarily be obtained directly from a spectral decomposition of $\xi$. Indeed, for some choices of $\xi \in \operatorname{Sep} \mathrm{D}(\mathcal{X}: \mathcal{Y})$ it may hold that every expression of $\xi$ in the form (6.12) requires that $\Sigma$ has cardinality strictly larger than $\operatorname{rank}(\xi)$. An upper bound on the size of the alphabet $\Sigma$ required for an expression of the form (6.12) to exist may, however, be obtained from Carathéodory's theorem (Theorem 1.9).

Proposition 6.6 Let $\xi \in \operatorname{SepD}(\mathcal{X}: \mathcal{Y})$ be a separable state, for $\mathcal{X}$ and $\mathcal{Y}$ being complex Euclidean spaces. There exists an alphabet $\Sigma$ such that $|\Sigma| \leq \operatorname{rank}(\xi)^{2}$, two collections of unit vectors $\left\{x_{a}: a \in \Sigma\right\} \subset \mathcal{X}$ and $\left\{y_{a}: a \in \Sigma\right\} \subset \mathcal{Y}$, and a probability vector $p \in \mathcal{P}(\Sigma)$ such that

$$
\begin{equation*}
\xi=\sum_{a \in \Sigma} p(a) x_{a} x_{a}^{*} \otimes y_{a} y_{a}^{*} \tag{6.18}
\end{equation*}
$$

Proof By Proposition 6.5 it holds that

$$
\begin{equation*}
\operatorname{SepD}(\mathcal{X}: \mathcal{Y})=\operatorname{conv}\left\{x x^{*} \otimes y y^{*}: x \in \mathcal{S}(\mathcal{X}), y \in \mathcal{S}(\mathcal{Y})\right\} \tag{6.19}
\end{equation*}
$$

from which it follows that $\xi$ is contained in the set

$$
\begin{equation*}
\operatorname{conv}\left\{x x^{*} \otimes y y^{*}: x \in \mathcal{S}(\mathcal{X}), y \in \mathcal{S}(\mathcal{Y}), \operatorname{im}\left(x x^{*} \otimes y y^{*}\right) \subseteq \operatorname{im}(\xi)\right\} \tag{6.20}
\end{equation*}
$$

Every density operator $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ satisfying $\operatorname{im}(\rho) \subseteq \operatorname{im}(\xi)$ is contained in the real affine subspace

$$
\begin{equation*}
\{H \in \operatorname{Herm}(\mathcal{X} \otimes \mathcal{Y}): \operatorname{im}(H) \subseteq \operatorname{im}(\xi), \operatorname{Tr}(H)=1\} \tag{6.21}
\end{equation*}
$$

of dimension $\operatorname{rank}(\xi)^{2}-1$, and therefore the proposition follows directly from Carathéodory's theorem.

By combining the previous proposition with Proposition 6.4, one obtains the following corollary.

Corollary 6.7 Let $R \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ be a nonzero separable operator, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$. There exists an alphabet $\Sigma$ such that $|\Sigma| \leq \operatorname{rank}(R)^{2}$, along with two collections of vectors $\left\{x_{a}: a \in \Sigma\right\} \subset \mathcal{X}$ and $\left\{y_{a}: a \in \Sigma\right\} \subset \mathcal{Y}$, such that

$$
\begin{equation*}
R=\sum_{a \in \Sigma} x_{a} x_{a}^{*} \otimes y_{a} y_{a}^{*} \tag{6.22}
\end{equation*}
$$

The last observation to be made about separable operators and states in this subsection is the following proposition, which establishes a basic topological property of the sets $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ and $\operatorname{SepD}(\mathcal{X}: \mathcal{Y})$.

Proposition 6.8 For every choice of complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, the set $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is compact and the set $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is closed.

Proof The unit spheres $\mathcal{S}(\mathcal{X})$ and $\mathcal{S}(\mathcal{Y})$ are compact, which implies that their Cartesian product $\mathcal{S}(\mathcal{X}) \times \mathcal{S}(\mathcal{Y})$ is also compact. The function

$$
\begin{equation*}
\phi: \mathcal{S}(\mathcal{X}) \times \mathcal{S}(\mathcal{Y}) \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}):(x, y) \mapsto x x^{*} \otimes y y^{*} \tag{6.23}
\end{equation*}
$$

is continuous, and therefore the set

$$
\begin{equation*}
\phi(\mathcal{S}(\mathcal{X}) \times \mathcal{S}(\mathcal{Y}))=\left\{x x^{*} \otimes y y^{*}: x \in \mathcal{S}(\mathcal{X}), y \in \mathcal{S}(\mathcal{Y})\right\} \tag{6.24}
\end{equation*}
$$

is compact. Because the convex hull of a compact set is necessarily compact, it follows that $\operatorname{Sep} \mathrm{D}(\mathcal{X}: \mathcal{Y})$ is compact.

As $\operatorname{SepD}(\mathcal{X}: \mathcal{Y})$ is compact, and does not include 0 , the cone it generates is closed, and therefore $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is closed.

## The Horodecki criterion

The next theorem provides an alternative characterization of separability, demonstrating that the property of separability for operators has a close connection with the property of positivity for maps.

Theorem 6.9 (Horodecki criterion) Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces and let $R \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ be a positive semidefinite operator. The following three statements are equivalent:

1. $R \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$.
2. For every choice of a complex Euclidean space $\mathcal{Z}$ and a positive map $\Phi \in \mathrm{T}(\mathcal{X}, \mathcal{Z})$ it holds that

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(R) \in \operatorname{Pos}(\mathcal{Z} \otimes \mathcal{Y}) \tag{6.25}
\end{equation*}
$$

3. For every positive and unital map $\Phi \in \mathrm{T}(\mathcal{X}, \mathcal{Y})$, it holds that

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(R) \in \operatorname{Pos}(\mathcal{Y} \otimes \mathcal{Y}) \tag{6.26}
\end{equation*}
$$

Proof Suppose first that $R \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$, so that

$$
\begin{equation*}
R=\sum_{a \in \Sigma} P_{a} \otimes Q_{a} \tag{6.27}
\end{equation*}
$$

for some choice of an alphabet $\Sigma$ and collections $\left\{P_{a}: a \in \Sigma\right\} \subset \operatorname{Pos}(\mathcal{X})$ and $\left\{Q_{a}: a \in \Sigma\right\} \subset \operatorname{Pos}(\mathcal{Y})$. For every complex Euclidean space $\mathcal{Z}$ and every positive map $\Phi \in \mathrm{T}(\mathcal{X}, \mathcal{Z})$ it holds that

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(R)=\sum_{a \in \Sigma} \Phi\left(P_{a}\right) \otimes Q_{a} \in \operatorname{Pos}(\mathcal{Z} \otimes \mathcal{Y}) \tag{6.28}
\end{equation*}
$$

by virtue of the fact that $\Phi\left(P_{a}\right)$ is a positive semidefinite operator for each $a \in \Sigma$. Statement 1 therefore implies statement 2 .

Statement 2 trivially implies statement 3.
Finally, the fact that statement 3 implies statement 1 will be proved in the contrapositive form. To this end, assume $R \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ is not a separable operator. As $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is a closed, convex cone within the real vector space $\operatorname{Herm}(\mathcal{X} \otimes \mathcal{Y})$, the hyperplane separation theorem (Theorem 1.11) implies that there must exist a Hermitian operator $H \in \operatorname{Herm}(\mathcal{X} \otimes \mathcal{Y})$ such that $\langle H, R\rangle<0$ and $\langle H, S\rangle \geq 0$ for every $S \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$. The operator $H$ will be used to define a positive and unital $\operatorname{map} \Phi \in \mathrm{T}(\mathcal{X}, \mathcal{Y})$ for which

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(R) \notin \operatorname{Pos}(\mathcal{Y} \otimes \mathcal{Y}) \tag{6.29}
\end{equation*}
$$

First, let $\Psi \in \mathrm{T}(\mathcal{Y}, \mathcal{X})$ be the unique map for which $J(\Psi)=H$, choose $\varepsilon>0$ to be a sufficiently small positive real number so that the inequality

$$
\begin{equation*}
\langle H, R\rangle+\varepsilon \operatorname{Tr}(R)<0 \tag{6.30}
\end{equation*}
$$

is satisfied, and define $\Xi \in \mathrm{T}(\mathcal{X}, \mathcal{Y})$ as

$$
\begin{equation*}
\Xi(X)=\Psi^{*}(X)+\varepsilon \operatorname{Tr}(X) \mathbb{1}_{\mathcal{Y}} \tag{6.31}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$. For arbitrarily chosen positive semidefinite operators $P \in \operatorname{Pos}(\mathcal{X})$ and $Q \in \operatorname{Pos}(\mathcal{Y})$, it is the case that

$$
\begin{equation*}
P \otimes \bar{Q} \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y}) \tag{6.32}
\end{equation*}
$$

and therefore

$$
\begin{equation*}
0 \leq\langle H, P \otimes \bar{Q}\rangle=\langle P \otimes \bar{Q}, J(\Psi)\rangle=\langle P, \Psi(Q)\rangle \tag{6.33}
\end{equation*}
$$

The fact that this inequality holds for every choice of $P \in \operatorname{Pos}(\mathcal{X})$ and $Q \in \operatorname{Pos}(\mathcal{Y})$ implies that $\Psi(Q) \in \operatorname{Pos}(\mathcal{X})$ for every choice of $Q \in \operatorname{Pos}(\mathcal{Y})$, and therefore $\Psi$ is a positive map. It follows from Proposition 2.18 that $\Psi^{*}$ is a positive map as well. For every nonzero positive semidefinite operator
$P \in \operatorname{Pos}(\mathcal{X})$, the operator $\Xi(P)$ is therefore equal to a positive semidefinite operator $\Psi^{*}(P)$ plus a positive multiple of the identity operator.

Now let $A=\Xi\left(\mathbb{1}_{\mathcal{X}}\right)$, which is necessarily a positive definite operator, and define $\Phi \in \mathrm{T}(\mathcal{X}, \mathcal{Y})$ as

$$
\begin{equation*}
\Phi(X)=A^{-\frac{1}{2}} \Xi(X) A^{-\frac{1}{2}} \tag{6.34}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$. It remains to verify that $\Phi$ is a positive and unital map for which (6.29) holds. The positivity of $\Phi$ follows from the fact that $\Xi$ is positive, and it holds that

$$
\begin{equation*}
\Phi\left(\mathbb{1}_{\mathcal{X}}\right)=A^{-\frac{1}{2}} \Xi\left(\mathbb{1}_{\mathcal{X}}\right) A^{-\frac{1}{2}}=A^{-\frac{1}{2}} A A^{-\frac{1}{2}}=\mathbb{1}_{\mathcal{Y}} \tag{6.35}
\end{equation*}
$$

establishing that $\Phi$ is unital. Finally, through the following computation, one may verify that the operator $\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{y})}\right)(R)$ is not positive semidefinite:

$$
\begin{align*}
&\langle\operatorname{vec}\left.(\sqrt{A}) \operatorname{vec}(\sqrt{A})^{*},\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(R)\right\rangle \\
&=\left\langle\operatorname{vec}\left(\mathbb{1}_{\mathcal{Y}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{Y}}\right)^{*},\left(\Xi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(R)\right\rangle \\
&=\left\langle J\left(\Xi^{*}\right), R\right\rangle  \tag{6.36}\\
& \quad=\left\langle J(\Psi)+\varepsilon \mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}, R\right\rangle \\
& \quad=\langle H, R\rangle+\varepsilon \operatorname{Tr}(R) \\
& \quad<0 .
\end{align*}
$$

This completes the proof.
One immediate application of Theorem 6.9 is that it provides a method for proving that certain positive semidefinite operators are not separable. The following example demonstrates this method for two families of states known as Werner states and isotropic states.

Example 6.10 Let $\Sigma$ be an alphabet, and let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces of the form $\mathcal{X}=\mathbb{C}^{\Sigma}$ and $\mathcal{Y}=\mathbb{C}^{\Sigma}$. The swap operator $W \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$ is the unique operator satisfying

$$
\begin{equation*}
W(x \otimes y)=y \otimes x \tag{6.37}
\end{equation*}
$$

for all vectors $x, y \in \mathbb{C}^{\Sigma}$. Equivalently, this operator is given by

$$
\begin{equation*}
W=\sum_{a, b \in \Sigma} E_{a, b} \otimes E_{b, a} \tag{6.38}
\end{equation*}
$$

The operator $W$ is both unitary and Hermitian, with eigenvalues 1 and -1 .

The eigenspace of $W$ corresponding to the eigenvalue 1 is spanned by the orthonormal collection

$$
\begin{equation*}
\left\{\frac{e_{a} \otimes e_{b}+e_{b} \otimes e_{a}}{\sqrt{2}}: a, b \in \Sigma, a<b\right\} \cup\left\{e_{a} \otimes e_{a}: a \in \Sigma\right\} \tag{6.39}
\end{equation*}
$$

where it has been assumed that a total ordering of the alphabet $\Sigma$ has been fixed, while the eigenspace corresponding to the eigenvalue -1 is spanned by the orthonormal collection

$$
\begin{equation*}
\left\{\frac{e_{a} \otimes e_{b}-e_{b} \otimes e_{a}}{\sqrt{2}}: a, b \in \Sigma, a<b\right\} . \tag{6.40}
\end{equation*}
$$

Let $n=|\Sigma|$, and define projection operators $\Delta_{0}, \Delta_{1}, \Pi_{0}, \Pi_{1} \in \operatorname{Proj}(\mathcal{X} \otimes \mathcal{Y})$ as follows:

$$
\begin{align*}
\Delta_{0} & =\frac{1}{n} \sum_{a, b \in \Sigma} E_{a, b} \otimes E_{a, b}, & \Pi_{0}=\frac{1}{2} \mathbb{1} \otimes \mathbb{1}+\frac{1}{2} W,  \tag{6.41}\\
\Delta_{1} & =\mathbb{1} \otimes \mathbb{1}-\Delta_{0}, & \Pi_{1}=\mathbb{1} \otimes \mathbb{1}-\Pi_{0} . \tag{6.42}
\end{align*}
$$

That these operators are indeed projection operators follows from the fact that they are Hermitian and square to themselves. Alternatively, one may observe that $\Delta_{0}=u u^{*}$ is the projection onto the one-dimensional subspace of $\mathcal{X} \otimes \mathcal{Y}$ spanned by the unit vector

$$
\begin{equation*}
u=\frac{1}{\sqrt{n}} \sum_{a \in \Sigma} e_{a} \otimes e_{a} \tag{6.43}
\end{equation*}
$$

$\Delta_{1}$ is the projection onto the orthogonal complement of this subspace, and $\Pi_{0}$ and $\Pi_{1}$ are the projection operators onto the subspaces spanned by the collections (6.39) and (6.40), respectively. (The images of $\Pi_{0}$ and $\Pi_{1}$ are also known as the symmetric and anti-symmetric subspaces of $\mathbb{C}^{\Sigma} \otimes \mathbb{C}^{\Sigma}$, and are considered in greater detail and generality in Chapter 7.) It holds that

$$
\begin{array}{ll}
\operatorname{rank}\left(\Delta_{0}\right)=1, & \operatorname{rank}\left(\Pi_{0}\right)=\binom{n+1}{2},  \tag{6.44}\\
\operatorname{rank}\left(\Delta_{1}\right)=n^{2}-1, & \operatorname{rank}\left(\Pi_{1}\right)=\binom{n}{2}
\end{array}
$$

States of the form

$$
\begin{equation*}
\lambda \Delta_{0}+(1-\lambda) \frac{\Delta_{1}}{n^{2}-1} \tag{6.45}
\end{equation*}
$$

are known as isotropic states, and states of the form

$$
\begin{equation*}
\lambda \frac{\Pi_{0}}{\binom{n+1}{2}}+(1-\lambda) \frac{\Pi_{1}}{\binom{n}{2}} \tag{6.46}
\end{equation*}
$$

are known as Werner states (for $\lambda \in[0,1]$ in both cases).
Now, let $\mathrm{T} \in \mathrm{T}(\mathcal{X})$ denote the transpose mapping, defined by the action $\mathrm{T}(X)=X^{\top}$ for all $X \in \mathrm{~L}(\mathcal{X})$. The mapping T is a positive map. Using the observation that

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{y})}\right)\left(\Delta_{0}\right)=\frac{1}{n} W \tag{6.47}
\end{equation*}
$$

which may be verified directly, as well as $\mathrm{T}\left(\mathbb{1}_{\mathcal{X}}\right)=\mathbb{1}_{\mathcal{X}}$ and $\mathrm{T}^{2}=\mathbb{1}_{\mathrm{L}(\mathcal{X})}$, the following relations may be obtained:

$$
\begin{align*}
& \left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\Delta_{0}\right)=\frac{1}{n} \Pi_{0}-\frac{1}{n} \Pi_{1}  \tag{6.48}\\
& \left(\mathrm{~T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\Delta_{1}\right)=\frac{n-1}{n} \Pi_{0}+\frac{n+1}{n} \Pi_{1}  \tag{6.49}\\
& \left(\mathrm{~T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\Pi_{0}\right)=\frac{n+1}{2} \Delta_{0}+\frac{1}{2} \Delta_{1}  \tag{6.50}\\
& \left(\mathrm{~T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\Pi_{1}\right)=-\frac{n-1}{2} \Delta_{0}+\frac{1}{2} \Delta_{1} \tag{6.51}
\end{align*}
$$

For $\lambda \in[0,1]$, the equations

$$
\begin{align*}
& \left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\lambda \Delta_{0}+(1-\lambda) \frac{\Delta_{1}}{n^{2}-1}\right) \\
& \quad=\left(\frac{1+\lambda n}{2}\right) \frac{\Pi_{0}}{\binom{n+1}{2}}+\left(\frac{1-\lambda n}{2}\right) \frac{\Pi_{1}}{\binom{n}{2}} \tag{6.52}
\end{align*}
$$

and

$$
\begin{align*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(\lambda & \left.\frac{\Pi_{0}}{\binom{n+1}{2}}+(1-\lambda) \frac{\Pi_{1}}{\binom{n}{2}}\right)  \tag{6.53}\\
& =\left(\frac{2 \lambda-1}{n}\right) \Delta_{0}+\left(1-\frac{2 \lambda-1}{n}\right) \frac{\Delta_{1}}{n^{2}-1}
\end{align*}
$$

are implied. It therefore holds that the isotropic state (6.45) is entangled (i.e., not separable) for $\lambda \in(1 / n, 1]$, while the Werner state (6.46) is entangled for $\lambda \in[0,1 / 2) .{ }^{1}$

1 It does indeed hold that the isotropic state (6.45) is separable for $\lambda \in[0,1 / n]$ and the Werner state (6.46) is separable for $\lambda \in[1 / 2,1]$. These facts are proved in Chapter 7 (q.v. Example 7.25).

A separable neighborhood of the identity operator
By means of the Horodecki criterion (Theorem 6.9), it may be proved that there exists a neighborhood of the identity operator $\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}$, for any choice of complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, in which every positive semidefinite operator is separable. Consequently, every density operator $\mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ that is sufficiently close to the completely mixed state is separable. In order to prove this fact, which is stated in more precise terms in Theorem 6.13 below, the following lemma will be used.

Lemma 6.11 Let $\Sigma$ be an alphabet, let $\mathcal{X}$ be a complex Euclidean space, let $\left\{X_{a, b}: a, b \in \Sigma\right\} \subset \mathrm{L}(\mathcal{X})$ be a collection of operators, and let $\mathcal{Y}=\mathbb{C}^{\Sigma}$. The operator

$$
\begin{equation*}
X=\sum_{a, b \in \Sigma} X_{a, b} \otimes E_{a, b} \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.54}
\end{equation*}
$$

satisfies

$$
\begin{equation*}
\|X\|^{2} \leq \sum_{a, b \in \Sigma}\left\|X_{a, b}\right\|^{2} \tag{6.55}
\end{equation*}
$$

Proof For each $a \in \Sigma$, define an operator $Y_{a} \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$ as

$$
\begin{equation*}
Y_{a}=\sum_{b \in \Sigma} X_{a, b} \otimes E_{a, b} \tag{6.56}
\end{equation*}
$$

By expanding the product $Y_{a} Y_{a}^{*}$ and applying the triangle inequality, the multiplicativity of the spectral norm under tensor products, and the spectral norm identity (1.178), one finds that

$$
\begin{equation*}
\left\|Y_{a} Y_{a}^{*}\right\|=\left\|\sum_{b \in \Sigma} X_{a, b} X_{a, b}^{*} \otimes E_{a, a}\right\| \leq \sum_{b \in \Sigma}\left\|X_{a, b} X_{a, b}^{*}\right\|=\sum_{b \in \Sigma}\left\|X_{a, b}\right\|^{2} \tag{6.57}
\end{equation*}
$$

Also observe that

$$
\begin{equation*}
X^{*} X=\sum_{a \in \Sigma} Y_{a}^{*} Y_{a} \tag{6.58}
\end{equation*}
$$

Therefore, by (6.57) together with the triangle inequality and the spectral norm identity (1.178), it holds that

$$
\begin{equation*}
\|X\|^{2}=\left\|X^{*} X\right\| \leq \sum_{a \in \Sigma}\left\|Y_{a}^{*} Y_{a}\right\| \leq \sum_{a, b \in \Sigma}\left\|X_{a, b}\right\|^{2} \tag{6.59}
\end{equation*}
$$

as required.
In addition, the following theorem (which is equivalent to Theorem 3.39) will be needed.

Theorem 6.12 Let $\Phi \in \mathrm{T}(\mathcal{X}, \mathcal{Y})$ be a positive and unital map, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$. It holds that

$$
\begin{equation*}
\|\Phi(X)\| \leq\|X\| \tag{6.60}
\end{equation*}
$$

for every operator $X \in \mathrm{~L}(\mathcal{X})$.
Proof By the assumption that $\Phi$ is positive and unital, Proposition 2.18 and Theorem 2.26 imply that $\Phi^{*}$ is positive and trace-preserving. For operators $X \in \mathrm{~L}(\mathcal{X})$ and $Y \in \mathrm{~L}(\mathcal{Y})$, one therefore has

$$
\begin{align*}
|\langle Y, \Phi(X)\rangle|= & \left|\left\langle\Phi^{*}(Y), X\right\rangle\right| \leq\|X\|\left\|\Phi^{*}(Y)\right\|_{1}  \tag{6.61}\\
& \leq\|X\|\|Y\|_{1}\left\|\Phi^{*}\right\|_{1}=\|X\|\|Y\|_{1}
\end{align*}
$$

where the final equality follows by Corollary 3.40 (to Theorem 3.39). By maximizing over all operators $Y \in \mathrm{~L}(\mathcal{Y})$ that satisfy $\|Y\|_{1} \leq 1$, one finds that $\|\Phi(X)\| \leq\|X\|$ for every $X \in \mathrm{~L}(\mathcal{X})$, as required.

Theorem 6.13 Let $H \in \operatorname{Herm}(\mathcal{X} \otimes \mathcal{Y})$ be a Hermitian operator satisfying $\|H\|_{2} \leq 1$, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$. It holds that

$$
\begin{equation*}
\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-H \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y}) \tag{6.62}
\end{equation*}
$$

Proof Let $\Phi \in \mathrm{T}(\mathcal{X}, \mathcal{Y})$ be an arbitrarily chosen positive and unital map. Let $\Sigma$ be the alphabet for which $\mathcal{Y}=\mathbb{C}^{\Sigma}$, and write

$$
\begin{equation*}
H=\sum_{a, b \in \Sigma} H_{a, b} \otimes E_{a, b} \tag{6.63}
\end{equation*}
$$

It holds that

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(H)=\sum_{a, b \in \Sigma} \Phi\left(H_{a, b}\right) \otimes E_{a, b} \tag{6.64}
\end{equation*}
$$

and therefore

$$
\begin{align*}
&\left\|\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(H)\right\|^{2} \leq \sum_{a, b \in \Sigma}\left\|\Phi\left(H_{a, b}\right)\right\|^{2} \\
& \quad \leq \sum_{a, b \in \Sigma}\left\|H_{a, b}\right\|^{2} \leq \sum_{a, b \in \Sigma}\left\|H_{a, b}\right\|_{2}^{2}=\|H\|_{2}^{2} \leq 1 \tag{6.65}
\end{align*}
$$

(The first inequality is implied by Lemma 6.11, and the second inequality is implied by Theorem 6.12.) The positivity of $\Phi$ implies that $\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(H)$ is Hermitian, and therefore $\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(H) \leq \mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}$. It follows that

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-H\right)=\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(H) \geq 0 \tag{6.66}
\end{equation*}
$$

Because (6.66) holds for all positive and unital maps $\Phi$, one concludes from Theorem 6.9 that $\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-H$ is separable.

## Bipartite operator entanglement rank

Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, and consider the collection of all positive semidefinite operators $R \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ for which there exists an alphabet $\Sigma$ and a collection of operators $\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{X})$ such that

$$
\begin{equation*}
R=\sum_{a \in \Sigma} \operatorname{vec}\left(A_{a}\right) \operatorname{vec}\left(A_{a}\right)^{*} \tag{6.67}
\end{equation*}
$$

and $\operatorname{rank}\left(A_{a}\right) \leq 1$ for each $a \in \Sigma$. An operator $A \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$ has rank at most 1 if and only if there exist vectors $u \in \mathcal{X}$ and $v \in \mathcal{Y}$ such that $\operatorname{vec}(A)=u \otimes v$, and from this observation it follows that the collection of operators $R$ just described coincides with $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$.

It is useful to generalize this notion, allowing for arbitrary upper-bounds on the rank of the operators $\left\{A_{a}: a \in \Sigma\right\}$, along the lines of the following definition.

Definition 6.14 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces and let $r \geq 1$ be a positive integer. The set $\operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y})$ is defined to be the set of all operators $R \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ for which there exists an alphabet $\Sigma$ and a collection of operators

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{X}) \tag{6.68}
\end{equation*}
$$

satisfying $\operatorname{rank}\left(A_{a}\right) \leq r$ for each $a \in \Sigma$, such that

$$
\begin{equation*}
R=\sum_{a \in \Sigma} \operatorname{vec}\left(A_{a}\right) \operatorname{vec}\left(A_{a}\right)^{*} \tag{6.69}
\end{equation*}
$$

An element $R \in \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y})$ is said to have entanglement rank bounded by $r$. The entanglement rank of $R \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$, with respect to the bipartition between $\mathcal{X}$ and $\mathcal{Y}$, is the minimum value of $r \geq 1$ such that $R \in \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y})$.

As indicated above, it holds that

$$
\begin{equation*}
\operatorname{Sep}(\mathcal{X}: \mathcal{Y})=\operatorname{Ent}_{1}(\mathcal{X}: \mathcal{Y}) \tag{6.70}
\end{equation*}
$$

and from Definition 6.14 it is immediate that

$$
\begin{equation*}
\operatorname{Ent}_{r-1}(\mathcal{X}: \mathcal{Y}) \subseteq \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y}) \tag{6.71}
\end{equation*}
$$

for every integer $r \geq 2$.
The containment (6.71) is proper, provided $r \leq \min \{\operatorname{dim}(\mathcal{X}), \operatorname{dim}(\mathcal{Y})\}$. To see that this is so, consider any operator $B \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$ having rank equal to $r$, and suppose that

$$
\begin{equation*}
\operatorname{vec}(B) \operatorname{vec}(B)^{*}=\sum_{a \in \Sigma} \operatorname{vec}\left(A_{a}\right) \operatorname{vec}\left(A_{a}\right)^{*} \tag{6.72}
\end{equation*}
$$

for some collection of operators $\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{X})$. As the operator represented by this equation has rank equal to 1 , it must hold that $A_{a}=\alpha_{a} B$ for each $a \in \Sigma$, for $\left\{\alpha_{a}: a \in \Sigma\right\}$ being a collection of complex numbers satisfying

$$
\begin{equation*}
\sum_{a \in \Sigma}\left|\alpha_{a}\right|^{2}=1 \tag{6.73}
\end{equation*}
$$

It is therefore not possible that (6.72) holds when each operator $A_{a}$ has rank strictly smaller than $r$, and therefore

$$
\begin{equation*}
\operatorname{vec}(B) \operatorname{vec}(B)^{*} \notin \operatorname{Ent}_{r-1}(\mathcal{X}: \mathcal{Y}) \tag{6.74}
\end{equation*}
$$

It is immediate, on the other hand, that $\operatorname{vec}(B) \operatorname{vec}(B)^{*} \in \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y})$.
Finally, one may observe that

$$
\begin{equation*}
\operatorname{Ent}_{n}(\mathcal{X}: \mathcal{Y})=\operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.75}
\end{equation*}
$$

for $n \geq \min \{\operatorname{dim}(\mathcal{X}), \operatorname{dim}(\mathcal{Y})\}$, as every operator $A \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$ has rank bounded by $n$ in this case.

The following simple proposition concerning entanglement rank will be useful in subsequent sections of this chapter.

Proposition 6.15 Let $B \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$ be an operator, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, and assume that $\|B\| \leq 1$. For every positive integer $r$ and every operator

$$
\begin{equation*}
P \in \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y}) \tag{6.76}
\end{equation*}
$$

having entanglement rank bounded by $r$, it holds that

$$
\begin{equation*}
\left\langle\operatorname{vec}(B) \operatorname{vec}(B)^{*}, P\right\rangle \leq r \operatorname{Tr}(P) \tag{6.77}
\end{equation*}
$$

Proof Under the assumption that $P$ has entanglement rank bounded by $r$, one may write

$$
\begin{equation*}
P=\sum_{a \in \Sigma} \operatorname{vec}\left(A_{a}\right) \operatorname{vec}\left(A_{a}\right)^{*} \tag{6.78}
\end{equation*}
$$

for an alphabet $\Sigma$ and a collection of operators $\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{X})$ for which $\operatorname{rank}\left(A_{a}\right) \leq r$ for every $a \in \Sigma$. For every operator $A \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$, one has

$$
\begin{equation*}
|\langle B, A\rangle|^{2} \leq\|A\|_{1}^{2} \leq \operatorname{rank}(A)\|A\|_{2}^{2} \tag{6.79}
\end{equation*}
$$

so that evaluating the inner product in the statement of the proposition
yields

$$
\begin{align*}
& \left\langle\operatorname{vec}(B) \operatorname{vec}(B)^{*}, P\right\rangle=\sum_{a \in \Sigma}\left|\left\langle B, A_{a}\right\rangle\right|^{2} \\
& \quad \leq \sum_{a \in \Sigma} \operatorname{rank}\left(A_{a}\right)\left\|A_{a}\right\|_{2}^{2} \leq r \sum_{a \in \Sigma}\left\|A_{a}\right\|_{2}^{2}=r \operatorname{Tr}(P) \tag{6.80}
\end{align*}
$$

as required.
Example 6.16 Let $\Sigma$ be an alphabet, let $n=|\Sigma|$, let $\mathcal{X}=\mathbb{C}^{\Sigma}$ and $\mathcal{Y}=\mathbb{C}^{\Sigma}$, and define a density operator $\tau \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ as

$$
\begin{equation*}
\tau=\frac{1}{n} \sum_{a, b \in \Sigma} E_{a, b} \otimes E_{a, b} \tag{6.81}
\end{equation*}
$$

The density operator $\tau$, which coincides with the isotropic state $\Delta_{0}$ defined in Example 6.10, is the canonical example of a maximally entangled state with respect to the spaces $\mathcal{X}$ and $\mathcal{Y}$. One may observe that

$$
\begin{equation*}
\tau=\frac{1}{n} \operatorname{vec}(\mathbb{1}) \operatorname{vec}(\mathbb{1})^{*} \tag{6.82}
\end{equation*}
$$

for $\mathbb{1}$ denoting the identity operator on $\mathbb{C}^{\Sigma}$, which may be viewed as an element of the set $\mathrm{L}(\mathcal{Y}, \mathcal{X})$ in the most straightforward way.

For every positive integer $r$ and every density operator

$$
\begin{equation*}
\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y}) \cap \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y}) \tag{6.83}
\end{equation*}
$$

having entanglement rank bounded by $r$, Proposition 6.15 implies that

$$
\begin{equation*}
\langle\tau, \rho\rangle=\frac{1}{n}\left\langle\operatorname{vec}(\mathbb{1}) \operatorname{vec}(\mathbb{1})^{*}, \rho\right\rangle \leq \frac{r}{n} \tag{6.84}
\end{equation*}
$$

One therefore has that every state of bounded entanglement rank must have a proportionately small inner product with the state $\tau$.

### 6.1.2 Separable maps and the LOCC paradigm

Separable maps are defined in an analogous way to separable operators, reflecting the natural correspondence between completely positive maps and positive semidefinite operators. The resulting notion of separability for maps, including channels, is algebraic in nature; and it cannot be said that it is directly motivated from a physical or operational viewpoint.

This notion of separability for channels is, however, closely connected to the more operationally motivated notion of channels implementable by local operations and classical communication, or LOCC for short. An LOCC channel is a channel that can be implemented by two individuals whose
local actions are unrestricted (corresponding to arbitrary measurements or channels), but whose communications with one another are restricted to be classical. This paradigm provides a foundation from which properties of entanglement are commonly studied, particularly in settings in which entanglement is viewed as a resource for information processing.

## Separable maps and channels

As suggested above, the notion of separability for maps is defined in an analogous way to separability for operators. The following definition states this in more precise terms.

Definition 6.17 Let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$ be complex Euclidean spaces. The set $\operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ is defined as the set of all completely positive maps of the form

$$
\begin{equation*}
\Xi \in \mathrm{CP}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W}) \tag{6.85}
\end{equation*}
$$

for which there exists an alphabet $\Sigma$ and collections of completely positive $\operatorname{maps}\left\{\Phi_{a}: a \in \Sigma\right\} \subset \mathrm{CP}(\mathcal{X}, \mathcal{Z})$ and $\left\{\Psi_{a}: a \in \Sigma\right\} \subset \mathrm{CP}(\mathcal{Y}, \mathcal{W})$ such that

$$
\begin{equation*}
\Xi=\sum_{a \in \Sigma} \Phi_{a} \otimes \Psi_{a} \tag{6.86}
\end{equation*}
$$

Elements of the set $\operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ are called separable maps.
As the following simple proposition states, separable maps are precisely those completely positive maps having Kraus representations for which the individual Kraus operators are tensor products of operators. A direct proof of this proposition is obtained by considering Kraus representations of the maps $\Phi_{a}$ and $\Psi_{a}$ in Definition 6.17, along the same lines as the proof of Proposition 6.5.

Proposition 6.18 Let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$ be complex Euclidean spaces and let $\Xi \in \mathrm{CP}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W})$ be a completely positive map. It holds that

$$
\begin{equation*}
\Phi \in \operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \tag{6.87}
\end{equation*}
$$

if and only if there exists an alphabet $\Sigma$ and collections of operators

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{X}, \mathcal{Z}) \quad \text { and } \quad\left\{B_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{W}) \tag{6.88}
\end{equation*}
$$

such that

$$
\begin{equation*}
\Xi(X)=\sum_{a \in \Sigma}\left(A_{a} \otimes B_{a}\right) X\left(A_{a} \otimes B_{a}\right)^{*} \tag{6.89}
\end{equation*}
$$

for every operator $X \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$.

Another straightforward proposition regarding separable maps follows, and like the previous proposition, it may be verified directly. It implies that the set of all separable maps is closed under composition.

Proposition 6.19 Let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}, \mathcal{W}, \mathcal{U}$, and $\mathcal{V}$ be complex Euclidean spaces, and suppose that $\Phi$ and $\Psi$ are separable maps of the form

$$
\begin{equation*}
\Phi \in \operatorname{SepCP}(\mathcal{X}, \mathcal{U}: \mathcal{Y}, \mathcal{V}) \quad \text { and } \quad \Psi \in \operatorname{SepCP}(\mathcal{U}, \mathcal{Z}: \mathcal{V}, \mathcal{W}) \tag{6.90}
\end{equation*}
$$

It holds that the composition $\Psi \Phi$ is separable:

$$
\begin{equation*}
\Psi \Phi \in \operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \tag{6.91}
\end{equation*}
$$

Similar to the analogous case for states, one defines the set of separable channels by simply restricting the definition of separability for completely positive maps to channels.

Definition 6.20 For complex Euclidean spaces $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$, one defines

$$
\begin{align*}
& \operatorname{SepC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \\
& \quad=\operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \cap \mathrm{C}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W}) \tag{6.92}
\end{align*}
$$

Elements of the set $\operatorname{SepC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ are referred to as separable channels.
It should be noted that, unlike the analogous case of states, separable channels need not be equal to convex combinations of product channels, as the following example illustrates.

Example 6.21 Let $\Sigma=\{0,1\}$ denote the binary alphabet, let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$ all be equal to $\mathbb{C}^{\Sigma}$, and define a channel $\Xi \in \mathrm{C}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W})$ by the equation

$$
\Xi\left(E_{a, b} \otimes E_{c, d}\right)= \begin{cases}E_{a, a} \otimes E_{a, a} & \text { if } a=b \text { and } c=d  \tag{6.93}\\ 0 & \text { if } a \neq b \text { or } c \neq d\end{cases}
$$

holding for all $a, b, c, d \in \Sigma$. It is the case that $\Xi$ is a separable channel, meaning that $\Xi \in \operatorname{SepC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$. Indeed, one may write

$$
\begin{equation*}
\Xi=\Phi_{0} \otimes \Psi_{0}+\Phi_{1} \otimes \Psi_{1} \tag{6.94}
\end{equation*}
$$

for completely positive maps defined as follows:

$$
\begin{array}{ll}
\Phi_{0}(X)=\left\langle E_{0,0}, X\right\rangle E_{0,0}, & \Psi_{0}(X)=\operatorname{Tr}(X) E_{0,0}  \tag{6.95}\\
\Phi_{1}(X)=\left\langle E_{1,1}, X\right\rangle E_{1,1}, & \Psi_{1}(X)=\operatorname{Tr}(X) E_{1,1}
\end{array}
$$

for every $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$.

It is not possible, however, to express the channel $\Xi$ in the form

$$
\begin{equation*}
\Xi=\sum_{a \in \Gamma} p(a) \Phi_{a} \otimes \Psi_{a} \tag{6.96}
\end{equation*}
$$

for any choice of an alphabet $\Gamma$, a probability vector $p \in \mathcal{P}(\Gamma)$, and two collections of channels

$$
\begin{equation*}
\left\{\Phi_{a}: a \in \Gamma\right\} \subset \mathrm{C}(\mathcal{X}, \mathcal{Z}) \quad \text { and } \quad\left\{\Psi_{a}: a \in \Gamma\right\} \subset \mathrm{C}(\mathcal{Y}, \mathcal{W}) \tag{6.97}
\end{equation*}
$$

To verify this claim, consider the fact that

$$
\begin{equation*}
\Xi\left(E_{0,0} \otimes \rho\right)=E_{0,0} \otimes E_{0,0} \quad \text { and } \quad \Xi\left(E_{1,1} \otimes \rho\right)=E_{1,1} \otimes E_{1,1} \tag{6.98}
\end{equation*}
$$

for every density operator $\rho \in \mathrm{D}(\mathcal{Y})$. If it were the case that (6.96) were true for each $\Phi_{a}$ and $\Psi_{a}$ being a channel, then one would necessarily have

$$
\begin{equation*}
\sum_{a \in \Gamma} p(a) \Phi_{a}\left(E_{0,0}\right) \otimes \Psi_{a}(\rho)=E_{0,0} \otimes E_{0,0} \tag{6.99}
\end{equation*}
$$

and therefore, by tracing over the space $\mathcal{Z}$,

$$
\begin{equation*}
\sum_{a \in \Sigma} p(a) \Psi_{a}(\rho)=E_{0,0} \tag{6.100}
\end{equation*}
$$

for every $\rho \in \mathrm{D}(\mathcal{Y})$. By similar reasoning, it would simultaneously hold that

$$
\begin{equation*}
\sum_{a \in \Sigma} p(a) \Phi_{a}\left(E_{1,1}\right) \otimes \Psi_{a}(\rho)=E_{1,1} \otimes E_{1,1} \tag{6.101}
\end{equation*}
$$

and therefore

$$
\begin{equation*}
\sum_{a \in \Sigma} p(a) \Psi_{a}(\rho)=E_{1,1} \tag{6.102}
\end{equation*}
$$

for every $\rho \in \mathrm{D}(\mathcal{Y})$. The equations (6.100) and (6.102) are in contradiction, implying that $\Xi$ is not equal to a convex combination of product channels.

Intuitively speaking, the situation represented by the previous example is quite simple. Channels that can be expressed as a convex combination of product channels correspond to transformations that may be implemented by means of local operations and shared randomness-no communication is needed to implement them, and such channels do not allow for a direct causal relationship to hold among the input and output systems across the bipartition with respect to which separability is considered. The channel $\Xi$, on the other hand, induces a direct causal relationship of this form.

As the following proposition states, a given completely positive map is separable if and only if its Choi representation is separable, with respect to the natural bipartition of the tensor product space over which it is defined.

Proposition 6.22 Let $\Xi \in \mathrm{CP}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W})$ be a completely positive map, for complex Euclidean spaces $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$, and define an isometry

$$
\begin{equation*}
V \in \mathrm{U}(\mathcal{Z} \otimes \mathcal{W} \otimes \mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{X} \otimes \mathcal{W} \otimes \mathcal{Y}) \tag{6.103}
\end{equation*}
$$

by the equation

$$
\begin{equation*}
V \operatorname{vec}(A \otimes B)=\operatorname{vec}(A) \otimes \operatorname{vec}(B) \tag{6.104}
\end{equation*}
$$

holding for all operators $A \in \mathrm{~L}(\mathcal{X}, \mathcal{Z})$ and $B \in \mathrm{~L}(\mathcal{Y}, \mathcal{W})$. It holds that

$$
\begin{equation*}
\Xi \in \operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \tag{6.105}
\end{equation*}
$$

if and only if

$$
\begin{equation*}
V J(\Xi) V^{*} \in \operatorname{Sep}(\mathcal{Z} \otimes \mathcal{X}: \mathcal{W} \otimes \mathcal{Y}) \tag{6.106}
\end{equation*}
$$

Proof Assume first that $\Xi$ is a separable map. By Proposition 6.18, there must exist an alphabet $\Sigma$ and two collections of operators,

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{X}, \mathcal{Z}) \quad \text { and } \quad\left\{B_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{W}) \tag{6.107}
\end{equation*}
$$

such that

$$
\begin{equation*}
\Xi(X)=\sum_{a \in \Sigma}\left(A_{a} \otimes B_{a}\right) X\left(A_{a} \otimes B_{a}\right)^{*} \tag{6.108}
\end{equation*}
$$

for every operator $X \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$. The Choi representation of $\Xi$ is therefore given by

$$
\begin{equation*}
J(\Xi)=\sum_{a \in \Sigma} \operatorname{vec}\left(A_{a} \otimes B_{a}\right) \operatorname{vec}\left(A_{a} \otimes B_{a}\right)^{*} \tag{6.109}
\end{equation*}
$$

so that

$$
\begin{equation*}
V J(\Xi) V^{*}=\sum_{a \in \Sigma} \operatorname{vec}\left(A_{a}\right) \operatorname{vec}\left(A_{a}\right)^{*} \otimes \operatorname{vec}\left(B_{a}\right) \operatorname{vec}\left(B_{a}\right)^{*} \tag{6.110}
\end{equation*}
$$

which is evidently contained in $\operatorname{Sep}(\mathcal{Z} \otimes \mathcal{X}: \mathcal{W} \otimes \mathcal{Y})$.
Conversely, if $\operatorname{VJ}(\Xi) V^{*}$ is separable, then it must be possible to express this operator in the form (6.110) for some choice of an alphabet $\Sigma$ and two collections of operators as in (6.107). It therefore follows that (6.109) is a Choi representation of $\Xi$, so that (6.108) holds for all $X \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$. The map $\Xi$ is therefore separable, which completes the proof.

Remark The isometry $V$ defined in Proposition 6.22 may alternatively be defined by the action

$$
\begin{equation*}
V(z \otimes w \otimes x \otimes y)=z \otimes x \otimes w \otimes y \tag{6.111}
\end{equation*}
$$

for every choice of vectors $x \in \mathcal{X}, y \in \mathcal{Y}, z \in \mathcal{Z}$, and $w \in \mathcal{W}$. In words, this
isometry represents a permutation of tensor factors, allowing a relationship concerning separability with respect to a particular bipartition to be stated precisely.

It is not uncommon in the theory of quantum information literature that statements of this nature are made without an explicit mention of such an isometry. This can sometimes simplify expressions and generally does not lead to any confusion - the isometry can usually be taken as being implicit, particularly in cases when the underlying complex Euclidean spaces have distinct names. In the interest of clarity and formality, however, this book will always represent such permutations of tensor factors explicitly.

Separable channels are not capable of creating entanglement: a separable channel applied to a separable state yields another separable state. More generally, separable maps cannot cause an increase in entanglement rank, as the following theorem establishes.

Theorem 6.23 Let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$ be complex Euclidean spaces and let $\Xi \in \operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ be a separable map. For every positive integer $r$ and every operator $P \in \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y})$, it holds that $\Xi(P) \in \operatorname{Ent}_{r}(\mathcal{Z}: \mathcal{W})$.

Proof For an operator $P \in \operatorname{Ent}_{r}(\mathcal{X}: \mathcal{Y})$ having entanglement rank bounded by $r$, there must exist an alphabet $\Gamma$ and a collection of operators

$$
\begin{equation*}
\left\{X_{b}: b \in \Gamma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{X}) \tag{6.112}
\end{equation*}
$$

satisfying $\operatorname{rank}\left(X_{b}\right) \leq r$ for every $b \in \Gamma$, such that

$$
\begin{equation*}
P=\sum_{b \in \Gamma} \operatorname{vec}\left(X_{b}\right) \operatorname{vec}\left(X_{b}\right)^{*} \tag{6.113}
\end{equation*}
$$

By Proposition 6.18, it follows that

$$
\begin{align*}
\Xi(P) & =\sum_{a \in \Sigma} \sum_{b \in \Gamma}\left(A_{a} \otimes B_{a}\right) \operatorname{vec}\left(X_{b}\right) \operatorname{vec}\left(X_{b}\right)^{*}\left(A_{a} \otimes B_{a}\right)^{*} \\
& =\sum_{a \in \Sigma} \sum_{b \in \Gamma} \operatorname{vec}\left(A_{a} X_{b} B_{a}^{\top}\right) \operatorname{vec}\left(A_{a} X_{b} B_{a}^{\top}\right)^{*} \tag{6.114}
\end{align*}
$$

for some choice of an alphabet $\Sigma$ and two collections of operators

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{X}, \mathcal{Z}) \quad \text { and } \quad\left\{B_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{W}) \tag{6.115}
\end{equation*}
$$

For every $a \in \Sigma$ and $b \in \Gamma$, it holds that

$$
\begin{equation*}
\operatorname{rank}\left(A_{a} X_{b} B_{a}^{\top}\right) \leq \operatorname{rank}\left(X_{b}\right) \leq r \tag{6.116}
\end{equation*}
$$

and therefore $\Xi(P) \in \operatorname{Ent}_{r}(\mathcal{Z}: \mathcal{W})$, as required.

Corollary 6.24 Let $\Xi \in \operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ be a separable map, for complex Euclidean spaces $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$. For every separable operator $P \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$, it holds that $\Xi(P)$ is also separable: $\Xi(P) \in \operatorname{Sep}(\mathcal{Z}: \mathcal{W})$.

## LOCC channels

As was stated at the beginning of the present subsection, LOCC channels represent transformations of quantum states that may be implemented by two individuals that communicate with one another classically and perform quantum channels and measurements on registers they hold locally.

For instance, one individual may apply a combination of channels and measurements to a collection of registers in their possession and then transmit the measurement outcomes to the other individual. Upon receiving this transmission, the other individual may apply a combination of channels and measurements, possibly depending on the communicated measurement outcomes, to a collection of registers in their possession. In general, LOCC channels represent the cumulative effect of composing any finite number of transformations of this sort. ${ }^{2}$

The following definition formalizes this notion. Naturally, it is possible to generalize this definition to three or more individuals, although this will not be done in this book.

Definition 6.25 Let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$ be complex Euclidean spaces and let $\Xi \in \mathrm{C}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W})$ be a channel. The channel $\Xi$ is an $L O C C$ channel under these conditions:

1. If there exists an alphabet $\Sigma$ and a collection

$$
\begin{equation*}
\left\{\Phi_{a}: a \in \Sigma\right\} \subset \mathrm{CP}(\mathcal{X}, \mathcal{Z}) \tag{6.117}
\end{equation*}
$$

of completely positive maps satisfying

$$
\begin{equation*}
\sum_{a \in \Sigma} \Phi_{a} \in \mathrm{C}(\mathcal{X}, \mathcal{Z}) \tag{6.118}
\end{equation*}
$$

along with a collection

$$
\begin{equation*}
\left\{\Psi_{a}: a \in \Sigma\right\} \subseteq \mathrm{C}(\mathcal{Y}, \mathcal{W}) \tag{6.119}
\end{equation*}
$$

of channels, such that

$$
\begin{equation*}
\Xi=\sum_{a \in \Sigma} \Phi_{a} \otimes \Psi_{a} \tag{6.120}
\end{equation*}
$$

then $\Xi$ is a one-way right $L O C C$ channel.
${ }^{2}$ One may consider variants of the definition that allow for an unbounded number of classical transmissions that terminate with probability 1 according to a chosen stopping rule. Only the finite case is considered in this book for simplicity.
2. If there exists an alphabet $\Sigma$ and a collection

$$
\begin{equation*}
\left\{\Psi_{a}: a \in \Sigma\right\} \subset \mathrm{CP}(\mathcal{Y}, \mathcal{W}) \tag{6.121}
\end{equation*}
$$

of completely positive maps satisfying

$$
\begin{equation*}
\sum_{a \in \Sigma} \Psi_{a} \in \mathrm{C}(\mathcal{Y}, \mathcal{W}) \tag{6.122}
\end{equation*}
$$

along with a collection

$$
\begin{equation*}
\left\{\Phi_{a}: a \in \Sigma\right\} \subseteq \mathrm{C}(\mathcal{X}, \mathcal{Z}) \tag{6.123}
\end{equation*}
$$

of channels, such that (6.120) holds, then $\Xi$ is a one-way left LOCC channel.
3. The channel $\Xi$ is an $L O C C$ channel if it is equal to a finite composition of one-way left and one-way right LOCC channels. That is, either $\Xi$ is a one-way left LOCC channel, a one-way right LOCC channel, or there exists an integer $m \geq 2$, complex Euclidean spaces $\mathcal{U}_{1}, \ldots, \mathcal{U}_{m-1}$ and $\mathcal{V}_{1}, \ldots, \mathcal{V}_{m-1}$, and channels

$$
\begin{gather*}
\Xi_{1} \in \mathrm{C}\left(\mathcal{X} \otimes \mathcal{Y}, \mathcal{U}_{1} \otimes \mathcal{V}_{1}\right) \\
\Xi_{2} \in \mathrm{C}\left(\mathcal{U}_{1} \otimes \mathcal{V}_{1}, \mathcal{U}_{2} \otimes \mathcal{V}_{2}\right) \\
\vdots  \tag{6.124}\\
\Xi_{m} \in \mathrm{C}\left(\mathcal{U}_{m-1} \otimes \mathcal{V}_{m-1}, \mathcal{Z} \otimes \mathcal{W}\right)
\end{gather*}
$$

each of which is either a one-way left LOCC channel or a one-way right LOCC channel, such that $\Xi$ is equal to the composition $\Xi=\Xi_{m} \cdots \Xi_{1}$.

The collection of all such LOCC channels is denoted $\operatorname{LOCC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$.
Remark In the definition above, one-way left and one-way right LOCC channels represent channels that can be implemented by local operations and one-way classical communication. In both cases, the channel $\Xi$ may be viewed as having resulted from actions performed by two individuals, Alice and Bob. Alice begins with a register X and Bob begins with Y , and as a result of their actions these registers are transformed into $Z$ and $W$, respectively.

In the case of a one-way right LOCC channel $\Xi$, the communication is from Alice to Bob (moving to the right, assuming Alice is on the left and Bob is on the right), with the alphabet $\Sigma$ representing the set of possible classical messages that may be transmitted. Alice's actions are described by a collection of completely positive maps

$$
\begin{equation*}
\left\{\Phi_{a}: a \in \Sigma\right\} \subset \operatorname{CP}(\mathcal{X}, \mathcal{Z}) \tag{6.125}
\end{equation*}
$$

that satisfies the constraint

$$
\begin{equation*}
\sum_{a \in \Sigma} \Phi_{a} \in \mathrm{C}(\mathcal{X}, \mathcal{Z}) \tag{6.126}
\end{equation*}
$$

In essence, this collection specifies an instrument (q.v. Section 2.3.2). Assuming the classical communication is represented by a classical register V having associated complex Euclidean space $\mathcal{V}=\mathbb{C}^{\Sigma}$, Alice's action would be described by the channel $\Phi \in \mathrm{C}(\mathcal{X}, \mathcal{Z} \otimes \mathcal{V})$ defined by

$$
\begin{equation*}
\Phi(X)=\sum_{a \in \Sigma} \Phi_{a}(X) \otimes E_{a, a} \tag{6.127}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X})$. The register V is sent to Bob, who observes its classical state (or, equivalently, measures V with respect to the standard basis) and transforms his register Y into W according to the channel $\Psi_{a} \in \mathrm{C}(\mathcal{Y}, \mathcal{W})$, for $a \in \Sigma$ being the classical state of V that was observed. Assuming that the register V is discarded after Bob applies the appropriate channel, the combined actions of Alice and Bob are described by $\Xi$.

For a one-way left LOCC channel $\Xi$, the situation is similar, with the roles of Alice and Bob switched.

It is apparent from Definition 6.25 , together with the fact that separable channels are closed under composition (Proposition 6.19), that every LOCC channel is a separable channel.

Proposition 6.26 For every choice of complex Euclidean spaces $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$, it holds that

$$
\begin{equation*}
\operatorname{LOCC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \subseteq \operatorname{SepC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \tag{6.128}
\end{equation*}
$$

### 6.1.3 Separable and LOCC measurements

As was explained in Section 2.3.1, one may associate a quantum-to-classical channel with each measurement, with the classical output of the channel representing the outcome of the measurement. Through an identification of this sort, the notions of separable and LOCC channels may be extended to measurements.

## Definitions of separable and LOCC measurements

The following definition of separable and LOCC measurements refers to an association of quantum-to-classical channels with measurements that has been adapted to a bipartite setting.

Definition 6.27 Let $\Sigma$ be an alphabet, let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, and let $\mu: \Sigma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ be a measurement. Define complex Euclidean spaces $\mathcal{Z}=\mathbb{C}^{\Sigma}$ and $\mathcal{W}=\mathbb{C}^{\Sigma}$, and define a channel

$$
\begin{equation*}
\Phi_{\mu} \in \mathrm{C}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W}) \tag{6.129}
\end{equation*}
$$

as

$$
\begin{equation*}
\Phi_{\mu}(X)=\sum_{a \in \Sigma}\langle\mu(a), X\rangle E_{a, a} \otimes E_{a, a} \tag{6.130}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$. The measurement $\mu$ is a separable measurement if

$$
\begin{equation*}
\Phi_{\mu} \in \operatorname{SepC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \tag{6.131}
\end{equation*}
$$

and $\mu$ is an LOCC measurement if

$$
\begin{equation*}
\Phi_{\mu} \in \operatorname{LOCC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W}) \tag{6.132}
\end{equation*}
$$

For a given measurement $\mu$, the channel $\Phi_{\mu}$ specified in Definition 6.27 is similar to the quantum-to-classical channel one would normally associate with $\mu$, except that two copies of the measurement outcome are produced rather than one. In a bipartite setting, this is a natural way of associating a quantum-to-classical channel with a measurement. If this measurement is performed on a pair of registers ( $\mathrm{X}, \mathrm{Y}$ ) by two individuals, Alice and Bob, where it is assumed that Alice holds X and Bob holds Y , the channel $\Phi_{\mu}$ represents the measurement $\mu$ under the assumption that both individuals learn the measurement outcome after the measurement is performed.

One alternative to Definition 6.27 is to replace the channel $\Phi_{\mu}$ by the quantum-to-classical channel that would ordinarily be associated with the measurement $\mu$, along with a specification of which side of the bipartition the measurement outcome is to fall (requiring this channel to be separable or LOCC, as in the stated definition). In essence, with respect to a situation in which Alice and Bob are performing the measurement $\mu$ as suggested above, such a definition specifies which of the individuals obtains the measurement outcome. This alternative creates an asymmetry in the definition, but is equivalent to Definition 6.27.

With respect to Definition 6.27, the separability of a given measurement is equivalent to the constraint that each measurement operator is separable, as the following proposition states.

Proposition 6.28 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, let $\Sigma$ be an alphabet, and let $\mu$ be a measurement of the form $\mu: \Sigma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$. It holds that $\mu$ is a separable measurement if and only if $\mu(a) \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ for every $a \in \Sigma$.

Proof Consider the Choi representation of the mapping $\Phi_{\mu}$, as specified in Definition 6.27, which is given by

$$
\begin{equation*}
J\left(\Phi_{\mu}\right)=\sum_{a \in \Sigma} E_{a, a} \otimes E_{a, a} \otimes \overline{\mu(a)} \tag{6.133}
\end{equation*}
$$

Along similar lines to the statement of Proposition 6.22, let

$$
\begin{equation*}
V \in \mathrm{U}(\mathcal{Z} \otimes \mathcal{W} \otimes \mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{X} \otimes \mathcal{W} \otimes \mathcal{Y}) \tag{6.134}
\end{equation*}
$$

be the isometry defined by the equation

$$
\begin{equation*}
V \operatorname{vec}(A \otimes B)=\operatorname{vec}(A) \otimes \operatorname{vec}(B) \tag{6.135}
\end{equation*}
$$

holding for all operators $A \in \mathrm{~L}(\mathcal{X}, \mathcal{Z})$ and $B \in \mathrm{~L}(\mathcal{Y}, \mathcal{W})$. If it is the case that $\mu(a) \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ for every $a \in \Sigma$, then it follows directly that

$$
\begin{equation*}
V J\left(\Phi_{\mu}\right) V^{*} \in \operatorname{Sep}(\mathcal{Z} \otimes \mathcal{X}: \mathcal{W} \otimes \mathcal{Y}) \tag{6.136}
\end{equation*}
$$

which implies that $\mu$ is a separable measurement by Proposition 6.22.
Now suppose that $\mu$ is a separable measurement, so that (6.136) holds. Define a mapping

$$
\begin{equation*}
\Xi_{a} \in \mathrm{~T}(\mathcal{Z} \otimes \mathcal{X} \otimes \mathcal{W} \otimes \mathcal{Y}, \mathcal{X} \otimes \mathcal{Y}) \tag{6.137}
\end{equation*}
$$

for each $a \in \Sigma$, as

$$
\begin{equation*}
\Xi_{a}(X)=\left(\left(e_{a}^{*} \otimes \mathbb{1}_{\mathcal{X}}\right) \otimes\left(e_{a}^{*} \otimes \mathbb{1}_{\mathcal{Y}}\right)\right) X\left(\left(e_{a} \otimes \mathbb{1}_{\mathcal{X}}\right) \otimes\left(e_{a} \otimes \mathbb{1}_{\mathcal{Y}}\right)\right) \tag{6.138}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{Z} \otimes \mathcal{X} \otimes \mathcal{W} \otimes \mathcal{Y})$. It is evident from this definition that $\Xi_{a}$ is a separable mapping for each $a \in \Sigma$, meaning

$$
\begin{equation*}
\Xi_{a} \in \operatorname{SepCP}(\mathcal{Z} \otimes \mathcal{X}, \mathcal{X}: \mathcal{W} \otimes \mathcal{Y}, \mathcal{Y}) \tag{6.139}
\end{equation*}
$$

It holds that

$$
\begin{equation*}
\overline{\mu(a)}=\Xi_{a}\left(V J\left(\Phi_{\mu}\right) V^{*}\right) \tag{6.140}
\end{equation*}
$$

for each $a \in \Sigma$, from which it follows that

$$
\begin{equation*}
\overline{\mu(a)} \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y}) \tag{6.141}
\end{equation*}
$$

by Corollary 6.24. This is equivalent to $\mu(a) \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ for each $a \in \Sigma$, as the entry-wise complex conjugate of every separable operator is evidently separable, which completes the proof.

For two complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, along with an alphabet $\Sigma$, it is the case that the set of all separable measurements of the form

$$
\begin{equation*}
\mu: \Sigma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.142}
\end{equation*}
$$

is a proper subset of the set of all measurements of the same form (aside from the trivial cases in which one of $\operatorname{dim}(\mathcal{X}), \operatorname{dim}(\mathcal{Y})$, or $|\Sigma|$ equals 1). As every LOCC channel is separable, it follows that every LOCC measurement is a separable measurement.

## One-way LOCC measurements

An interesting restricted type of LOCC measurement is one in which only one-way communication is permitted. The following definition formalizes this type of measurement.

Definition 6.29 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, let $\Sigma$ be an alphabet, and let

$$
\begin{equation*}
\mu: \Sigma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.143}
\end{equation*}
$$

be a measurement. The measurement $\mu$ is a one-way LOCC measurement if either of the following two conditions is met:

1. There exists an alphabet $\Gamma$ and a measurement $\nu: \Gamma \rightarrow \operatorname{Pos}(\mathcal{X})$, along with a measurement $\pi_{b}: \Sigma \rightarrow \operatorname{Pos}(\mathcal{Y})$ for each $b \in \Gamma$, such that the equation

$$
\begin{equation*}
\mu(a)=\sum_{b \in \Gamma} \nu(b) \otimes \pi_{b}(a) \tag{6.144}
\end{equation*}
$$

holds for every $a \in \Sigma$. In this case the measurement $\mu$ is said to be a one-way right LOCC measurement.
2. There exists an alphabet $\Gamma$ and a measurement $\nu: \Gamma \rightarrow \operatorname{Pos}(\mathcal{Y})$, along with a measurement $\pi_{b}: \Sigma \rightarrow \operatorname{Pos}(\mathcal{X})$ for each $b \in \Gamma$, such that the equation

$$
\begin{equation*}
\mu(a)=\sum_{b \in \Gamma} \pi_{b}(a) \otimes \nu(b) \tag{6.145}
\end{equation*}
$$

holds for every $a \in \Sigma$. In this case the measurement $\mu$ is said to be a one-way left LOCC measurement.

Limitations on state discrimination by separable measurements
One may consider the problem of state discrimination, as was discussed in Chapter 3, in which measurements are restricted to be separable or LOCC measurements. Many examples of sets of orthogonal pure states are known that cannot be discriminated by separable or LOCC measurements without error. The following theorem provides one class of examples, and implies that there exist relatively small sets of orthogonal pure states having this characteristic.

Theorem 6.30 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, let $n=\operatorname{dim}(\mathcal{Y})$, and assume $n \leq \operatorname{dim}(\mathcal{X})$. Also let

$$
\begin{equation*}
\left\{U_{1}, \ldots, U_{m}\right\} \in \mathrm{U}(\mathcal{Y}, \mathcal{X}) \tag{6.146}
\end{equation*}
$$

be an orthogonal collection of isometries, and let $u_{k} \in \mathcal{X} \otimes \mathcal{Y}$ be the unit vector defined as

$$
\begin{equation*}
u_{k}=\frac{1}{\sqrt{n}} \operatorname{vec}\left(U_{k}\right) \tag{6.147}
\end{equation*}
$$

for each $k \in\{1, \ldots, m\}$. For every separable measurement of the form

$$
\begin{equation*}
\mu:\{1, \ldots, m\} \rightarrow \operatorname{Sep}(\mathcal{X}: \mathcal{Y}) \tag{6.148}
\end{equation*}
$$

it holds that

$$
\begin{equation*}
\sum_{k=1}^{m}\left\langle\mu(k), u_{k} u_{k}^{*}\right\rangle \leq \operatorname{dim}(\mathcal{X}) \tag{6.149}
\end{equation*}
$$

Proof Under the assumption that $\mu$ is a separable measurement, one may write

$$
\begin{equation*}
\mu(k)=\sum_{a \in \Sigma} P_{k, a} \otimes Q_{k, a} \tag{6.150}
\end{equation*}
$$

for each $k \in\{1, \ldots, m\}$, for some choice of an alphabet $\Sigma$ and collections of positive semidefinite operators as follows:

$$
\begin{align*}
& \left\{P_{k, a}: k \in\{1, \ldots, m\}, a \in \Sigma\right\} \subset \operatorname{Pos}(\mathcal{X})  \tag{6.151}\\
& \left\{Q_{k, a}: k \in\{1, \ldots, m\}, a \in \Sigma\right\} \subset \operatorname{Pos}(\mathcal{Y})
\end{align*}
$$

(No generality lost in using the same alphabet $\Sigma$ in the expressions (6.150) for each choice of $k$, as one is free to choose $\Sigma$ to be as large as is needed, and to set $P_{k, a}=0$ or $Q_{k, a}=0$ for some choices of $k$ and $a$ as necessary.) It holds that

$$
\begin{array}{r}
\left\langle\mu(k), \operatorname{vec}\left(U_{k}\right) \operatorname{vec}\left(U_{k}\right)^{*}\right\rangle=\sum_{a \in \Sigma} \operatorname{Tr}\left(U_{k}^{*} P_{k, a} U_{k} Q_{k, a}^{\top}\right) \\
\leq \sum_{a \in \Sigma}\left\|U_{k}^{*} P_{k, a} U_{k}\right\|_{2}\left\|Q_{k, a}\right\|_{2} \leq \sum_{a \in \Sigma}\left\|P_{k, a}\right\|_{1}\left\|Q_{k, a}\right\|_{1}  \tag{6.152}\\
=\sum_{a \in \Sigma} \operatorname{Tr}\left(P_{k, a}\right) \operatorname{Tr}\left(Q_{k, a}\right)=\operatorname{Tr}(\mu(k))
\end{array}
$$

and therefore

$$
\begin{equation*}
\sum_{k=1}^{m}\left\langle\mu(k), \operatorname{vec}\left(U_{k}\right) \operatorname{vec}\left(U_{k}\right)^{*}\right\rangle \leq \sum_{k=1}^{m} \operatorname{Tr}(\mu(k))=n \operatorname{dim}(\mathcal{X}) \tag{6.153}
\end{equation*}
$$

The theorem follows by dividing both sides of this inequality by $n$.

For any set of unit vectors $\left\{u_{1}, \ldots, u_{m}\right\}$ as described by this theorem, for which $m>\operatorname{dim}(\mathcal{X})$, one therefore has that

$$
\begin{equation*}
\frac{1}{m} \sum_{k=1}^{m}\left\langle\mu(k), u_{k} u_{k}^{*}\right\rangle \leq \frac{\operatorname{dim}(\mathcal{X})}{m}<1 \tag{6.154}
\end{equation*}
$$

Consequently, for one of the $m$ pure states associated with these vectors being selected uniformly at random, any separable measurement that aims to discriminate these states must err with probability strictly greater than 0 .

## LOCC discrimination of any pair of orthogonal pure states

Although Theorem 6.30 establishes that there exist relatively small sets of orthogonal pure states that cannot be perfectly discriminated by separable measurements, the same cannot be said about pairs of orthogonal pure states. Indeed, every pair of orthogonal pure states can be discriminated without error by a one-way LOCC measurement. The following lemma is used to prove this fact.

Lemma 6.31 Let $\mathcal{X}$ be a complex Euclidean space of dimension $n$ and let $X \in \mathrm{~L}(\mathcal{X})$ be an operator satisfying $\operatorname{Tr}(X)=0$. There exists an orthonormal basis $\left\{x_{1}, \ldots, x_{n}\right\}$ of $\mathcal{X}$ such that $x_{k}^{*} X x_{k}=0$ for all $k \in\{1, \ldots, n\}$.

Proof The proof is by induction on $n$. The base case $n=1$ is immediate, so it will be assumed that $n \geq 2$ for the rest of the proof. It will also be assumed that $\mathcal{X}=\mathbb{C}^{n}$, which causes no loss of generality.

For every integer $k \in\{1, \ldots, n\}$, it holds that $\lambda_{k}(X) \in \mathcal{N}(X)$, where $\mathcal{N}(X)$ denotes the numerical range of $X$. By the Toeplitz-Hausdorff theorem (Theorem 3.54), the numerical range is convex, and therefore

$$
\begin{equation*}
0=\frac{1}{n} \operatorname{Tr}(X)=\frac{1}{n} \sum_{k=1}^{n} \lambda_{k}(X) \in \mathcal{N}(X) \tag{6.155}
\end{equation*}
$$

By the definition of the numerical range, there must therefore exist a unit vector $x_{n} \in \mathcal{X}$ such that $x_{n}^{*} X x_{n}=0$.

Let $V \in \mathrm{U}\left(\mathbb{C}^{n-1}, \mathbb{C}^{n}\right)$ be any isometry that satisfies $x_{n} \perp \operatorname{im}(V)$, which is equivalent to

$$
\begin{equation*}
V V^{*}=\mathbb{1}-x_{n} x_{n}^{*} \tag{6.156}
\end{equation*}
$$

It holds that

$$
\begin{equation*}
\operatorname{Tr}\left(V^{*} X V\right)=\operatorname{Tr}\left(\left(\mathbb{1}-x_{n} x_{n}^{*}\right) X\right)=\operatorname{Tr}(X)-x_{n}^{*} X x_{n}=0 \tag{6.157}
\end{equation*}
$$

As $V^{*} X V \in \mathrm{~L}\left(\mathbb{C}^{n-1}\right)$, the hypothesis of induction implies that there exists
an orthonormal basis $\left\{u_{1}, \ldots, u_{n-1}\right\}$ of $\mathbb{C}^{n-1}$ such that

$$
\begin{equation*}
u_{k}^{*}\left(V^{*} X V\right) u_{k}=0 \tag{6.158}
\end{equation*}
$$

for all $k \in\{1, \ldots, n-1\}$. Define $x_{k}=V u_{k}$ for each $k \in\{1, \ldots, n-1\}$, and observe that $\left\{x_{1}, \ldots, x_{n-1}\right\}$ is an orthonormal set, with each element $x_{k}$ of this set satisfying $x_{k}^{*} X x_{k}=0$. As $V$ is an isometry and $x_{n} \perp \operatorname{im}(X)$, it follows that $\left\{x_{1}, \ldots, x_{n}\right\}$ is an orthonormal basis of $\mathcal{X}$ having the property stated by the lemma.

Theorem 6.32 Let $u_{0}, u_{1} \in \mathcal{X} \otimes \mathcal{Y}$ be orthogonal unit vectors, for $\mathcal{X}$ and $\mathcal{Y}$ being complex Euclidean spaces. There exists a one-way LOCC measurement

$$
\begin{equation*}
\mu:\{0,1\} \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.159}
\end{equation*}
$$

such that

$$
\begin{equation*}
\left\langle\mu(0), u_{0} u_{0}^{*}\right\rangle=1=\left\langle\mu(1), u_{1} u_{1}^{*}\right\rangle . \tag{6.160}
\end{equation*}
$$

Proof Let $n=\operatorname{dim}(\mathcal{Y})$ and let $A_{0}, A_{1} \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$ be the unique operators satisfying $u_{0}=\operatorname{vec}\left(A_{0}\right)$ and $u_{1}=\operatorname{vec}\left(A_{1}\right)$. The orthogonality of the vectors $u_{0}$ and $u_{1}$ is equivalent to the condition $\operatorname{Tr}\left(A_{0}^{*} A_{1}\right)=0$. By Lemma 6.31, there exists an orthonormal basis $\left\{x_{1}, \ldots, x_{n}\right\}$ of $\mathcal{Y}$ with the property that $x_{k}^{*} A_{0}^{*} A_{1} x_{k}=0$, which is equivalent to the condition that

$$
\begin{equation*}
\left\langle A_{0} x_{k} x_{k}^{*} A_{0}^{*}, A_{1} x_{k} x_{k}^{*} A_{1}^{*}\right\rangle=0, \tag{6.161}
\end{equation*}
$$

for every $k \in\{1, \ldots, n\}$.
Define a measurement $\nu:\{1, \ldots, n\} \rightarrow \operatorname{Pos}(\mathcal{Y})$ as

$$
\begin{equation*}
\nu(k)=\overline{x_{k}} x_{k}^{\top} \tag{6.162}
\end{equation*}
$$

for each $k \in\{1, \ldots, n\}$. By the equation (6.161), one has that there must exist a measurement $\pi_{k}:\{0,1\} \rightarrow \operatorname{Pos}(\mathcal{X})$, for each $k \in\{1, \ldots, n\}$, such that

$$
\begin{equation*}
\left\langle\pi_{k}(0), A_{1} x_{k} x_{k}^{*} A_{1}^{*}\right\rangle=0=\left\langle\pi_{k}(1), A_{0} x_{k} x_{k}^{*} A_{0}^{*}\right\rangle . \tag{6.163}
\end{equation*}
$$

Finally, define $\mu:\{0,1\} \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ as

$$
\begin{equation*}
\mu(a)=\sum_{k=1}^{n} \pi_{k}(a) \otimes \nu(k) \tag{6.164}
\end{equation*}
$$

for each $a \in\{0,1\}$, which is a one-way measurement with respect to the
second condition of Definition 6.29. It holds that

$$
\begin{align*}
\left\langle\mu(0), u_{1} u_{1}^{*}\right\rangle & =\sum_{k=1}^{n}\left\langle\pi_{k}(0),\left(\mathbb{1} \otimes x_{k}^{\top}\right) \operatorname{vec}\left(A_{1}\right) \operatorname{vec}\left(A_{1}\right)^{*}\left(\mathbb{1} \otimes \overline{x_{k}}\right)\right\rangle  \tag{6.165}\\
& =\sum_{k=1}^{n}\left\langle\pi_{k}(0), A_{1} x_{k} x_{k}^{*} A_{1}^{*}\right\rangle=0
\end{align*}
$$

and through a similar calculation one finds that $\left\langle\mu(1), u_{0} u_{0}^{*}\right\rangle=0$, which completes the proof.

Remark The preceding proof may be adapted in a straightforward way to prove that there exists a one-way LOCC measurement respecting the first condition of Definition 6.29, as opposed to the second, that satisfies the requirements of the theorem.

### 6.2 Manipulation of entanglement

As presented in the previous section, entanglement is defined as a lack of separability: for two complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, a bipartite state $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ that is not contained in the set $\operatorname{Sep}(\mathcal{X}: \mathcal{Y})$ is entangled with respect to the bipartition between $\mathcal{X}$ and $\mathcal{Y}$. This definition is qualitative, in the sense that it does not provide a measure of how much entanglement is present in a given state or suggest how two entangled states might relate to one another. The present section discusses such notions, and develops basic concepts and techniques relating to quantitative aspects of entanglement.

### 6.2.1 Entanglement transformation

The next theorem establishes a necessary and sufficient condition under which two individuals may transform one pure state into another by means of local operations and classical communication. The condition concerns the reductions of the initial and final pure states to one of the two individuals, requiring that the reduction of the initial state is majorized by the reduction of the final state. This condition is not only equivalent to the existence of an LOCC (or even a separable) channel transforming the initial state to the final state, but also implies that the transformation can be accomplished with one-way classical communication, from either of the two individuals to the other. The theorem offers a tool through which two fundamental ways of quantifying how much entanglement exists in a given state, called the entanglement cost and the distillable entanglement, may be analyzed for pure states.

Theorem 6.33 (Nielsen's theorem) Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces and let $u, v \in \mathcal{X} \otimes \mathcal{Y}$ be unit vectors. The following statements are equivalent:

1. $\operatorname{Tr}_{\mathcal{Y}}\left(u u^{*}\right) \prec \operatorname{Tr}_{\mathcal{Y}}\left(v v^{*}\right)$.
2. There exists an alphabet $\Sigma$ and collections of operators

$$
\begin{equation*}
\left\{U_{a}: a \in \Sigma\right\} \subset \mathrm{U}(\mathcal{X}) \quad \text { and } \quad\left\{B_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}) \tag{6.166}
\end{equation*}
$$

satisfying

$$
\begin{equation*}
\sum_{a \in \Sigma} B_{a}^{*} B_{a}=\mathbb{1}_{\mathcal{Y}} \tag{6.167}
\end{equation*}
$$

and

$$
\begin{equation*}
v v^{*}=\sum_{a \in \Sigma}\left(U_{a} \otimes B_{a}\right) u u^{*}\left(U_{a} \otimes B_{a}\right)^{*} \tag{6.168}
\end{equation*}
$$

3. There exists an alphabet $\Sigma$ and collections of operators

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{X}) \quad \text { and } \quad\left\{V_{a}: a \in \Sigma\right\} \subset \mathrm{U}(\mathcal{Y}) \tag{6.169}
\end{equation*}
$$

satisfying

$$
\begin{equation*}
\sum_{a \in \Sigma} A_{a}^{*} A_{a}=\mathbb{1}_{\mathcal{X}} \tag{6.170}
\end{equation*}
$$

and

$$
\begin{equation*}
v v^{*}=\sum_{a \in \Sigma}\left(A_{a} \otimes V_{a}\right) u u^{*}\left(A_{a} \otimes V_{a}\right)^{*} \tag{6.171}
\end{equation*}
$$

4. There exists a separable channel ${ }^{3} \Phi \in \operatorname{SepC}(\mathcal{X}: \mathcal{Y})$ such that

$$
\begin{equation*}
v v^{*}=\Phi\left(u u^{*}\right) \tag{6.172}
\end{equation*}
$$

Proof Let $X, Y \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$ be the unique operators for which $u=\operatorname{vec}(X)$ and $v=\operatorname{vec}(Y)$, and let

$$
\begin{equation*}
X=\sum_{k=1}^{r} s_{k} x_{k} y_{k}^{*} \tag{6.173}
\end{equation*}
$$

be a singular value decomposition of $X$, for $r=\operatorname{rank}(X)$.
Assume first that statement 1 holds, which is equivalent to $X X^{*} \prec Y Y^{*}$. There must therefore exist an alphabet $\Sigma$, a probability vector $p \in \mathcal{P}(\Sigma)$, and a collection of unitary operators $\left\{W_{a}: a \in \Sigma\right\} \subset \mathrm{U}(\mathcal{X})$ such that

$$
\begin{equation*}
X X^{*}=\sum_{a \in \Sigma} p(a) W_{a} Y Y^{*} W_{a}^{*} \tag{6.174}
\end{equation*}
$$

[^0]Let $\mathcal{Z}=\mathbb{C}^{\Sigma}$ and define an operator $Z \in \mathrm{~L}(\mathcal{Y} \otimes \mathcal{Z}, \mathcal{X})$ as

$$
\begin{equation*}
Z=\sum_{a \in \Sigma} \sqrt{p(a)} W_{a} Y \otimes e_{a}^{*} . \tag{6.175}
\end{equation*}
$$

It holds that

$$
\begin{equation*}
Z Z^{*}=\sum_{a \in \Sigma} p(a) W_{a} Y Y^{*} W_{a}^{*}=X X^{*}, \tag{6.176}
\end{equation*}
$$

and therefore $Z$ and $X$ agree on their singular values, and on the possible choices for their left singular vectors. It follows that one may write

$$
\begin{equation*}
Z=\sum_{k=1}^{r} s_{k} x_{k} w_{k}^{*} \tag{6.177}
\end{equation*}
$$

for $\left\{w_{1}, \ldots, w_{r}\right\} \subset \mathcal{Y} \otimes \mathcal{Z}$ being an orthonormal collection of vectors. Let $V \in \mathrm{U}(\mathcal{Y}, \mathcal{Y} \otimes \mathcal{Z})$ be an isometry for which $V y_{k}=w_{k}$ for all $k \in\{1, \ldots, r\}$, so that $X V^{*}=Z$.

Now, define operators

$$
\begin{equation*}
U_{a}=W_{a}^{*} \quad \text { and } \quad B_{a}=\left(\mathbb{1}_{\mathcal{Y}} \otimes e_{a}^{*}\right) \bar{V} \tag{6.178}
\end{equation*}
$$

for each $a \in \Sigma$. As $V$ is an isometry, so too is $\bar{V}$, and therefore

$$
\begin{equation*}
\left.\sum_{a \in \Sigma} B_{a}^{*} B_{a}=\sum_{a \in \Sigma} V^{\top} \mathbb{1}_{\mathcal{Y}} \otimes E_{a, a}\right) \bar{V}=V^{\top} \bar{V}=\mathbb{1}_{\mathcal{Y}} . \tag{6.179}
\end{equation*}
$$

It holds that

$$
\begin{equation*}
W_{a}^{*} X B_{a}^{\top}=W_{a}^{*} X V^{*}\left(\mathbb{1}_{\mathcal{Y}} \otimes e_{a}\right)=W_{a}^{*} Z\left(\mathbb{1}_{\mathcal{Y}} \otimes e_{a}\right)=\sqrt{p(a)} Y \tag{6.180}
\end{equation*}
$$

for each $a \in \Sigma$, and therefore

$$
\begin{align*}
\sum_{a \in \Sigma}\left(U_{a} \otimes\right. & \left.B_{a}\right) u u^{*}\left(U_{a} \otimes B_{a}\right)^{*} \\
& =\sum_{a \in \Sigma} \operatorname{vec}\left(W_{a}^{*} X B_{a}^{\top}\right) \operatorname{vec}\left(W_{a}^{*} X B_{a}^{\top}\right)^{*}  \tag{6.181}\\
& =\sum_{a \in \Sigma} p(a) \operatorname{vec}(Y) \operatorname{vec}(Y)^{*} \\
& =v v^{*} .
\end{align*}
$$

It has been established that statement 1 implies statement 2 .
The fact that statement 1 implies statement 3 is established by a similar argument with the roles of $\mathcal{X}$ and $\mathcal{Y}$ exchanged, along with the observation that $\operatorname{Tr} \mathcal{Y}_{\mathcal{Y}}\left(u u^{*}\right) \prec \operatorname{Tr} \mathcal{Y}\left(v v^{*}\right)$ is equivalent to $\operatorname{Tr} \mathcal{X}\left(u u^{*}\right) \prec \operatorname{Tr}_{\mathcal{X}}\left(v v^{*}\right)$.

Statements 2 and 3 each imply statement 4 directly, as the mappings defined by the actions

$$
\begin{align*}
u u^{*} \mapsto \sum_{a \in \Sigma}\left(U_{a} \otimes B_{a}\right) u u^{*}\left(U_{a} \otimes B_{a}\right)^{*},  \tag{6.182}\\
u u^{*} \mapsto \sum_{a \in \Sigma}\left(A_{a} \otimes V_{a}\right) u u^{*}\left(A_{a} \otimes V_{a}\right)^{*}
\end{align*}
$$

are both separable channels.
Finally, assume statement 4 holds, letting $\Phi \in \operatorname{SepC}(\mathcal{X}: \mathcal{Y})$ be a fixed separable channel for which $\Phi\left(u u^{*}\right)=v v^{*}$. It will be proved that

$$
\begin{equation*}
\lambda\left(X X^{*}\right) \prec \lambda\left(Y Y^{*}\right) ; \tag{6.183}
\end{equation*}
$$

by Theorem 4.32, this relation is equivalent to $X X^{*} \prec Y Y^{*}$, which in turn is equivalent to statement 1 . Let $n=\operatorname{dim}(\mathcal{X})$, and observe that

$$
\begin{equation*}
\sum_{k=1}^{n} \lambda_{k}\left(X X^{*}\right)=\operatorname{Tr}\left(X X^{*}\right)=1=\operatorname{Tr}\left(Y Y^{*}\right)=\sum_{k=1}^{n} \lambda_{k}\left(Y Y^{*}\right) \tag{6.184}
\end{equation*}
$$

by the assumption that $u$ and $v$ are unit vectors. By Theorem 4.30, one finds that the relation (6.183) will therefore follow from the inequality

$$
\begin{equation*}
\sum_{k=m}^{n} \lambda_{k}\left(Y Y^{*}\right) \leq \sum_{k=m}^{n} \lambda_{k}\left(X X^{*}\right) \tag{6.185}
\end{equation*}
$$

holding for every choice of $m \in\{1, \ldots, n\}$.
By the separability of the channel $\Phi$, there must exist an alphabet $\Sigma$ and two collections of operators

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{X}) \quad \text { and } \quad\left\{B_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}) \tag{6.186}
\end{equation*}
$$

with $\left\{A_{a} \otimes B_{a}: a \in \Sigma\right\}$ being a set of Kraus operators of $\Phi$, for which

$$
\begin{equation*}
v v^{*}=\sum_{a \in \Sigma}\left(A_{a} \otimes B_{a}\right) u u^{*}\left(A_{a} \otimes B_{a}\right)^{*} . \tag{6.187}
\end{equation*}
$$

As $v v^{*}$ is a rank-one operator, it follows that there must exist a probability vector $p \in \mathcal{P}(\Sigma)$ such that

$$
\begin{equation*}
\left(A_{a} \otimes B_{a}\right) u u^{*}\left(A_{a} \otimes B_{a}\right)^{*}=p(a) v v^{*} \tag{6.188}
\end{equation*}
$$

which is equivalent to

$$
\begin{equation*}
\operatorname{vec}\left(A_{a} X B_{a}^{\top}\right) \operatorname{vec}\left(A_{a} X B_{a}^{\top}\right)^{*}=p(a) \operatorname{vec}(Y) \operatorname{vec}(Y)^{*} \tag{6.189}
\end{equation*}
$$

for each $a \in \Sigma$. By taking the partial trace over $\mathcal{Y}$, it follows that

$$
\begin{equation*}
A_{a} X B_{a}^{\top} \overline{B_{a}} X^{*} A_{a}^{*}=p(a) Y Y^{*} \tag{6.190}
\end{equation*}
$$

for each $a \in \Sigma$, and therefore

$$
\begin{equation*}
\sum_{k=m}^{n} \lambda_{k}\left(Y Y^{*}\right)=\sum_{k=m}^{n} \sum_{a \in \Sigma} \lambda_{k}\left(A_{a} X B_{a}^{\top} \overline{B_{a}} X^{*} A_{a}^{*}\right) \tag{6.191}
\end{equation*}
$$

for each $m \in\{1, \ldots, n\}$.
Next, for each choice of $a \in \Sigma$ and $m \in\{1, \ldots, n\}$, let $\Pi_{a, m} \in \operatorname{Proj}(\mathcal{X})$ be the projection operator onto the orthogonal complement of the subspace of $\mathcal{X}$ spanned by the set $\left\{A_{a} x_{1}, \ldots, A_{a} x_{m-1}\right\}$, where one is to assume $x_{k}=0$ for $k>r$. By the definition of these projection operators, it is evident that

$$
\begin{equation*}
\left\langle\Pi_{a, m}, A_{a} X B_{a}^{\top} \overline{B_{a}} X^{*} A_{a}^{*}\right\rangle=\left\langle\Pi_{a, m}, A_{a} X_{m} B_{a}^{\top} \overline{B_{a}} X_{m}^{*} A_{a}^{*}\right\rangle \tag{6.192}
\end{equation*}
$$

for every $a \in \Sigma$ and $m \in\{1, \ldots, n\}$, where

$$
\begin{equation*}
X_{m}=\sum_{k=m}^{r} s_{k} x_{k} y_{k}^{*} \tag{6.193}
\end{equation*}
$$

and one is to interpret that $X_{m}=0$ for $m>r$. Because each operator $\Pi_{a, m}$ is a projection, and the operator $A_{a} X_{m} B_{a}^{\top} \overline{B_{a}} X_{m}^{*} A_{a}^{*}$ is positive semidefinite, it follows that

$$
\begin{equation*}
\left\langle\Pi_{a, m}, A_{a} X_{m} B_{a}^{\top} \overline{B_{a}} X_{m}^{*} A_{a}^{*}\right\rangle \leq \operatorname{Tr}\left(A_{a} X_{m} B_{a}^{\top} \overline{B_{a}} X_{m}^{*} A_{a}^{*}\right) \tag{6.194}
\end{equation*}
$$

Using the fact that $\Phi$ is a channel, and therefore preserves trace, one finds that

$$
\begin{align*}
& \sum_{a \in \Sigma} \operatorname{Tr}\left(A_{a} X_{m} B_{a}^{\top} \overline{B_{a}} X_{m}^{*} A_{a}^{*}\right)=\operatorname{Tr}\left(\Phi\left(\operatorname{vec}\left(X_{m}\right) \operatorname{vec}\left(X_{m}\right)^{*}\right)\right) \\
& \quad=\operatorname{Tr}\left(\operatorname{vec}\left(X_{m}\right) \operatorname{vec}\left(X_{m}\right)^{*}\right)=\operatorname{Tr}\left(X_{m} X_{m}^{*}\right)=\sum_{k=m}^{n} \lambda_{k}\left(X X^{*}\right) \tag{6.195}
\end{align*}
$$

for each $m \in\{1, \ldots, n\}$.
Finally, as it necessarily holds that $\operatorname{rank}\left(\Pi_{a, m}\right) \geq n-m+1$ for every $a \in \Sigma$ and $m \in\{1, \ldots, n\}$, it follows that

$$
\begin{equation*}
\left\langle\Pi_{a, m}, A_{a} X B_{a}^{\top} \overline{B_{a}} X^{*} A_{a}^{*}\right\rangle \geq \sum_{k=m}^{n} \lambda_{k}\left(A_{a} X B_{a}^{\top} \overline{B_{a}} X^{*} A_{a}^{*}\right) \tag{6.196}
\end{equation*}
$$

By combining $(6.191),(6.192),(6.194),(6.195)$, and (6.196), one finds that

$$
\begin{equation*}
\sum_{k=m}^{n} \lambda_{k}\left(Y Y^{*}\right) \leq \sum_{k=m}^{n} \lambda_{k}\left(X X^{*}\right) \tag{6.197}
\end{equation*}
$$

which establishes (6.183), and therefore completes the proof.

Theorem 6.33 implies the following corollary, which characterizes the pure state transformations, possibly involving different complex Euclidean spaces, that may be realized by LOCC channels.

Corollary 6.34 Let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$ be complex Euclidean spaces and let $x \in \mathcal{X} \otimes \mathcal{Y}$ and $y \in \mathcal{Z} \otimes \mathcal{W}$ be unit vectors. The following statements are equivalent:

1. For $\rho=\operatorname{Tr}_{\mathcal{Y}}\left(x x^{*}\right), \sigma=\operatorname{Tr}_{\mathcal{W}}\left(y y^{*}\right)$, and $r=\min \{\operatorname{rank}(\rho), \operatorname{rank}(\sigma)\}$, it holds that

$$
\begin{equation*}
\lambda_{1}(\rho)+\cdots+\lambda_{m}(\rho) \leq \lambda_{1}(\sigma)+\cdots+\lambda_{m}(\sigma) \tag{6.198}
\end{equation*}
$$

for every $m \in\{1, \ldots, r\}$.
2. There exists a one-way right $L O C C$ channel $\Phi \in \operatorname{LOCC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ for which it holds that $\Phi\left(x x^{*}\right)=y y^{*}$.
3. There exists a one-way left LOCC channel $\Phi \in \operatorname{LOCC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ for which it holds that $\Phi\left(x x^{*}\right)=y y^{*}$.
4. There exists a separable channel $\Phi \in \operatorname{SepC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ for which it holds that $\Phi\left(x x^{*}\right)=y y^{*}$.

Proof Define four isometries, $A_{0} \in \mathrm{U}(\mathcal{X}, \mathcal{X} \oplus \mathcal{Z}), B_{0} \in \mathrm{U}(\mathcal{Y}, \mathcal{Y} \oplus \mathcal{W})$, $A_{1} \in \mathrm{U}(\mathcal{Z}, \mathcal{X} \oplus \mathcal{Z})$, and $B_{1} \in \mathrm{U}(\mathcal{W}, \mathcal{Y} \oplus \mathcal{W})$, as follows:

$$
\begin{align*}
& A_{0} x=x \oplus 0, \quad A_{1} z=0 \oplus z \\
& B_{0} y=y \oplus 0, \quad B_{1} w=0 \oplus w \tag{6.199}
\end{align*}
$$

for every choice of vectors $x \in \mathcal{X}, y \in \mathcal{Y}, z \in \mathcal{Z}$, and $w \in \mathcal{W}$. Also define four channels, $\Psi_{0} \in \mathrm{C}(\mathcal{X} \oplus \mathcal{Z}, \mathcal{X}), \Lambda_{0} \in \mathrm{C}(\mathcal{Y} \oplus \mathcal{W}, \mathcal{Y}), \Psi_{1} \in \mathrm{C}(\mathcal{X} \oplus \mathcal{Z}, \mathcal{Z})$, and $\Lambda_{1} \in \mathrm{C}(\mathcal{Y} \oplus \mathcal{W}, \mathcal{W})$, as

$$
\begin{align*}
\Psi_{0}(X) & =A_{0}^{*} X A_{0}+\left\langle\mathbb{1}_{\mathcal{X} \oplus \mathcal{Z}}-A_{0} A_{0}^{*}, X\right\rangle \tau_{0} \\
\Lambda_{0}(Y) & =B_{0}^{*} Y B_{0}+\left\langle\mathbb{1}_{\mathcal{Y} \oplus \mathcal{W}}-B_{0} B_{0}^{*}, Y\right\rangle \xi_{0} \\
\Psi_{1}(X) & =A_{1}^{*} X A_{1}+\left\langle\mathbb{1}_{\mathcal{X} \oplus \mathcal{Z}}-A_{1} A_{1}^{*}, X\right\rangle \tau_{1}  \tag{6.200}\\
\Lambda_{1}(Y) & =B_{1}^{*} Y B_{1}+\left\langle\mathbb{1}_{\mathcal{Y} \oplus \mathcal{W}}-B_{1} B_{1}^{*}, Y\right\rangle \xi_{1}
\end{align*}
$$

for all $X \in \mathrm{~L}(\mathcal{X} \oplus \mathcal{Z})$ and $Y \in \mathrm{~L}(\mathcal{Y} \oplus \mathcal{W})$, where $\tau_{0} \in \mathrm{D}(\mathcal{X}), \xi_{0} \in \mathrm{D}(\mathcal{Y})$, $\tau_{1} \in \mathrm{D}(\mathcal{Z})$, and $\xi_{1} \in \mathrm{D}(\mathcal{W})$ are fixed, but otherwise arbitrarily selected, density operators.

Assume first that statement 1 holds. One concludes that

$$
\begin{equation*}
A_{0} \rho A_{0}^{*} \prec A_{1} \sigma A_{1}^{*}, \tag{6.201}
\end{equation*}
$$

and therefore the four equivalent statements of Theorem 6.33 hold for the vectors

$$
\begin{equation*}
u=\left(A_{0} \otimes B_{0}\right) x \quad \text { and } \quad v=\left(A_{1} \otimes B_{1}\right) y \tag{6.202}
\end{equation*}
$$

There must therefore exist a one-way right LOCC channel $\Xi$, of the form specified in the statement of Theorem 6.33 , such that $\Xi\left(u u^{*}\right)=v v^{*}$. Define $\Phi \in \mathrm{C}(\mathcal{X} \otimes \mathcal{Y}, \mathcal{Z} \otimes \mathcal{W})$ as

$$
\begin{equation*}
\Phi(X)=\left(\left(\Psi_{1} \otimes \Lambda_{1}\right) \Xi\right)\left(\left(A_{0} \otimes B_{0}\right) X\left(A_{0} \otimes B_{0}\right)^{*}\right) \tag{6.203}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$. It holds that $\Phi$ is a one-way right LOCC channel satisfying $\Phi\left(x x^{*}\right)=y y^{*}$, and therefore statement 1 implies statement 2 . The fact that statement 1 implies statement 3 is similar.

Statements 2 and 3 trivially imply that statement 4 holds.
Finally, assume statement 4 holds. Define a channel $\Xi$ as

$$
\begin{equation*}
\Xi(X)=\left(A_{1} \otimes B_{1}\right)\left(\Phi\left(\Psi_{0} \otimes \Lambda_{0}\right)\right)(X)\left(A_{1} \otimes B_{1}\right)^{*} \tag{6.204}
\end{equation*}
$$

for all $X \in \mathrm{~L}((\mathcal{X} \oplus \mathcal{Z}) \otimes(\mathcal{Y} \oplus \mathcal{W}))$. The channel $\Xi$ is separable and satisfies

$$
\begin{equation*}
\Xi\left(u u^{*}\right)=v v^{*} \tag{6.205}
\end{equation*}
$$

for vectors $u$ and $v$ as in (6.202). The four equivalent statements listed in Theorem 6.33 therefore hold for $u$ and $v$, which implies

$$
\begin{align*}
& \operatorname{Tr} \mathcal{Y} \oplus \mathcal{W}( \\
&\left.\left(A_{0} \otimes B_{0}\right) x x^{*}\left(A_{0} \otimes B_{0}\right)^{*}\right)  \tag{6.206}\\
& \prec \operatorname{Tr}_{\mathcal{Y} \oplus \mathcal{W}}\left(\left(A_{1} \otimes B_{1}\right) y y^{*}\left(A_{1} \otimes B_{1}\right)^{*}\right) .
\end{align*}
$$

This relation is equivalent to (6.201), which implies that statement 1 holds, and completes the proof.

### 6.2.2 Distillable entanglement and entanglement cost

Let $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ be a state, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$. There are various ways in which one may quantify the amount of entanglement that is present in $\rho$, with respect to the bipartition between $\mathcal{X}$ and $\mathcal{Y}$. The distillable entanglement and entanglement cost represent two such measures. The distillable entanglement concerns the rate at which copies of the state $\rho$ can be converted into copies of the maximally entangled two-qubit state

$$
\begin{equation*}
\tau=\frac{1}{2} \sum_{a, b \in\{0,1\}} E_{a, b} \otimes E_{a, b} \tag{6.207}
\end{equation*}
$$

with high accuracy by means of an LOCC channel. The entanglement cost refers to the reverse process; it is the rate at which approximate copies of $\rho$
may be produced from copies of $\tau$ by an LOCC channel. In both cases, it is the asymptotic behavior of these processes, as the number of copies of each state grows, that is taken as the measure of entanglement.

For every bipartite state, the distillable entanglement is upper-bounded by the entanglement cost, with the two measures coinciding for pure states. In general, however, the two quantities may differ, with the entanglement cost being strictly larger than the distillable entanglement in some cases.

## Notation related to distillable entanglement and entanglement cost

The following notation will be useful when discussing both the distillable entanglement and entanglement cost of a bipartite state $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$.

First, for a given positive integer $n$, representing the number of copies of the state $\rho$ to be manipulated for either the distillable entanglement or entanglement cost, one may define an isometry

$$
\begin{equation*}
U_{n} \in \mathrm{U}\left((\mathcal{X} \otimes \mathcal{Y})^{\otimes n}, \mathcal{X}^{\otimes n} \otimes \mathcal{Y}^{\otimes n}\right) \tag{6.208}
\end{equation*}
$$

by the action

$$
\begin{equation*}
U_{n}\left(\operatorname{vec}\left(A_{1}\right) \otimes \cdots \otimes \operatorname{vec}\left(A_{n}\right)\right)=\operatorname{vec}\left(A_{1} \otimes \cdots \otimes A_{n}\right) \tag{6.209}
\end{equation*}
$$

for all operators $A_{1}, \ldots, A_{n} \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$. Equivalently, $U_{n}$ is defined by the action

$$
\begin{align*}
& U_{n}\left(\left(x_{1} \otimes y_{1}\right) \otimes \cdots \otimes\left(x_{n} \otimes y_{n}\right)\right)  \tag{6.210}\\
& \quad=\left(x_{1} \otimes \cdots \otimes x_{n}\right) \otimes\left(y_{1} \otimes \cdots \otimes y_{n}\right)
\end{align*}
$$

for all vectors $x_{1}, \ldots, x_{n} \in \mathcal{X}$ and $y_{1}, \ldots, y_{n} \in \mathcal{Y}$. This isometry has the effect of re-ordering the tensor factors of the space $(\mathcal{X} \otimes \mathcal{Y})^{\otimes n}$ so that it takes the form of a bipartite tensor product space $\mathcal{X}^{\otimes n} \otimes \mathcal{Y}^{\otimes n}$ that allows for notions concerning entanglement and separability to be conveniently stated.

Next, the binary alphabet will be denoted $\Gamma=\{0,1\}$, and the state

$$
\begin{equation*}
\tau=\frac{1}{2} \sum_{a, b \in\{0,1\}} E_{a, b} \otimes E_{a, b} \tag{6.211}
\end{equation*}
$$

is to be considered as an element of the set $\mathrm{D}(\mathcal{Z} \otimes \mathcal{W})$, for $\mathcal{Z}=\mathbb{C}^{\Gamma}$ and $\mathcal{W}=\mathbb{C}^{\Gamma}$. Similar to above, one may define an isometry

$$
\begin{equation*}
V_{m} \in \mathrm{U}\left((\mathcal{Z} \otimes \mathcal{W})^{\otimes m}, \mathcal{Z}^{\otimes m} \otimes \mathcal{W}^{\otimes m}\right) \tag{6.212}
\end{equation*}
$$

playing an analogous role to the isometry $U_{n}$, but for the spaces $\mathcal{Z}$ and $\mathcal{W}$ in place of $\mathcal{X}$ and $\mathcal{Y}$. This isometry is defined by the action

$$
\begin{equation*}
V_{m}\left(\operatorname{vec}\left(B_{1}\right) \otimes \cdots \otimes \operatorname{vec}\left(B_{m}\right)\right)=\operatorname{vec}\left(B_{1} \otimes \cdots \otimes B_{m}\right) \tag{6.213}
\end{equation*}
$$

for all operators $B_{1}, \ldots, B_{m} \in \mathrm{~L}(\mathcal{W}, \mathcal{Z})$. Equivalently, $V_{m}$ is defined by the action

$$
\begin{align*}
& V_{m}\left(\left(z_{1} \otimes w_{1}\right) \otimes \cdots \otimes\left(z_{m} \otimes w_{m}\right)\right)  \tag{6.214}\\
& \quad=\left(z_{1} \otimes \cdots \otimes z_{m}\right) \otimes\left(w_{1} \otimes \cdots \otimes w_{m}\right)
\end{align*}
$$

for all vectors $z_{1}, \ldots, z_{m} \in \mathcal{Z}$ and $w_{1}, \ldots, w_{m} \in \mathcal{W}$.

## Definitions of distillable entanglement and entanglement cost

With respect to the notation introduced above, the distillable entanglement and entanglement cost are defined as follows.

Definition 6.35 Let $X$ and $Y$ be registers and let $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ be a state of $(\mathrm{X}, \mathrm{Y})$. With respect to the state $\rho$, the distillable entanglement $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y})$ of the pair $(X, Y)$ is the supremum value of all nonnegative real numbers $\alpha \geq 0$ for which the following statement holds: there exists a sequence of LOCC channels $\left(\Psi_{1}, \Psi_{2}, \ldots\right)$, where

$$
\begin{equation*}
\Psi_{n} \in \operatorname{LOCC}\left(\mathcal{X}^{\otimes n}, \mathcal{Z}^{\otimes m}: \mathcal{Y}^{\otimes n}, \mathcal{W}^{\otimes m}\right) \tag{6.215}
\end{equation*}
$$

for $m=\lfloor\alpha n\rfloor$, such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \mathrm{~F}\left(V_{m} \tau^{\otimes m} V_{m}^{*}, \Psi_{n}\left(U_{n} \rho^{\otimes n} U_{n}^{*}\right)\right)=1 \tag{6.216}
\end{equation*}
$$

Definition 6.36 Let $X$ and $Y$ be registers and let $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ be a state of $(\mathrm{X}, \mathrm{Y})$. With respect to the state $\rho$, the entanglement cost $\mathrm{E}_{\mathrm{C}}(\mathrm{X}: \mathrm{Y})$ of the pair $(\mathbf{X}, \mathrm{Y})$ is the infimum value of all nonnegative real numbers $\alpha \geq 0$ for which the following statement holds: there exists a sequence of LOCC channels $\left(\Phi_{1}, \Phi_{2}, \ldots\right)$, where

$$
\begin{equation*}
\Phi_{n} \in \operatorname{LOCC}\left(\mathcal{Z}^{\otimes m}, \mathcal{X}^{\otimes n}: \mathcal{W}^{\otimes m}, \mathcal{Y}^{\otimes n}\right) \tag{6.217}
\end{equation*}
$$

for $m=\lfloor\alpha n\rfloor$, such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \mathrm{~F}\left(U_{n} \rho^{\otimes n} U_{n}^{*}, \Phi_{n}\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right)\right)=1 \tag{6.218}
\end{equation*}
$$

It is intuitive that the entanglement cost should be at least as large as the distillable entanglement, for any choice of $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$, for otherwise one could repeatedly distill copies of the state $\tau$ from copies of a given state $\rho$, use them to produce more copies of $\rho$, and repeat this process indefinitely, eventually producing any desired number of copies of $\tau$ from a finite number of copies of $\rho$. Such an "entanglement factory" must surely not be possible through local operations and classical communication alone. The following proposition confirms this intuition.

Proposition 6.37 Let X and Y be registers. With respect to every state of the pair $(\mathrm{X}, \mathrm{Y})$ it holds that $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y}) \leq \mathrm{E}_{\mathrm{C}}(\mathrm{X}: \mathrm{Y})$.

Proof Suppose that $n, m$, and $k$ are nonnegative integers, and

$$
\begin{gather*}
\Phi_{n} \in \operatorname{LOCC}\left(\mathcal{Z}^{\otimes m}, \mathcal{X}^{\otimes n}: \mathcal{W}^{\otimes m}, \mathcal{Y}^{\otimes n}\right) \\
\Psi_{n} \in \operatorname{LOCC}\left(\mathcal{X}^{\otimes n}, \mathcal{Z}^{\otimes k}: \mathcal{Y}^{\otimes n}, \mathcal{W}^{\otimes k}\right) \tag{6.219}
\end{gather*}
$$

are LOCC channels. The composition $\Psi_{n} \Phi_{n}$ is an LOCC, and therefore separable, channel. It holds that

$$
\begin{equation*}
V_{m} \tau^{\otimes m} V_{m}^{*} \in \operatorname{Ent}_{2^{m}}\left(\mathcal{Z}^{\otimes m}: \mathcal{W}^{\otimes m}\right) \tag{6.220}
\end{equation*}
$$

and Theorem 6.23 implies that

$$
\begin{equation*}
\left(\Psi_{n} \Phi_{n}\right)\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right) \in \operatorname{Ent}_{2^{m}}\left(\mathcal{Z}^{\otimes k}: \mathcal{W}^{\otimes k}\right) \tag{6.221}
\end{equation*}
$$

By Proposition 6.15, one finds that

$$
\begin{align*}
& \mathrm{F}\left(\left(\Psi_{n} \Phi_{n}\right)\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right), V_{k} \tau^{\otimes k} V_{k}^{*}\right)^{2}  \tag{6.222}\\
& \quad=\left\langle\left(\Psi_{n} \Phi_{n}\right)\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right), V_{k} \tau^{\otimes k} V_{k}^{*}\right\rangle \leq 2^{m-k}
\end{align*}
$$

Now, let $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ be any state of the pair $(\mathrm{X}, \mathrm{Y})$, and suppose $\alpha$ and $\beta$ are nonnegative real numbers satisfying the requirements of the definitions of entanglement cost and distillable entanglement, respectively, for the state $\rho$. For all $\varepsilon>0$, there must therefore exist a sufficiently large positive integer $n$ such that, for $m=\lfloor\alpha n\rfloor$ and $k=\lfloor\beta n\rfloor$, there exist LOCC channels of the form (6.219) for which the following bounds hold:

$$
\begin{align*}
& \mathrm{F}\left(\Phi_{n}\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right), U_{n} \rho^{\otimes n} U_{n}^{*}\right)>1-\varepsilon  \tag{6.223}\\
& \mathrm{F}\left(\Psi_{n}\left(U_{n} \rho^{\otimes n} V_{n}^{*}\right), V_{k} \tau^{\otimes k} V_{k}^{*}\right)>1-\varepsilon
\end{align*}
$$

Therefore, by Theorem 3.29, together with the monotonicity of the fidelity function under the action of channels (Theorem 3.27), one may conclude that

$$
\begin{equation*}
\mathrm{F}\left(\left(\Psi_{n} \Phi_{n}\right)\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right), V_{k} \tau^{\otimes k} V_{k}^{*}\right)>1-4 \varepsilon . \tag{6.224}
\end{equation*}
$$

Taking $\varepsilon<1 / 16$, one concludes that

$$
\begin{equation*}
\mathrm{F}\left(\left(\Psi_{n} \Phi_{n}\right)\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right), V_{k} \tau^{\otimes k} V_{k}^{*}\right)^{2}>\frac{1}{2} \tag{6.225}
\end{equation*}
$$

and therefore $m \geq k$ by (6.222). As this is so for all sufficiently large $n$, it follows that $\beta \leq \alpha$. One concludes that $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y}) \leq \mathrm{E}_{\mathrm{C}}(\mathrm{X}: \mathrm{Y})$.

## Pure state entanglement

The next theorem demonstrates that the entanglement cost and distillable entanglement are equal for bipartite pure states; in both cases, the value of these measures agrees with the von Neumann entropy of the states obtained by restricting the given pure state to either part of its bipartition.

Theorem 6.38 Let X and Y be registers. With respect to every pure state of the pair $(\mathrm{X}, \mathrm{Y})$, one has

$$
\begin{equation*}
E_{D}(X: Y)=H(X)=H(Y)=E_{C}(X: Y) \tag{6.226}
\end{equation*}
$$

Proof Let $u \in \mathcal{X} \otimes \mathcal{Y}$ be a unit vector, and consider the pure state $u u^{*}$ of the pair $(\mathrm{X}, \mathrm{Y})$. The equality $\mathrm{H}(\mathrm{X})=\mathrm{H}(\mathrm{Y})$ was discussed in Section 5.1.2. Specifically, by means of the Schmidt decomposition, one may write

$$
\begin{equation*}
u=\sum_{a \in \Sigma} \sqrt{p(a)} x_{a} \otimes y_{a} \tag{6.227}
\end{equation*}
$$

for some choice of an alphabet $\Sigma$, a probability vector $p \in \mathcal{P}(\Sigma)$, and two orthonormal collections $\left\{x_{a}: a \in \Sigma\right\} \subset \mathcal{X}$ and $\left\{y_{a}: a \in \Sigma\right\} \subset \mathcal{Y}$. It holds that

$$
\begin{equation*}
\operatorname{Tr}_{\mathcal{Y}}\left(u u^{*}\right)=\sum_{a \in \Sigma} p(a) x_{a} x_{a}^{*} \quad \text { and } \quad \operatorname{Tr}_{\mathcal{X}}\left(u u^{*}\right)=\sum_{a \in \Sigma} p(a) y_{a} y_{a}^{*} \tag{6.228}
\end{equation*}
$$

which implies that $\mathrm{H}(\mathrm{X})=\mathrm{H}(p)=\mathrm{H}(\mathrm{Y})$.
Next, recall that, for every positive integer $n$ and positive real number $\varepsilon>0$, the set of $\varepsilon$-typical strings $T_{n, \varepsilon}$ with respect to $p$ contains those strings $a_{1} \cdots a_{n} \in \Sigma^{n}$ for which

$$
\begin{equation*}
2^{-n(\mathrm{H}(p)+\varepsilon)}<p\left(a_{1}\right) \cdots p\left(a_{n}\right)<2^{-n(\mathrm{H}(p)-\varepsilon)} \tag{6.229}
\end{equation*}
$$

With this set in mind, one may define a vector $v_{n, \varepsilon} \in \mathcal{X}^{\otimes n} \otimes \mathcal{Y}^{\otimes n}$, for every positive integer $n$ and positive real number $\varepsilon>0$, as

$$
\begin{equation*}
v_{n, \varepsilon}=\sum_{a_{1} \cdots a_{n} \in T_{n, \varepsilon}} \sqrt{p\left(a_{1}\right) \cdots p\left(a_{n}\right)} x_{a_{1} \cdots a_{n}} \otimes y_{a_{1} \cdots a_{n}} \tag{6.230}
\end{equation*}
$$

where the shorthand notations

$$
\begin{equation*}
x_{a_{1} \cdots a_{n}}=x_{a_{1}} \otimes \cdots \otimes x_{a_{n}} \quad \text { and } \quad y_{a_{1} \cdots a_{n}}=y_{a_{1}} \otimes \cdots \otimes y_{a_{n}} \tag{6.231}
\end{equation*}
$$

have been used for the sake of brevity. Also define a normalized version of the vector $v_{n, \varepsilon}$ as

$$
\begin{equation*}
w_{n, \varepsilon}=\frac{v_{n, \varepsilon}}{\left\|v_{n, \varepsilon}\right\|} \tag{6.232}
\end{equation*}
$$

Observe that

$$
\begin{equation*}
2^{-n(\mathrm{H}(p)+\varepsilon)}<\lambda_{k}\left(\operatorname{Tr}_{\mathcal{Y} \otimes n}\left(v_{n, \varepsilon} v_{n, \varepsilon}^{*}\right)\right)<2^{-n(\mathrm{H}(p)-\varepsilon)} \tag{6.233}
\end{equation*}
$$

and therefore

$$
\begin{equation*}
\frac{2^{-n(\mathrm{H}(p)+\varepsilon)}}{\left\|v_{n, \varepsilon}\right\|^{2}}<\lambda_{k}\left(\operatorname{Tr}_{\mathcal{Y} \otimes n}\left(w_{n, \varepsilon} w_{n, \varepsilon}^{*}\right)\right)<\frac{2^{-n(\mathrm{H}(p)-\varepsilon)}}{\left\|v_{n, \varepsilon}\right\|^{2}} \tag{6.234}
\end{equation*}
$$

for $k=1, \ldots,\left|T_{n, \varepsilon}\right|$, while the remaining eigenvalues are zero in both cases.
Now, consider the entanglement cost of the pair $(\mathrm{X}, \mathrm{Y})$ with respect to the state $u u^{*}$. Let $\alpha$ be any real number such that $\alpha>\mathrm{H}(p)$, let $\varepsilon>0$ be sufficiently small so that $\alpha>\mathrm{H}(p)+2 \varepsilon$, and consider any choice of $n>1 / \varepsilon$. For $m=\lfloor\alpha n\rfloor$, it holds that $m \geq n(\mathrm{H}(p)+\varepsilon)$, and moreover

$$
\begin{equation*}
\lambda_{k}\left(\operatorname{Tr}_{\mathcal{W} \otimes m}\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right)\right)=2^{-m} \tag{6.235}
\end{equation*}
$$

for $k=1, \ldots, 2^{m}$. As

$$
\begin{equation*}
2^{-m} \leq 2^{-n(\mathrm{H}(p)+\varepsilon)} \leq \frac{2^{-n(\mathrm{H}(p)+\varepsilon)}}{\left\|v_{n, \varepsilon}\right\|^{2}} \tag{6.236}
\end{equation*}
$$

it follows that

$$
\begin{equation*}
\sum_{j=1}^{k} \lambda_{j}\left(\operatorname{Tr}_{\mathcal{W} \otimes m}\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right)\right) \leq \sum_{j=1}^{k} \lambda_{j}\left(\operatorname{Tr}_{\mathcal{Y} \otimes n}\left(w_{n, \varepsilon} w_{n, \varepsilon}^{*}\right)\right) \tag{6.237}
\end{equation*}
$$

for every $k \in\left\{1, \ldots, 2^{m}\right\}$. It follows by Corollary 6.34 to Nielsen's theorem (Theorem 6.33) that there exists an LOCC channel

$$
\begin{equation*}
\Phi_{n} \in \operatorname{LOCC}\left(\mathcal{Z}^{\otimes m}, \mathcal{X}^{\otimes n}: \mathcal{W}^{\otimes m}, \mathcal{Y}^{\otimes n}\right) \tag{6.238}
\end{equation*}
$$

such that

$$
\begin{equation*}
\Phi_{n}\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right)=w_{n, \varepsilon} w_{n, \varepsilon}^{*} \tag{6.239}
\end{equation*}
$$

As

$$
\begin{equation*}
\mathrm{F}\left(U_{n}\left(u u^{*}\right)^{\otimes n} U_{n}^{*}, w_{n, \varepsilon} w_{n, \varepsilon}^{*}\right)^{2}=\sum_{a_{1} \cdots a_{n} \in T_{n, \varepsilon}} p\left(a_{1}\right) \cdots p\left(a_{n}\right) \tag{6.240}
\end{equation*}
$$

which approaches 1 in the limit as $n$ approaches infinity, it follows that $\mathrm{E}_{\mathrm{C}}(\mathrm{X}: \mathrm{Y}) \leq \alpha$. As this is so for all $\alpha>\mathrm{H}(p)$, the inequality $\mathrm{E}_{\mathrm{C}}(\mathrm{X}: \mathrm{Y}) \leq \mathrm{H}(p)$ follows.

Next, consider the distillable entanglement of ( $\mathrm{X}, \mathrm{Y}$ ) with respect to the state $u u^{*}$. If $\mathrm{H}(p)=0$, then there is nothing to prove, as the distillable entanglement is trivially nonnegative, so it will be assumed hereafter that $\mathrm{H}(p)>0$. Let $\alpha$ be a real number such that $\alpha<\mathrm{H}(p)$, and let $\varepsilon \in(0,1)$ be
sufficiently small so that $\alpha<\mathrm{H}(p)-2 \varepsilon$. Consider any choice of a positive integer $n \geq-\log (1-\varepsilon) / \varepsilon$, and let $m=\lfloor\alpha n\rfloor$. It holds that

$$
\begin{equation*}
m \leq n(\mathrm{H}(p)-\varepsilon)+\log (1-\varepsilon) \tag{6.241}
\end{equation*}
$$

and therefore

$$
\begin{equation*}
\frac{2^{-n(\mathrm{H}(p)-\varepsilon)}}{1-\varepsilon} \leq 2^{-m} \tag{6.242}
\end{equation*}
$$

As the quantity

$$
\begin{equation*}
\left\|v_{n, \varepsilon}\right\|^{2}=\sum_{a_{1} \cdots a_{n} \in T_{n, \varepsilon}} p\left(a_{1}\right) \cdots p\left(a_{n}\right) \tag{6.243}
\end{equation*}
$$

approaches 1 in the limit as $n$ approaches infinity, it follows that

$$
\begin{equation*}
\frac{2^{-n(\mathrm{H}(p)-\varepsilon)}}{\left\|v_{n, \varepsilon}\right\|^{2}} \leq 2^{-m} \tag{6.244}
\end{equation*}
$$

for all but finitely many positive integers $n$.
Now, consider any choice of $n$ for which (6.244) holds (where $m=\lfloor\alpha n\rfloor$ as before). One therefore has

$$
\begin{equation*}
\sum_{j=1}^{k} \lambda_{j}\left(\operatorname{Tr}_{\mathcal{Y}^{\otimes n}}\left(w_{n, \varepsilon} w_{n, \varepsilon}^{*}\right)\right) \leq \sum_{j=1}^{k} \lambda_{j}\left(\operatorname{Tr}_{\mathcal{W}^{\otimes m}}\left(V_{m} \tau^{\otimes m} V_{m}^{*}\right)\right) \tag{6.245}
\end{equation*}
$$

for every $k \in\left\{1, \ldots, 2^{m}\right\}$. Again using Corollary 6.34 , one has that there must exist an LOCC channel

$$
\begin{equation*}
\Phi_{n} \in \operatorname{LOCC}\left(\mathcal{X}^{\otimes n}, \mathcal{Z}^{\otimes m}: \mathcal{Y}^{\otimes n}, \mathcal{W}^{\otimes m}\right) \tag{6.246}
\end{equation*}
$$

such that

$$
\begin{equation*}
\Phi_{n}\left(w_{n, \varepsilon} w_{n, \varepsilon}^{*}\right)=V_{m} \tau^{\otimes m} V_{m}^{*} \tag{6.247}
\end{equation*}
$$

Making use of the monotonicity of the fidelity function under the action of any channel (Theorem 3.27), one finds that

$$
\begin{align*}
& \mathrm{F}\left(\Phi_{n}\left(U_{n}\left(u u^{*}\right)^{\otimes n} U_{n}^{*}\right), V_{m} \tau^{\otimes m} V_{m}^{*}\right)^{2} \\
& \quad=\mathrm{F}\left(\Phi_{n}\left(U_{n}\left(u u^{*}\right)^{\otimes n} U_{n}^{*}\right), \Phi_{n}\left(w_{n, \varepsilon} w_{n, \varepsilon}^{*}\right)\right)^{2} \\
& \quad \geq \mathrm{F}\left(U_{n}\left(u u^{*}\right)^{\otimes n} U_{n}^{*}, w_{n, \varepsilon} w_{n, \varepsilon}^{*}\right)^{2}  \tag{6.248}\\
& \quad=\sum_{a_{1} \cdots a_{n} \in T_{n, \varepsilon}} p\left(a_{1}\right) \cdots p\left(a_{n}\right) .
\end{align*}
$$

The quantity on the right-hand side of this inequality approaches 1 in the
limit as $n$ approaches infinity, from which it follows that $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y}) \geq \alpha$. As this is so for all $\alpha<\mathrm{H}(p)$, one concludes that $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y}) \geq \mathrm{H}(p)$.

It has been proved that

$$
\begin{equation*}
\mathrm{E}_{\mathrm{C}}(\mathrm{X}: \mathrm{Y}) \leq \mathrm{H}(p) \leq \mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y}) \tag{6.249}
\end{equation*}
$$

The inequality $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y}) \leq \mathrm{E}_{\mathrm{C}}(\mathrm{X}: \mathrm{Y})$ holds by Proposition 6.37, so the proof is complete.

Remark For a given unit vector $u \in \mathcal{X} \otimes \mathcal{Y}$, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, the quantity in (6.226) is known as the entanglement entropy of the pure state $u u^{*}$.

### 6.2.3 Bound entanglement and partial transposition

Informally speaking, Theorem 6.38 implies that all pure state entanglement is equivalent in the bipartite setting. A bipartite pure state is entangled if and only if it has positive entanglement entropy. Moreover, given any two entangled pure states, one necessarily has that an approximate conversion between a large number of copies of the first state and the second state is possible through the use of an LOCC channel, at a rate determined by the ratio of the entanglement entropies of the two states.

The situation is more complex for mixed states. One respect in which this is so is that there exist entangled states having no distillable entanglement. The entanglement in such states, which is referred to as bound entanglement, can never be converted into pure state entanglement through the use of an LOCC channel. The fact that states of this sort exist may be proved through the use of properties of the transpose mapping.

The partial transpose and separability
For any complex Euclidean space $\mathcal{X}$, the transpose mapping $\mathrm{T} \in \mathrm{T}(\mathcal{X})$ is defined as

$$
\begin{equation*}
\mathrm{T}(X)=X^{\top} \tag{6.250}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X})$. As this is a positive map, it follows by the Horodecki criterion (Theorem 6.9) that

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(R) \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.251}
\end{equation*}
$$

for every separable operator $R \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y})$. If $P \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ is a positive
semidefinite operator for which

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P) \notin \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.252}
\end{equation*}
$$

then one may therefore conclude that $P$ is not separable.
The converse of this statement does not hold in general. Given a positive semidefinite operator $P \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ for which

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P) \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.253}
\end{equation*}
$$

one may not conclude that $P$ is separable; an example of a non-separable operator possessing the property $(6.253)$ is described below.

It is the case, however, that an operator $P \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ satisfying the condition (6.253) is highly constrained, in some sense, with respect to the way it is entangled. With this idea in mind, one defines the sets of PPT operators and PPT states (short for positive partial transpose operators and states) as follows.

Definition 6.39 For any choice of complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, the set $\operatorname{PPT}(\mathcal{X}: \mathcal{Y})$ is defined as the set of all operators $P \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ that satisfy

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P) \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.254}
\end{equation*}
$$

Elements of the set $\operatorname{PPT}(\mathcal{X}: \mathcal{Y})$ are called $P P T$ operators, while elements of the set $\operatorname{PPT}(\mathcal{X}: \mathcal{Y}) \cap \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ are called PPT states.

Unextendable product sets and non-separable PPT operators
One method by which non-separable PPT operators may be constructed involves the notion of an unextendable product set. For complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, an orthonormal collection of vectors of the form

$$
\begin{equation*}
\mathcal{A}=\left\{u_{1} \otimes v_{1}, \ldots, u_{m} \otimes v_{m}\right\} \tag{6.255}
\end{equation*}
$$

for unit vectors $u_{1}, \ldots, u_{m} \in \mathcal{X}$ and $v_{1}, \ldots, v_{m} \in \mathcal{Y}$, is an unextendable product set if two properties hold:

1. $\mathcal{A}$ spans a proper subspace of $\mathcal{X} \otimes \mathcal{Y}$. (Equivalently, $m<\operatorname{dim}(\mathcal{X} \otimes \mathcal{Y})$.)
2. For every choice of vectors $x \in \mathcal{X}$ and $y \in \mathcal{Y}$ satisfying $x \otimes y \perp \mathcal{A}$, it must hold that $x \otimes y=0$.

Example 6.40 Define unit vectors $u_{1}, \ldots, u_{5} \in \mathcal{X}$ and $v_{1}, \ldots, v_{5} \in \mathcal{Y}$, for $\mathcal{X}=\mathbb{C}^{3}$ and $\mathcal{Y}=\mathbb{C}^{3}$, as follows:

$$
\begin{align*}
u_{1} & =e_{1}, & v_{1} & =\frac{1}{\sqrt{2}}\left(e_{1}-e_{2}\right), \\
u_{2} & =e_{3}, & v_{2} & =\frac{1}{\sqrt{2}}\left(e_{2}-e_{3}\right), \\
u_{3} & =\frac{1}{\sqrt{2}}\left(e_{1}-e_{2}\right), & v_{3} & =e_{3}, \\
u_{4} & =\frac{1}{\sqrt{2}}\left(e_{2}-e_{3}\right), & v_{4} & =e_{1},  \tag{6.256}\\
u_{5} & =\frac{1}{\sqrt{3}}\left(e_{1}+e_{2}+e_{3}\right), & v_{5} & =\frac{1}{\sqrt{3}}\left(e_{1}+e_{2}+e_{3}\right) .
\end{align*}
$$

It therefore holds, for each $k \in\{1, \ldots, 5\}$, that $u_{k} \otimes v_{k}=\operatorname{vec}\left(A_{k}\right)$ for

$$
\begin{gather*}
A_{1}=\frac{1}{\sqrt{2}}\left(\begin{array}{ccc}
1 & -1 & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right), \quad A_{2}=\frac{1}{\sqrt{2}}\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
0 & 1 & -1
\end{array}\right) \\
A_{3}=\frac{1}{\sqrt{2}}\left(\begin{array}{ccc}
0 & 0 & 1 \\
0 & 0 & -1 \\
0 & 0 & 0
\end{array}\right), \quad A_{4}=\frac{1}{\sqrt{2}}\left(\begin{array}{ccc}
0 & 0 & 0 \\
1 & 0 & 0 \\
-1 & 0 & 0
\end{array}\right)  \tag{6.257}\\
A_{5}=\frac{1}{3}\left(\begin{array}{lll}
1 & 1 & 1 \\
1 & 1 & 1 \\
1 & 1 & 1
\end{array}\right) .
\end{gather*}
$$

The set

$$
\begin{equation*}
\mathcal{A}=\left\{u_{1} \otimes v_{1}, \ldots, u_{5} \otimes v_{5}\right\} \tag{6.258}
\end{equation*}
$$

is orthonormal by inspection. If $x \in \mathcal{X}$ and $y \in \mathcal{Y}$ satisfy

$$
\begin{equation*}
\left\langle x \otimes y, u_{k} \otimes v_{k}\right\rangle=\left\langle x, u_{k}\right\rangle\left\langle y, v_{k}\right\rangle=0 \tag{6.259}
\end{equation*}
$$

for $k=1, \ldots, 5$, then one must have $\left\langle x, u_{k}\right\rangle=0$ for at least 3 distinct choices of $k \in\{1, \ldots, 5\}$ or $\left\langle y, v_{k}\right\rangle=0$ for at least 3 distinct choices of $k \in\{1, \ldots, 5\}$. As every 3 distinct choices of $u_{k}$ span all of $\mathcal{X}$ and every 3 distinct choices of $v_{k}$ span all of $\mathcal{Y}$, it follows that $x \otimes y=0$. The set $\mathcal{A}$ is therefore an unextendable product set.

The projection onto the subspace orthogonal to an unextendable product set must be both PPT and entangled, as the following theorem states.

Theorem 6.41 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, let

$$
\begin{equation*}
\mathcal{A}=\left\{u_{1} \otimes v_{1}, \ldots, u_{m} \otimes v_{m}\right\} \tag{6.260}
\end{equation*}
$$

be an unextendable product set in $\mathcal{X} \otimes \mathcal{Y}$, and define

$$
\begin{equation*}
\Pi=\sum_{k=1}^{m} u_{k} u_{k}^{*} \otimes v_{k} v_{k}^{*} . \tag{6.261}
\end{equation*}
$$

It holds that

$$
\begin{equation*}
\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y}) \backslash \operatorname{Sep}(\mathcal{X}: \mathcal{Y}) . \tag{6.262}
\end{equation*}
$$

Proof From the assumption that $\mathcal{A}$ is an orthonormal set, one may conclude that $\left\{\overline{u_{1}} \otimes v_{1}, \ldots, \overline{u_{k}} \otimes v_{k}\right\}$ is an orthonormal set as well. It follows that

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\nu)}\right)(\Pi)=\sum_{k=1}^{m} \overline{u_{k}} u_{k}^{\top} \otimes v_{k} v_{k}^{*} \tag{6.263}
\end{equation*}
$$

is a projection operator, and therefore

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(\Pi) \leq \mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}} \tag{6.264}
\end{equation*}
$$

As

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}\right)=\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}, \tag{6.265}
\end{equation*}
$$

one obtains the inclusion

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi\right) \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.266}
\end{equation*}
$$

It therefore holds that

$$
\begin{equation*}
\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y}) . \tag{6.267}
\end{equation*}
$$

Now, toward a contradiction, assume that

$$
\begin{equation*}
\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi \in \operatorname{Sep}(\mathcal{X}: \mathcal{Y}) \tag{6.268}
\end{equation*}
$$

which implies that

$$
\begin{equation*}
\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi=\sum_{a \in \Sigma} x_{a} x_{a}^{*} \otimes y_{a} y_{a}^{*} \tag{6.269}
\end{equation*}
$$

for some choice of an alphabet $\Sigma$ and collections $\left\{x_{a}: a \in \Sigma\right\} \subset \mathcal{X}$ and $\left\{y_{a}: a \in \Sigma\right\} \subset \mathcal{Y}$. It holds that

$$
\begin{align*}
& \sum_{k=1}^{m} \sum_{a \in \Sigma}\left|\left\langle x_{a} \otimes y_{a}, u_{k} \otimes v_{k}\right\rangle\right|^{2} \\
& \quad=\sum_{k=1}^{m}\left(u_{k} \otimes v_{k}\right)^{*}\left(\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi\right)\left(u_{k} \otimes v_{k}\right)=0, \tag{6.270}
\end{align*}
$$

and therefore

$$
\begin{equation*}
\left\langle x_{a} \otimes y_{a}, u_{k} \otimes v_{k}\right\rangle=0 \tag{6.271}
\end{equation*}
$$

for every $a \in \Sigma$ and $k \in\{1, \ldots, m\}$. As $\mathcal{A}$ is an unextendable product set, it follows that $x_{a} \otimes y_{a}=0$ for every $a \in \Sigma$, and therefore

$$
\begin{equation*}
\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi=0 \tag{6.272}
\end{equation*}
$$

This, however, is in contradiction with the assumption $m<\operatorname{dim}(\mathcal{X} \otimes \mathcal{Y})$. It follows that

$$
\begin{equation*}
\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{Y}}-\Pi \notin \operatorname{Sep}(\mathcal{X}: \mathcal{Y}) \tag{6.273}
\end{equation*}
$$

which completes the proof.

## PPT states have no distillable entanglement

PPT states may not always be separable, but they exhibit similar properties to separable states in some respects. One such respect is that their overlap with every maximally entangled state is small. The next proposition, which is reminiscent of Proposition 6.15, is representative of this fact.

Proposition 6.42 Let $A \in \mathrm{~L}(\mathcal{Y}, \mathcal{X})$ be an operator satisfying $\|A\| \leq 1$, for $\mathcal{X}$ and $\mathcal{Y}$ being complex Euclidean spaces. For every $P \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y})$ it holds that

$$
\begin{equation*}
\left\langle\operatorname{vec}(A) \operatorname{vec}(A)^{*}, P\right\rangle \leq \operatorname{Tr}(P) \tag{6.274}
\end{equation*}
$$

Proof The transpose mapping is its own adjoint and inverse, and therefore

$$
\begin{align*}
& \left\langle\operatorname{vec}(A) \operatorname{vec}(A)^{*}, P\right\rangle  \tag{6.275}\\
& \quad=\left\langle\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\operatorname{vec}(A) \operatorname{vec}(A)^{*}\right),\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P)\right\rangle .
\end{align*}
$$

It holds that

$$
\begin{equation*}
\operatorname{vec}(A)=\left(\mathbb{1}_{\mathcal{X}} \otimes A^{\top}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \tag{6.276}
\end{equation*}
$$

which implies that

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\operatorname{vec}(A) \operatorname{vec}(A)^{*}\right)=\left(\mathbb{1}_{\mathcal{X}} \otimes A^{\top}\right) W\left(\mathbb{1}_{\mathcal{X}} \otimes \bar{A}\right) \tag{6.277}
\end{equation*}
$$

for $W \in \mathrm{U}(\mathcal{X} \otimes \mathcal{X})$ denoting the swap operator on $\mathcal{X} \otimes \mathcal{X}$. The operator represented by (6.277) has spectral norm at most 1, and therefore

$$
\begin{align*}
&\left\langle\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)\left(\operatorname{vec}(A) \operatorname{vec}(A)^{*}\right)\right.\left.\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P)\right\rangle \\
& \leq\left\|\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P)\right\|_{1} . \tag{6.278}
\end{align*}
$$

Finally, because $P \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y})$, together with the observation that the transpose mapping preserves trace, one has

$$
\begin{equation*}
\left\|\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P)\right\|_{1}=\operatorname{Tr}(P) \tag{6.279}
\end{equation*}
$$

The proposition follows from (6.275), (6.278), and (6.279).
Example 6.43 Similar to Example 6.16, let $\Sigma$ be an alphabet, let $n=|\Sigma|$, and let $\mathcal{X}=\mathbb{C}^{\Sigma}$ and $\mathcal{Y}=\mathbb{C}^{\Sigma}$. Define a density operator $\tau \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ as

$$
\begin{equation*}
\tau=\frac{1}{n} \sum_{a, b \in \Sigma} E_{a, b} \otimes E_{a, b}=\frac{1}{n} \operatorname{vec}(\mathbb{1}) \operatorname{vec}(\mathbb{1})^{*} \tag{6.280}
\end{equation*}
$$

where $\mathbb{1}$ denotes the identity operator on $\mathbb{C}^{\Sigma}$, which may be viewed as an element of the set $\mathrm{L}(\mathcal{Y}, \mathcal{X})$. For every PPT state

$$
\begin{equation*}
\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y}) \cap \operatorname{PPT}(\mathcal{X}: \mathcal{Y}) \tag{6.281}
\end{equation*}
$$

it holds that

$$
\begin{equation*}
\langle\tau, \rho\rangle=\frac{1}{n}\left\langle\operatorname{vec}(\mathbb{1}) \operatorname{vec}(\mathbb{1})^{*}, \rho\right\rangle \leq \frac{1}{n} \tag{6.282}
\end{equation*}
$$

by Proposition 6.42. Thus, with respect to their overlap with the maximally entangled state $\tau$, one has that PPT operators are bounded in a similar way to separable operators.

Proposition 6.42 , when combined with the next proposition stating that separable maps (and therefore LOCC channels) map PPT operators to PPT operators, leads to a proof that PPT states have distillable entanglement equal to zero.

Proposition 6.44 Let $\mathcal{X}, \mathcal{Y}, \mathcal{Z}$, and $\mathcal{W}$ be complex Euclidean spaces, let $P \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y})$ be a PPT operator, and let $\Phi \in \operatorname{SepCP}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ be a separable map. It holds that $\Phi(P) \in \operatorname{PPT}(\mathcal{Z}: \mathcal{W})$.

Proof For any choice of operators $A \in \mathrm{~L}(\mathcal{X}, \mathcal{Z})$ and $B \in \mathrm{~L}(\mathcal{Y}, \mathcal{W})$, the assumption $P \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y})$ implies that

$$
\begin{align*}
& \left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{W})}\right)\left((A \otimes B) P(A \otimes B)^{*}\right) \\
& \quad=(\bar{A} \otimes B)\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P)(\bar{A} \otimes B)^{*} \in \operatorname{Pos}(\mathcal{Z} \otimes \mathcal{W}) \tag{6.283}
\end{align*}
$$

(In this equation, T refers to the transpose mapping on $\mathcal{Z}$ or $\mathcal{X}$, as the context dictates.) As $\Phi$ is separable, one has

$$
\begin{equation*}
\Phi(X)=\sum_{a \in \Sigma}\left(A_{a} \otimes B_{a}\right) X\left(A_{a} \otimes B_{a}\right)^{*} \tag{6.284}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X} \otimes \mathcal{Y})$, for some choice of an alphabet $\Sigma$ and collections
of operators $\left\{A_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{X}, \mathcal{Z})$ and $\left\{B_{a}: a \in \Sigma\right\} \subset \mathrm{L}(\mathcal{Y}, \mathcal{W})$. Consequently, one has that

$$
\begin{equation*}
\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{W})}\right)(\Phi(P))=\sum_{a \in \Sigma}\left(\overline{A_{a}} \otimes B_{a}\right)\left(\mathrm{T} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Y})}\right)(P)\left(\overline{A_{a}} \otimes B_{a}\right)^{*} \tag{6.285}
\end{equation*}
$$

is positive semidefinite, and therefore $\Phi(P) \in \operatorname{PPT}(\mathcal{Z}: \mathcal{W})$, as required.
Theorem 6.45 Let X and Y be registers and consider a PPT state

$$
\begin{equation*}
\rho \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y}) \cap \mathrm{D}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.286}
\end{equation*}
$$

of the pair $(\mathrm{X}, \mathrm{Y})$. With respect to the state $\rho$, it holds that $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y})=0$.
Proof Let $\Gamma=\{0,1\}$, let $\mathcal{Z}=\mathbb{C}^{\Gamma}$ and $\mathcal{W}=\mathbb{C}^{\Gamma}$, and let $\tau \in \mathrm{D}(\mathcal{Z} \otimes \mathcal{W})$ be defined as

$$
\begin{equation*}
\tau=\frac{1}{2} \sum_{a, b \in \Gamma} E_{a, b} \otimes E_{a, b} \tag{6.287}
\end{equation*}
$$

Suppose $\alpha>0$, let $n$ be any positive integer for which $m=\lfloor\alpha n\rfloor \geq 1$, and consider any LOCC channel $\Phi \in \operatorname{LOCC}\left(\mathcal{X}^{\otimes n}, \mathcal{Z}^{\otimes m}: \mathcal{Y}^{\otimes n}, \mathcal{W}^{\otimes m}\right)$. Recall the operators $U_{n}$ and $V_{m}$ as defined by (6.210) and (6.213). It holds that

$$
\begin{equation*}
U_{n} \rho^{\otimes n} U_{n}^{*} \in \operatorname{PPT}\left(\mathcal{X}^{\otimes n}: \mathcal{Y}^{\otimes n}\right) \tag{6.288}
\end{equation*}
$$

and therefore

$$
\begin{equation*}
\Phi\left(U_{n} \rho^{\otimes n} U_{n}^{*}\right) \in \operatorname{PPT}\left(\mathcal{Z}^{\otimes m}: \mathcal{W}^{\otimes m}\right) \tag{6.289}
\end{equation*}
$$

by Proposition 6.44. One may therefore conclude from Proposition 6.42 that

$$
\begin{equation*}
\mathrm{F}\left(V_{m} \tau^{\otimes m} V_{m}^{*}, \Phi\left(U_{n} \rho^{\otimes n} U_{n}^{*}\right)\right) \leq 2^{-\frac{m}{2}} \leq \frac{1}{\sqrt{2}} \tag{6.290}
\end{equation*}
$$

The number $\alpha$ therefore fails to satisfy the requirements of Definition 6.35. It follows that $\mathrm{E}_{\mathrm{D}}(\mathrm{X}: \mathrm{Y})=0$.

### 6.3 Phenomena associated with entanglement

This section discusses a few notions generally associated with entanglement: teleportation, dense coding, and non-classical correlations. These notions serve as representatives of the sorts of operational effects that entanglement may induce.

### 6.3.1 Teleportation and dense coding

In quantum information theory, teleportation has traditionally referred to a protocol by which a single-qubit quantum channel is implemented through the use of a maximally entangled pair of qubits along with two classical bits of communication. Informally speaking, teleportation suggests the following transformation:

1 pair of maximally entangled qubits
+2 bits of classical communication
$\rightarrow 1$ qubit of quantum communication.
The dense coding protocol offers a trade-off of resources that is, in some sense, complementary to teleportation. Again traditionally speaking, it is a protocol by which a two-bit classical channel is implemented through the use of a maximally entangled pair of qubits and a single-qubit quantum channel. In this case, the suggested transformation is as follows:

1 pair of maximally entangled qubits
+1 qubit of quantum communication
$\rightarrow 2$ bits of classical communication.
In both cases, the maximally entangled pair of qubits is consumed by the conversion between two classical bits and one qubit of communication; in essence, the entangled pair of qubits functions as a resource allowing for this conversion.

In the discussion that follows, teleportation and dense coding will be considered in greater generality. The traditional protocols suggested above will emerge as specific instances of more general classes of protocols.

## Teleportation

Consider the following scenario in which two individuals, Alice and Bob, aim to implement an ideal quantum channel through the combined use of entanglement and classical communication.

Scenario 6.46 (Teleportation) Alice holds a register X and Bob holds Y. Both registers have the same classical state set $\Sigma$, and the state of the pair $(\mathrm{X}, \mathrm{Y})$ is given by the maximally entangled state

$$
\begin{equation*}
\tau=\frac{1}{|\Sigma|} \sum_{b, c \in \Sigma} E_{b, c} \otimes E_{b, c} \tag{6.291}
\end{equation*}
$$

Alice obtains a new register Z, whose classical state set is also $\Sigma$, and she wishes to transmit $\mathbf{Z}$ to Bob. Alice and Bob attempt to accomplish this task
using classical communication together with the shared entangled state $\tau$, by means of a protocol as follows:

1. Alice performs a measurement $\mu: \Gamma \rightarrow \operatorname{Pos}(\mathcal{Z} \otimes \mathcal{X})$ on the pair $(Z, X)$, where $\Gamma$ is an arbitrarily chosen alphabet, and sends the outcome $a \in \Gamma$ of this measurement to Bob.
2. For $\left\{\Psi_{a}: a \in \Gamma\right\} \subseteq \mathrm{C}(\mathcal{Y}, \mathcal{Z})$ being a collection of channels indexed by $\Gamma$, Bob applies the channel $\Psi_{a}$ to Y , for whichever symbol $a \in \Gamma$ was sent to him by Alice, transforming this register into a new register Z .

An analysis will reveal that this protocol accomplishes the task at hand for a suitable choice for Alice's measurement and Bob's collection of channels.

Remark One may consider more general scenarios along similar lines to Scenario 6.46. For instance, $X, Y$, and $Z$ might not share the same classical state set, the initial state of the pair $(X, Y)$ might be initialized to a different state than $\tau$, and Alice and Bob might aim to implement a channel different from the identity channel. The discussion that follows, however, will focus on the setting described in Scenario 6.46 in the interest of simplicity.

For any choice of Alice's measurement $\mu$ and Bob's collection of channels $\left\{\Psi_{a}: a \in \Gamma\right\}$, the channel $\Phi \in \mathrm{C}(\mathcal{Z})$ that is implemented by the protocol described in Scenario 6.46 may be expressed as

$$
\begin{equation*}
\Phi(Z)=\frac{1}{|\Sigma|} \sum_{a \in \Gamma} \sum_{b, c \in \Sigma}\left\langle\mu(a), Z \otimes E_{b, c}\right\rangle \Psi_{a}\left(E_{b, c}\right) \tag{6.292}
\end{equation*}
$$

for all $Z \in \mathrm{~L}(\mathcal{Z})$.
The following theorem provides a characterization of those measurements and collections of channels for which the channel $\Phi$ is equal to the identity channel, which represents an ideal transmission of quantum information from Alice to Bob. (The statement of the theorem refers to a single complex Euclidean space $\mathcal{X}$, rather than $\mathcal{X}, \mathcal{Y}$, and $\mathcal{Z}$, and includes the assumption that none of the measurement operators of $\mu$ are identically zero, as this allows for a simpler statement and proof of the characterization.)

Theorem 6.47 Let $\Sigma$ and $\Gamma$ be alphabets and let $\mathcal{X}=\mathbb{C}^{\Sigma}$ be a complex Euclidean space. Also let

$$
\begin{equation*}
\mu: \Gamma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{X}) \tag{6.293}
\end{equation*}
$$

be a measurement such that $\mu(a) \neq 0$ for every $a \in \Gamma$, and let

$$
\begin{equation*}
\left\{\Psi_{a}: a \in \Gamma\right\} \subseteq \mathrm{C}(\mathcal{X}) \tag{6.294}
\end{equation*}
$$

be a collection of channels. The following two statements are equivalent:

1. It holds that

$$
\begin{equation*}
X=\frac{1}{|\Sigma|} \sum_{a \in \Gamma} \sum_{b, c \in \Sigma}\left\langle\mu(a), X \otimes E_{b, c}\right\rangle \Psi_{a}\left(E_{b, c}\right) \tag{6.295}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$.
2. There exists a collection $\left\{U_{a}: a \in \Gamma\right\} \subset \mathrm{U}(\mathcal{X})$ of unitary operators and a probability vector $p \in \mathcal{P}(\Gamma)$ such that

$$
\begin{equation*}
\mu(a)=p(a)|\Sigma| \operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*} \quad \text { and } \quad \Psi_{a}(X)=U_{a} X U_{a}^{*} \tag{6.296}
\end{equation*}
$$

for every choice of $a \in \Gamma$ and $X \in \mathrm{~L}(\mathcal{X})$.
The proof of Theorem 6.47 will make use of the following proposition, which establishes that a channel of the form $\Phi \in \mathrm{C}(\mathcal{X})$, for any complex Euclidean space $\mathcal{X}$, can be the inverse of a completely positive map only if it is a unitary channel.

Proposition 6.48 Let $\mathcal{X}$ be a complex Euclidean space, let $\Phi \in \mathrm{C}(\mathcal{X})$ be a channel, and let $\Psi \in \mathrm{CP}(\mathcal{X})$ be a completely positive map for which $\Phi \Psi=\mathbb{1}_{\mathrm{L}(\mathcal{X})}$. There exists a unitary operator $U \in \mathrm{U}(\mathcal{X})$ such that

$$
\begin{equation*}
\Phi(X)=U^{*} X U \quad \text { and } \quad \Psi(X)=U X U^{*} \tag{6.297}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X})$.
Proof As $\Psi$ is completely positive, and evidently nonzero, its Choi operator $J(\Psi)$ is a nonzero positive semidefinite operator. By the spectral theorem (Corollary 1.4), it is therefore possible to write

$$
\begin{equation*}
J(\Psi)=\sum_{k=1}^{r} \operatorname{vec}\left(A_{k}\right) \operatorname{vec}\left(A_{k}\right)^{*} \tag{6.298}
\end{equation*}
$$

for $r=\operatorname{rank}(J(\Psi))$ and $\left\{A_{1}, \ldots, A_{r}\right\} \subset \mathrm{L}(\mathcal{X})$ being an orthogonal collection of nonzero operators. Consequently, one has

$$
\begin{align*}
& \sum_{k=1}^{r}\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)\left(\operatorname{vec}\left(A_{k}\right) \operatorname{vec}\left(A_{k}\right)^{*}\right)  \tag{6.299}\\
& \quad=\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)(J(\Psi))=J(\Phi \Psi)=\operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right)^{*}
\end{align*}
$$

As $\operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right)^{*}$ has rank equal to one, and each operator

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)\left(\operatorname{vec}\left(A_{k}\right) \operatorname{vec}\left(A_{k}\right)^{*}\right) \tag{6.300}
\end{equation*}
$$

is positive semidefinite (by the complete positivity of $\Phi$ ), it follows that there must exist a probability vector $\left(p_{1}, \ldots, p_{r}\right)$ such that

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)\left(\operatorname{vec}\left(A_{k}\right) \operatorname{vec}\left(A_{k}\right)^{*}\right)=p_{k} \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right)^{*} \tag{6.301}
\end{equation*}
$$

for each $k \in\{1, \ldots, r\}$. Because $\Phi$ preserves trace, it follows that

$$
\begin{equation*}
\left(A_{k}^{*} A_{k}\right)^{\top}=\left(\operatorname{Tr} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)\left(\operatorname{vec}\left(A_{k}\right) \operatorname{vec}\left(A_{k}\right)^{*}\right)=p_{k} \mathbb{1}_{\mathcal{X}} \tag{6.302}
\end{equation*}
$$

and therefore $A_{k}=\sqrt{p_{k}} U_{k}$ for some choice of a unitary operator $U_{k} \in \mathrm{U}(\mathcal{X})$, for each $k \in\{1, \ldots, r\}$. This implies that

$$
\begin{align*}
\left(\mathbb{1}_{\mathcal{X}} \otimes U_{k}^{\top}\right) J(\Phi)\left(\mathbb{1}_{\mathcal{X}} \otimes U_{k}^{\top}\right)^{*} & =\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)\left(\operatorname{vec}\left(U_{k}\right) \operatorname{vec}\left(U_{k}\right)^{*}\right)  \tag{6.303}\\
& =\operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right)^{*}
\end{align*}
$$

and therefore

$$
\begin{equation*}
J(\Phi)=\operatorname{vec}\left(U_{k}^{*}\right) \operatorname{vec}\left(U_{k}^{*}\right)^{*} \tag{6.304}
\end{equation*}
$$

again for each $k \in\{1, \ldots, r\}$. As $\left\{A_{1}, \ldots, A_{r}\right\}$ is a collection of nonzero, orthogonal operators, and is therefore linearly independent, one concludes that $r=1$ and $p_{1}=1$; and by setting $U=U_{1}$ the proposition is proved.

Proof of Theorem 6.47 Assume that statement 1 holds. For each $a \in \Gamma$, define a map $\Xi_{a} \in \mathrm{~T}(\mathcal{X})$ as

$$
\begin{equation*}
\Xi_{a}(X)=\frac{1}{|\Sigma|} \sum_{b, c \in \Sigma}\left\langle\mu(a), X \otimes E_{b, c}\right\rangle E_{b, c} \tag{6.305}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X})$. The Choi operator of $\Xi_{a}$ is given by

$$
\begin{equation*}
J\left(\Xi_{a}\right)=\frac{1}{|\Sigma|} W \overline{\mu(a)} W \tag{6.306}
\end{equation*}
$$

for $W \in \mathrm{U}(\mathcal{X} \otimes \mathcal{X})$ denoting the swap operator. As $J\left(\Xi_{a}\right) \in \operatorname{Pos}(\mathcal{X} \otimes \mathcal{X})$ for each $a \in \Gamma$, it follows that $\Xi_{a}$ is completely positive, and moreover is nonzero by the assumption that $\mu(a)$ is nonzero. Statement 1 may now be expressed as

$$
\begin{equation*}
\sum_{a \in \Gamma} \Psi_{a} \Xi_{a}=\mathbb{1}_{\mathrm{L}(\mathcal{X})}, \tag{6.307}
\end{equation*}
$$

which is equivalent to

$$
\begin{equation*}
\sum_{a \in \Gamma} J\left(\Psi_{a} \Xi_{a}\right)=\operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right)^{*} \tag{6.308}
\end{equation*}
$$

As the composition $\Psi_{a} \Xi_{a}$ is necessarily completely positive and nonzero for each $a \in \Gamma$, and the operator $\operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right)^{*}$ has rank equal to 1 , it follows that there must exist a probability vector $p \in \mathcal{P}(\Gamma)$ such that

$$
\begin{equation*}
J\left(\Psi_{a} \Xi_{a}\right)=p(a) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{X}}\right)^{*} \tag{6.309}
\end{equation*}
$$

for each $a \in \Gamma$. Consequently,

$$
\begin{equation*}
\frac{\left(\Psi_{a} \Xi_{a}\right)(X)}{p(a)}=X \tag{6.310}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$. By Proposition 6.48 , there must exist a collection of unitary operators $\left\{U_{a}: a \in \Gamma\right\} \subset \mathrm{U}(\mathcal{X})$ such that

$$
\begin{equation*}
\Psi_{a}(X)=U_{a} X U_{a}^{*} \quad \text { and } \quad \frac{1}{p(a)} \Xi_{a}(X)=U_{a}^{*} X U_{a} \tag{6.311}
\end{equation*}
$$

for every $a \in \Gamma$ and $X \in \mathrm{~L}(\mathcal{X})$. Thus,

$$
\begin{equation*}
\frac{1}{|\Sigma|} W \overline{\mu(a)} W=J\left(\Xi_{a}\right)=p(a) \operatorname{vec}\left(U_{a}^{*}\right) \operatorname{vec}\left(U_{a}^{*}\right)^{*} \tag{6.312}
\end{equation*}
$$

and because $W \operatorname{vec}(Y)=\operatorname{vec}\left(Y^{\boldsymbol{\top}}\right)$ for every $Y \in \mathrm{~L}(\mathcal{X})$, one therefore has

$$
\begin{equation*}
\mu(a)=p(a)|\Sigma| \operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*} \tag{6.313}
\end{equation*}
$$

for each $a \in \Gamma$. Statement 1 therefore implies statement 2 .
Now assume statement 2 holds. As $\mu$ is assumed to be a measurement, it must be the case that

$$
\begin{equation*}
\sum_{a \in \Gamma} p(a) \operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*}=\frac{1}{|\Sigma|} \mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{X}} \tag{6.314}
\end{equation*}
$$

The operator represented by the equation (6.314) coincides with the Choi operator $J(\Omega)$ of the completely depolarizing channel $\Omega \in \mathrm{C}(\mathcal{X})$. It follows that one may write

$$
\begin{equation*}
\Omega(X)=\sum_{a \in \Gamma} p(a) U_{a} X U_{a}^{*} \tag{6.315}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$. Because the natural representation of the completely depolarizing channel is given by

$$
\begin{equation*}
K(\Omega)=\frac{1}{|\Sigma|} \sum_{b, c \in \Sigma} E_{b, c} \otimes E_{b, c} \tag{6.316}
\end{equation*}
$$

one finds that

$$
\begin{equation*}
\sum_{a \in \Gamma} p(a) \overline{U_{a}} \otimes U_{a}=\overline{K(\Omega)}=\frac{1}{|\Sigma|} \sum_{b, c \in \Sigma} E_{b, c} \otimes E_{b, c} \tag{6.317}
\end{equation*}
$$

by Proposition 2.20.
Now consider the channel $\Phi \in \mathrm{C}(\mathcal{X})$ defined by

$$
\begin{equation*}
\Phi(X)=\frac{1}{|\Sigma|} \sum_{a \in \Gamma} \sum_{b, c \in \Sigma}\left\langle\mu(a), X \otimes E_{b, c}\right\rangle \Psi_{a}\left(E_{b, c}\right) \tag{6.318}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$. Making use of the expression (6.317), one may write

$$
\begin{equation*}
\Phi(X)=\sum_{a, b \in \Gamma} p(b)\left\langle\mu(a), X \otimes \overline{U_{b}}\right\rangle \Psi_{a}\left(U_{b}\right) \tag{6.319}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$. By substituting according to (6.296), one obtains

$$
\begin{align*}
\Phi(X) & =|\Sigma| \sum_{a, b \in \Gamma} p(a) p(b) \operatorname{vec}\left(U_{a}\right)^{*}\left(X \otimes \overline{U_{b}}\right) \operatorname{vec}\left(U_{a}\right) U_{a} U_{b} U_{a}^{*} \\
& =|\Sigma| \sum_{a, b \in \Gamma} p(a) p(b)\left\langle U_{a} U_{b} U_{a}^{*}, X\right\rangle U_{a} U_{b} U_{a}^{*} \tag{6.320}
\end{align*}
$$

The natural representation $K(\Phi)$ of the channel $\Phi$ is therefore given by

$$
\begin{align*}
& |\Sigma| \sum_{a, b \in \Gamma} p(a) p(b) \operatorname{vec}\left(U_{a} U_{b} U_{a}^{*}\right) \operatorname{vec}\left(U_{a} U_{b} U_{a}^{*}\right)^{*} \\
& \quad=\sum_{a \in \Gamma} p(a)\left(U_{a} \otimes \overline{U_{a}}\right)\left(|\Sigma| \sum_{b \in \Gamma} p(b) \operatorname{vec}\left(U_{b}\right) \operatorname{vec}\left(U_{b}\right)^{*}\right)\left(U_{a} \otimes \overline{U_{a}}\right)^{*}  \tag{6.321}\\
& \quad=\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{X}}
\end{align*}
$$

where the last equality has made use of (6.314). It follows that $\Phi$ is equal to the identity channel, and therefore statement 2 implies statement 1.

Theorem 6.47 implies that every mixed-unitary representation of the completely depolarizing channel gives rise to a teleportation protocol, as the following corollary makes precise.
Corollary 6.49 Let $\Sigma$ and $\Gamma$ be alphabets, let $\mathcal{X}=\mathbb{C}^{\Sigma}$, let

$$
\begin{equation*}
\left\{U_{a}: a \in \Gamma\right\} \subset \mathrm{U}(\mathcal{X}) \tag{6.322}
\end{equation*}
$$

be a collection of unitary operators, let $p \in \mathcal{P}(\Gamma)$ be a probability vector, and assume that

$$
\begin{equation*}
\Omega(X)=\sum_{a \in \Gamma} p(a) U_{a} X U_{a}^{*} \tag{6.323}
\end{equation*}
$$

for every $X \in \mathrm{~L}(\mathcal{X})$, where $\Omega \in \mathrm{C}(\mathcal{X})$ denotes the completely depolarizing channel with respect to the space $\mathcal{X}$. For $\mu: \Gamma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{X})$ defined as

$$
\begin{equation*}
\mu(a)=p(a)|\Sigma| \operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*} \tag{6.324}
\end{equation*}
$$

for each $a \in \Gamma$, one has that $\mu$ is a measurement, and moreover

$$
\begin{equation*}
X=\frac{1}{|\Sigma|} \sum_{a \in \Gamma} \sum_{b, c \in \Sigma}\left\langle\mu(a), X \otimes E_{b, c}\right\rangle U_{a} E_{b, c} U_{a}^{*} \tag{6.325}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X})$.

Proof There is no loss of generality in assuming that $p(a) \neq 0$ for every $a \in \Gamma$, for otherwise one could define an alphabet $\Gamma_{0}=\{a \in \Gamma: p(a) \neq 0\}$, verify that the corollary holds in this case, and observe that the statement of the corollary is equivalent when $\Gamma$ is replaced by $\Gamma_{0}$ in this way.

It is evident that $\mu$ is a measurement, as each $\mu(a)$ is positive semidefinite and it holds that

$$
\begin{equation*}
\sum_{a \in \Gamma} \mu(a)=\sum_{a \in \Gamma} p(a)|\Sigma| \operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*}=|\Sigma| J(\Omega)=\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{X}} \tag{6.326}
\end{equation*}
$$

By defining $\Psi_{a}(X)=U_{a} X U_{a}^{*}$ for every $X \in \mathrm{~L}(\mathcal{X})$ and $a \in \Gamma$, one has that statement 2 of Theorem 6.47 is satisfied. This implies that statement 1 of that theorem holds, which is equivalent to (6.325), and therefore completes the proof.

Example 6.50 Let $\Gamma=\Sigma \times \Sigma$, for $\Sigma=\{0,1\}$ denoting the binary alphabet. Elements of $\Gamma$ will be viewed as binary strings of length 2 for convenience. Define $p \in \mathcal{P}(\Gamma)$ as

$$
\begin{equation*}
p(00)=p(01)=p(10)=p(11)=\frac{1}{4} \tag{6.327}
\end{equation*}
$$

and define unitary operators $U_{00}, U_{01}, U_{10}, U_{11} \in \mathrm{U}\left(\mathbb{C}^{\Sigma}\right)$ as follows:

$$
\begin{align*}
& U_{00}=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right), \quad U_{01}=\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right) \\
& U_{10}=\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right), \quad U_{11}=\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right) \tag{6.328}
\end{align*}
$$

The operators $U_{00}, U_{01}, U_{10}, U_{11}$ coincide with the discrete Weyl operators acting on the space $\mathbb{C}^{\Sigma}$, and (as explained in Section 4.1.2) provide a mixedunitary realization of the completely depolarizing channel $\Omega \in \mathrm{C}\left(\mathbb{C}^{\Sigma}\right)$ :

$$
\begin{equation*}
\frac{1}{4} \sum_{a, b \in \Sigma} U_{a b} X U_{a b}^{*}=\frac{\operatorname{Tr}(X)}{2} \mathbb{1} \tag{6.329}
\end{equation*}
$$

for every $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$. Consequently, by taking $\mu: \Gamma \rightarrow \operatorname{Pos}\left(\mathbb{C}^{\Sigma} \otimes \mathbb{C}^{\Sigma}\right)$ to be the measurement defined as

$$
\begin{array}{ll}
\mu(00)=\frac{\operatorname{vec}\left(U_{00}\right) \operatorname{vec}\left(U_{00}\right)^{*}}{2}, & \mu(01)=\frac{\operatorname{vec}\left(U_{01}\right) \operatorname{vec}\left(U_{01}\right)^{*}}{2}  \tag{6.330}\\
\mu(10)=\frac{\operatorname{vec}\left(U_{10}\right) \operatorname{vec}\left(U_{10}\right)^{*}}{2}, & \mu(11)=\frac{\operatorname{vec}\left(U_{11}\right) \operatorname{vec}\left(U_{11}\right)^{*}}{2}
\end{array}
$$

or equivalently $\mu(a b)=u_{a b} u_{a b}^{*}$ for

$$
\begin{array}{ll}
u_{00}=\frac{e_{00}+e_{11}}{\sqrt{2}}, & u_{01}=\frac{e_{00}-e_{11}}{\sqrt{2}},  \tag{6.331}\\
u_{10}=\frac{e_{01}+e_{10}}{\sqrt{2}}, & u_{11}=\frac{e_{01}-e_{10}}{\sqrt{2}},
\end{array}
$$

and setting $\Psi_{a b}(X)=U_{a b} X U_{a b}^{*}$ for each $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$ and $a, b \in \Sigma$, one obtains a teleportation protocol as described in Scenario 6.46. Indeed, the resulting protocol is equivalent to the traditional notion of teleportation in which an ideal single-qubit channel is implemented using a maximally entangled pair of qubits along with two classical bits of communication. The set $\left\{u_{00}, u_{01}, u_{10}, u_{11}\right\}$ is typically called the Bell basis, and $\mu$ represents a measurement with respect to this basis.

Example 6.51 The previous example may be generalized as follows. Let $\Sigma=\mathbb{Z}_{n}$ for any positive integer $n$, let $\Gamma=\Sigma \times \Sigma$, and let the collection $\left\{U_{a b}: a, b \in \Sigma\right\} \subset \mathrm{U}\left(\mathbb{C}^{\Sigma}\right)$ of unitary operators be in correspondence with the discrete Weyl operators acting on $\mathbb{C}^{\Sigma}$. By taking $\mu: \Gamma \rightarrow \operatorname{Pos}\left(\mathbb{C}^{\Sigma} \otimes \mathbb{C}^{\Sigma}\right)$ to be the measurement defined as

$$
\begin{equation*}
\mu(a b)=\frac{\operatorname{vec}\left(U_{a b}\right) \operatorname{vec}\left(U_{a b}\right)^{*}}{n} \tag{6.332}
\end{equation*}
$$

for each $a, b \in \Sigma$, and setting $\Psi_{a b}(X)=U_{a b} X U_{a b}^{*}$ for each $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$, one again obtains a teleportation protocol as described in Scenario 6.46.

In the teleportation protocols described in the previous two examples, the number of distinct classical symbols that must be transmitted is equal to the square of the number of classical states in the quantum system that is teleported. This is optimal, as the following corollary states.

Corollary 6.52 Let $\Sigma$ and $\Gamma$ be alphabets, let $\mu: \Gamma \rightarrow \operatorname{Pos}\left(\mathbb{C}^{\Sigma} \otimes \mathbb{C}^{\Sigma}\right)$ be a measurement, and let $\left\{\Psi_{a}: a \in \Gamma\right\} \subseteq \mathrm{C}\left(\mathbb{C}^{\Sigma}\right)$ be a collection of channels such that

$$
\begin{equation*}
X=\frac{1}{|\Sigma|} \sum_{a \in \Gamma} \sum_{b, c \in \Sigma}\left\langle\mu(a), X \otimes E_{b, c}\right\rangle \Psi_{a}\left(E_{b, c}\right) \tag{6.333}
\end{equation*}
$$

for every $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$. It holds that $|\Gamma| \geq|\Sigma|^{2}$.
Proof By Theorem 6.47, it follows that

$$
\begin{equation*}
\mu(a)=p(a)|\Sigma| \operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*} \tag{6.334}
\end{equation*}
$$

for each $a \in \Gamma$, for some choice of a probability vector $p \in \mathcal{P}(\Gamma)$ and a
collection of unitary operators $\left\{U_{a}: a \in \Gamma\right\} \subset \mathrm{U}\left(\mathbb{C}^{\Sigma}\right)$. Each operator $\mu(a)$ has rank at most one, while

$$
\begin{equation*}
\sum_{a \in \Gamma} \mu(a)=\mathbb{1}_{\Sigma} \otimes \mathbb{1}_{\Sigma} \tag{6.335}
\end{equation*}
$$

has rank $|\Sigma|^{2}$. It follows that $|\Gamma| \geq|\Sigma|^{2}$ as required.

## Dense coding

Along similar lines to the discussion of teleportation above, a scenario in which Alice and Bob aim to implement an ideal classical channel through shared entanglement and quantum communication may be considered.

Scenario 6.53 (Dense coding) Alice holds a register X and Bob holds Y. Both registers have the same classical state set $\Sigma$, and the state of the pair $(\mathrm{X}, \mathrm{Y})$ is given by the maximally entangled state

$$
\begin{equation*}
\tau=\frac{1}{|\Sigma|} \sum_{b, c \in \Sigma} E_{b, c} \otimes E_{b, c} \tag{6.336}
\end{equation*}
$$

Alice obtains a classical register $Z$ having classical state set $\Gamma$. She wishes to transmit the classical state $Z$ to Bob by means of a protocol as follows:

1. Alice applies one of a collection of channels

$$
\begin{equation*}
\left\{\Psi_{a}: a \in \Gamma\right\} \subseteq \mathrm{C}(\mathcal{X}) \tag{6.337}
\end{equation*}
$$

to $\mathbf{X}$, with the channel applied being indexed by the classical state $a \in \Gamma$ of $Z$. The register $X$ is then sent to Bob.
2. Bob performs a measurement

$$
\begin{equation*}
\mu: \Gamma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y}) \tag{6.338}
\end{equation*}
$$

on the pair $(\mathbf{X}, \mathrm{Y})$. The outcome $b \in \Gamma$ is interpreted as the result of the transmission from Alice.

It is not surprising that protocols of this sort exist that function as desired, meaning that Bob's measurement outcome $b \in \Gamma$ corresponds precisely to the classical state $a \in \Gamma$ of Alice's register Z. Indeed, when $\Gamma$ is no larger than $\Sigma$, the task is trivially accomplished. What is more interesting is that there are protocols of this form that work perfectly in the case that $\Gamma$ is as large as $\Sigma \times \Sigma$.

The following proposition establishes that a dense coding protocol may be derived from an arbitrary mixed-unitary realization of the completely depolarizing channel, provided the unitary operators are drawn uniformly from a set indexed by $\Sigma \times \Sigma$.

Proposition 6.54 Let $\Sigma$ be an alphabet, let $\mathcal{X}=\mathbb{C}^{\Sigma}$, and let

$$
\begin{equation*}
\tau=\frac{1}{|\Sigma|} \sum_{c, d \in \Sigma} E_{c, d} \otimes E_{c, d} \tag{6.339}
\end{equation*}
$$

Assume $\left\{U_{a b}: a b \in \Sigma \times \Sigma\right\} \subset \mathrm{U}(\mathcal{X})$ is a collection of unitary operators such that

$$
\begin{equation*}
\Omega(X)=\frac{1}{|\Sigma|^{2}} \sum_{a b \in \Sigma \times \Sigma} U_{a b} X U_{a b}^{*} \tag{6.340}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X})$, where $\Omega \in \mathrm{C}(\mathcal{X})$ is the completely depolarizing channel with respect to the space $\mathcal{X}$. For $\left\{\Psi_{a b}: a b \in \Sigma \times \Sigma\right\} \subseteq \mathrm{C}(\mathcal{X})$ being a collection of channels defined as

$$
\begin{equation*}
\Psi_{a b}(X)=U_{a b} X U_{a b}^{*} \tag{6.341}
\end{equation*}
$$

for each $a b \in \Sigma \times \Sigma$ and $X \in \mathrm{~L}(\mathcal{X})$, and for $\mu: \Sigma \times \Sigma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{X})$ being defined as

$$
\begin{equation*}
\mu(a b)=\frac{\operatorname{vec}\left(U_{a b}\right) \operatorname{vec}\left(U_{a b}\right)^{*}}{|\Sigma|} \tag{6.342}
\end{equation*}
$$

for each $a b \in \Sigma \times \Sigma$, it holds that $\mu$ is a measurement and

$$
\left\langle\mu(c d),\left(\Psi_{a b} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)(\tau)\right\rangle= \begin{cases}1 & \text { if } a b=c d  \tag{6.343}\\ 0 & \text { if } a b \neq c d\end{cases}
$$

for all $a, b, c, d \in \Sigma$.
Proof It holds that

$$
\begin{equation*}
\sum_{a b \in \Sigma \times \Sigma} \mu(a b)=|\Sigma| J(\Omega)=\mathbb{1}_{\mathcal{X}} \otimes \mathbb{1}_{\mathcal{X}} \tag{6.344}
\end{equation*}
$$

As each operator $\mu(a b)$ is evidently positive semidefinite, it follows that $\mu$ is a measurement. For each $a b \in \Sigma \times \Sigma$, one has

$$
\begin{align*}
& \left\langle\mu(a b),\left(\Psi_{a b} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)(\tau)\right\rangle \\
& \quad=\frac{1}{|\Sigma|^{2}}\left\langle\operatorname{vec}\left(U_{a b}\right) \operatorname{vec}\left(U_{a b}\right)^{*}, \operatorname{vec}\left(U_{a b}\right) \operatorname{vec}\left(U_{a b}\right)^{*}\right\rangle=1 \tag{6.345}
\end{align*}
$$

Because $\left(\Psi_{a b} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)(\tau)$ is a density operator for each $a b \in \Sigma \times \Sigma$, it follows that

$$
\begin{equation*}
\left\langle\mu(c d),\left(\Psi_{a b} \otimes \mathbb{1}_{\mathrm{L}(\mathcal{X})}\right)(\tau)\right\rangle=0 \tag{6.346}
\end{equation*}
$$

for $c d \neq a b$, which completes the proof.

Example 6.55 As in Example 6.50, let $\Sigma=\{0,1\}$ and define unitary operators $U_{00}, U_{01}, U_{10}, U_{11} \in \mathrm{U}\left(\mathbb{C}^{\Sigma}\right)$ as follows:

$$
\begin{align*}
& U_{00}=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right), \quad U_{01}=\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right)  \tag{6.347}\\
& U_{10}=\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right), \quad U_{11}=\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right)
\end{align*}
$$

As the operators $U_{00}, U_{01}, U_{10}, U_{11}$ provide a mixed-unitary realization of the completely depolarizing channel, by taking $\mu: \Sigma \times \Sigma \rightarrow \operatorname{Pos}\left(\mathbb{C}^{\Sigma} \otimes \mathbb{C}^{\Sigma}\right)$ to be the measurement defined as

$$
\begin{align*}
& \mu(00)=\frac{\operatorname{vec}\left(U_{00}\right) \operatorname{vec}\left(U_{00}\right)^{*}}{2}, \quad \mu(01)=\frac{\operatorname{vec}\left(U_{01}\right) \operatorname{vec}\left(U_{01}\right)^{*}}{2}  \tag{6.348}\\
& \mu(10)=\frac{\operatorname{vec}\left(U_{10}\right) \operatorname{vec}\left(U_{10}\right)^{*}}{2}, \quad \mu(11)=\frac{\operatorname{vec}\left(U_{11}\right) \operatorname{vec}\left(U_{11}\right)^{*}}{2}
\end{align*}
$$

and setting $\Psi_{a b}(X)=U_{a b} X U_{a b}^{*}$ for each $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$, as in Example 6.50, one obtains a dense coding protocol as described in Scenario 6.53. The resulting protocol is equivalent to the traditional notion of dense coding in which an ideal two-bit classical channel is implemented using a maximally entangled pair of qubits along with one qubit of quantum communication.

In analogy to the more general type of teleportation protocol described previously, one may consider the capabilities of dense coding protocols for arbitrary choices of an alphabet $\Gamma$, as opposed to $\Gamma=\Sigma \times \Sigma$. In particular, suppose Alice's channels are given by the collection

$$
\begin{equation*}
\left\{\Psi_{a}: a \in \Gamma\right\} \subseteq \mathrm{C}(\mathcal{X}) \tag{6.349}
\end{equation*}
$$

for an arbitrary alphabet $\Gamma$, and that the symbol $a \in \Gamma$ Alice wishes to send to Bob is randomly selected according to a probability vector $p \in \mathcal{P}(\Gamma)$. The state of the pair $(\mathrm{X}, \mathrm{Y})$ prior to Bob's measurement is described by the ensemble $\eta: \Gamma \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{X})$ defined as

$$
\begin{equation*}
\eta(a)=\frac{p(a)}{|\Sigma|} \sum_{b, c \in \Sigma} \Psi_{a}\left(E_{b, c}\right) \otimes E_{b, c} \tag{6.350}
\end{equation*}
$$

for all $a \in \Gamma$. The following theorem provides a characterization of when the Holevo information $\chi(\eta)$ of this ensemble attains its maximum possible value, which is $2 \log (|\Sigma|)$.

Theorem 6.56 Let $\Sigma$ and $\Gamma$ be alphabets, let $p \in \mathcal{P}(\Gamma)$ be a probability vector such that $p(a) \neq 0$ for all $a \in \Gamma$, and let

$$
\begin{equation*}
\left\{\Psi_{a}: a \in \Gamma\right\} \subseteq \mathrm{C}\left(\mathbb{C}^{\Sigma}\right) \tag{6.351}
\end{equation*}
$$

be a collection of channels. The following two statements are equivalent:

1. For the ensemble $\eta: \Gamma \rightarrow \operatorname{Pos}\left(\mathbb{C}^{\Sigma} \otimes \mathbb{C}^{\Sigma}\right)$ defined as

$$
\begin{equation*}
\eta(a)=\frac{p(a)}{|\Sigma|} \sum_{b, c \in \Sigma} \Psi_{a}\left(E_{b, c}\right) \otimes E_{b, c} \tag{6.352}
\end{equation*}
$$

for all $a \in \Gamma$, one has that $\chi(\eta)=2 \log (|\Sigma|)$.
2. There exists a collection $\left\{U_{a}: a \in \Gamma\right\} \subset \mathrm{U}\left(\mathbb{C}^{\Sigma}\right)$ of unitary operators such that

$$
\begin{equation*}
\Psi_{a}(X)=U_{a} X U_{a}^{*} \tag{6.353}
\end{equation*}
$$

for every choice of $a \in \Gamma$ and $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$, and moreover it holds that

$$
\begin{equation*}
\Omega(X)=\sum_{a \in \Gamma} p(a) U_{a} X U_{a}^{*} \tag{6.354}
\end{equation*}
$$

for all $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$, where $\Omega \in \mathrm{C}\left(\mathbb{C}^{\Sigma}\right)$ denotes the completely depolarizing channel defined with respect to the space $\mathbb{C}^{\Sigma}$.

Proof The Holevo information of the ensemble $\eta$ defined by (6.352) is

$$
\begin{align*}
\chi(\eta)=\mathrm{H}\left(\sum_{a \in \Gamma}\right. & \left.\frac{p(a)}{|\Sigma|} \sum_{b, c \in \Sigma} \Psi_{a}\left(E_{b, c}\right) \otimes E_{b, c}\right)  \tag{6.355}\\
& -\sum_{a \in \Gamma} p(a) \mathrm{H}\left(\frac{1}{|\Sigma|} \sum_{b, c \in \Sigma} \Psi_{a}\left(E_{b, c}\right) \otimes E_{b, c}\right),
\end{align*}
$$

which may alternatively be written as

$$
\begin{equation*}
\chi(\eta)=\mathrm{H}\left(\sum_{a \in \Gamma} p(a) \frac{J\left(\Psi_{a}\right)}{|\Sigma|}\right)-\sum_{a \in \Gamma} p(a) \mathrm{H}\left(\frac{J\left(\Psi_{a}\right)}{|\Sigma|}\right) . \tag{6.356}
\end{equation*}
$$

Under the assumption that $\chi(\eta)=2 \log (|\Sigma|)$, it must hold that

$$
\begin{equation*}
\mathrm{H}\left(\sum_{a \in \Gamma} p(a) \frac{J\left(\Psi_{a}\right)}{|\Sigma|}\right)=2 \log (|\Sigma|) \quad \text { and } \quad \mathrm{H}\left(\frac{J\left(\Psi_{a}\right)}{|\Sigma|}\right)=0 \tag{6.357}
\end{equation*}
$$

for each $a \in \Gamma$. The rank of $J\left(\Psi_{a}\right)$ is therefore equal to 1 for each $a \in \Gamma$, and as each $\Psi_{a}$ is a channel it follows that there must exist a collection of unitary operators

$$
\begin{equation*}
\left\{U_{a}: a \in \Gamma\right\} \subset \mathrm{U}\left(\mathbb{C}^{\Sigma}\right) \tag{6.358}
\end{equation*}
$$

such that (6.353) holds for each $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$ and each $a \in \Gamma$. The first equation of (6.357) is equivalent to

$$
\begin{equation*}
\sum_{a \in \Gamma} p(a) \frac{J\left(\Psi_{a}\right)}{|\Sigma|}=\frac{\mathbb{1} \otimes \mathbb{1}}{|\Sigma|^{2}} \tag{6.359}
\end{equation*}
$$

which implies

$$
\begin{equation*}
\sum_{a \in \Gamma} p(a) \operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*}=\frac{\mathbb{1} \otimes \mathbb{1}}{|\Sigma|}=J(\Omega) \tag{6.360}
\end{equation*}
$$

and therefore

$$
\begin{equation*}
\sum_{a \in \Gamma} p(a) U_{a} X U_{a}^{*}=\Omega(X) \tag{6.361}
\end{equation*}
$$

for all $X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma}\right)$. Statement 1 therefore implies statement 2 .
Under the assumption that statement 2 holds, the Holevo information of $\eta$ may be calculated directly:

$$
\begin{aligned}
\chi(\eta)= & \mathrm{H}\left(\sum_{a \in \Gamma} \frac{p(a)}{|\Sigma|} \sum_{b, c \in \Sigma} \Psi_{a}\left(E_{b, c}\right) \otimes E_{b, c}\right) \\
& -\sum_{a \in \Gamma} p(a) \mathrm{H}\left(\frac{1}{|\Sigma|} \sum_{b, c \in \Sigma} \Psi_{a}\left(E_{b, c}\right) \otimes E_{b, c}\right) \\
= & \mathrm{H}\left(\frac{\mathbb{1} \otimes \mathbb{1}}{|\Sigma|^{2}}\right)-\sum_{a \in \Gamma} p(a) \mathrm{H}\left(\frac{\operatorname{vec}\left(U_{a}\right) \operatorname{vec}\left(U_{a}\right)^{*}}{|\Sigma|}\right) \\
= & 2 \log (|\Sigma|) .
\end{aligned}
$$

Statement 2 therefore implies statement 1, which completes the proof.

### 6.3.2 Non-classical correlations

The definition of entanglement, as the absence of separability, is not directly related to an observable physical phenomenon. Entanglement is, however, fundamentally connected with the correlations that may exist among the outcomes of measurements performed on two or more separate parts of a physical system. To describe this connection, it is helpful to consider the following scenario.

Scenario 6.57 Two individuals, Alice and Bob, share a compound register ( $\mathrm{X}, \mathrm{Y}$ ), with Alice holding X and Bob holding Y . Two events occur:

1. Alice receives an input symbol, drawn from a fixed alphabet $\Sigma_{\mathrm{A}}$, and she must produce an output symbol from a fixed alphabet $\Gamma_{\mathrm{A}}$.
2. Bob receives an input symbol, drawn from a fixed alphabet $\Sigma_{\mathrm{B}}$, and he must produce an output symbol from a fixed alphabet $\Gamma_{\mathrm{B}}$.

Alice and Bob cannot communicate with one another at any point after they have received their input symbols. The output symbols they produce may, in general, be probabilistic, possibly resulting from measurements made on whichever one of the registers X or Y is in the possession of the individual performing the measurement.

The discussion that follows is primarily concerned with the collections of output distributions that may be produced by Alice and Bob, as described in the scenario above, through measurements on a shared entangled state, as compared with the correlations that may result from the initial state of $(\mathrm{X}, \mathrm{Y})$ being separable.

## Correlation operators

The output distributions produced by Alice and Bob in a particular instance of Scenario 6.57 , ranging over all pairs of input symbols, may collectively be described by a single operator

$$
\begin{equation*}
C \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.363}
\end{equation*}
$$

defined so that $C((a, c),(b, d))$ is the probability that Alice and Bob output $(c, d) \in \Gamma_{\mathrm{A}} \times \Gamma_{\mathrm{B}}$, assuming they are given the input pair $(a, b) \in \Sigma_{\mathrm{A}} \times \Sigma_{\mathrm{B}}$. Such an operator must satisfy certain constraints. For instance, to carry the interpretation that $C$ represents a collection of probability distributions, each entry must be a nonnegative real number, and it must hold that

$$
\begin{equation*}
\sum_{(c, d) \in \Gamma_{\mathrm{A}} \times \Gamma_{\mathrm{B}}} C((a, c),(b, d))=1 \tag{6.364}
\end{equation*}
$$

for every pair $(a, b) \in \Sigma_{\mathrm{A}} \times \Sigma_{\mathrm{B}}$. Additional constraints are imposed by the assumption that Alice and Bob are separated and cannot communicate.

Definition 6.58 Let $\Sigma_{A}, \Sigma_{\mathrm{B}}, \Gamma_{\mathrm{A}}$, and $\Gamma_{\mathrm{B}}$ be alphabets, and let

$$
\begin{equation*}
C \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.365}
\end{equation*}
$$

be an operator.

1. The operator $C$ is a deterministic correlation operator if

$$
\begin{equation*}
C=\sum_{(a, b) \in \Sigma_{\mathrm{A}} \times \Sigma_{\mathrm{B}}} E_{a, b} \otimes E_{f(a), g(b)}, \tag{6.366}
\end{equation*}
$$

or equivalently

$$
C((a, c),(b, d))= \begin{cases}1 & \text { if } c=f(a) \text { and } d=g(b)  \tag{6.367}\\ 0 & \text { otherwise }\end{cases}
$$

for some choice of functions $f: \Sigma_{\mathrm{A}} \rightarrow \Gamma_{\mathrm{A}}$ and $g: \Sigma_{\mathrm{B}} \rightarrow \Gamma_{\mathrm{B}}$. It is said that $C$ is a probabilistic correlation operator if $C$ is equal to a convex combination of deterministic correlation operators.
2. The operator $C$ is a quantum correlation operator if there exist complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, a state $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$, and two collections of measurements $\left\{\mu_{a}: a \in \Sigma_{\mathrm{A}}\right\}$ and $\left\{\nu_{b}: b \in \Sigma_{\mathrm{B}}\right\}$, taking the form

$$
\begin{equation*}
\mu_{a}: \Gamma_{\mathrm{A}} \rightarrow \operatorname{Pos}(\mathcal{X}) \quad \text { and } \quad \nu_{b}: \Gamma_{\mathrm{B}} \rightarrow \operatorname{Pos}(\mathcal{Y}) \tag{6.368}
\end{equation*}
$$

such that

$$
\begin{equation*}
C((a, c),(b, d))=\left\langle\mu_{a}(c) \otimes \nu_{b}(d), \rho\right\rangle \tag{6.369}
\end{equation*}
$$

for every $a \in \Sigma_{\mathrm{A}}, b \in \Sigma_{\mathrm{B}}, c \in \Gamma_{\mathrm{A}}$, and $d \in \Gamma_{\mathrm{B}}$.
Example 6.59 Let $\Sigma_{\mathrm{A}}, \Sigma_{\mathrm{B}}, \Gamma_{\mathrm{A}}$, and $\Gamma_{\mathrm{B}}$ all be equal to the binary alphabet $\Sigma=\{0,1\}$, let $\mathcal{X}=\mathbb{C}^{\Sigma}$ and $\mathcal{Y}=\mathbb{C}^{\Sigma}$, define $\tau \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ to be the maximally entangled state

$$
\begin{equation*}
\tau=\frac{1}{2} \sum_{a, b \in \Sigma} E_{a, b} \otimes E_{a, b} \tag{6.370}
\end{equation*}
$$

and define measurements $\mu_{0}, \mu_{1}: \Gamma_{\mathrm{A}} \rightarrow \operatorname{Pos}(\mathcal{X})$ and $\nu_{0}, \nu_{1}: \Gamma_{\mathrm{B}} \rightarrow \operatorname{Pos}(\mathcal{Y})$ as

$$
\begin{array}{ll}
\mu_{0}(0)=\Pi_{0}, & \mu_{0}(1)=\Pi_{\pi / 2} \\
\mu_{1}(0)=\Pi_{\pi / 4}, & \mu_{1}(1)=\Pi_{3 \pi / 4}  \tag{6.371}\\
\nu_{0}(0)=\Pi_{\pi / 8}, & \nu_{0}(1)=\Pi_{5 \pi / 8} \\
\nu_{1}(0)=\Pi_{7 \pi / 8}, & \nu_{1}(1)=\Pi_{3 \pi / 8}
\end{array}
$$

for

$$
\Pi_{\theta}=\left(\begin{array}{cc}
\cos ^{2}(\theta) & \cos (\theta) \sin (\theta)  \tag{6.372}\\
\cos (\theta) \sin (\theta) & \sin ^{2}(\theta)
\end{array}\right)
$$

Equivalently, these measurement operators are as described in Figure 6.1.
For this choice of $\tau$, and because each of the measurement operators above have real number entries, it holds that

$$
\begin{equation*}
\left\langle\mu_{a}(c) \otimes \nu_{b}(d), \tau\right\rangle=\frac{1}{2}\left\langle\mu_{a}(c), \nu_{b}(d)\right\rangle \tag{6.373}
\end{equation*}
$$

$$
\begin{array}{ll}
\mu_{0}(0)=\left(\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right), & \mu_{0}(1)=\left(\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right), \\
\mu_{1}(0)=\left(\begin{array}{cc}
\frac{1}{2} & \frac{1}{2} \\
\frac{1}{2} & \frac{1}{2}
\end{array}\right), & \mu_{1}(1)=\left(\begin{array}{cc}
\frac{1}{2} & -\frac{1}{2} \\
-\frac{1}{2} & \frac{1}{2}
\end{array}\right), \\
\nu_{0}(0)=\left(\begin{array}{cc}
\frac{2+\sqrt{2}}{4} & \frac{\sqrt{2}}{4} \\
\frac{\sqrt{2}}{4} & \frac{2-\sqrt{2}}{4}
\end{array}\right), & \nu_{0}(1)=\left(\begin{array}{cc}
\frac{2-\sqrt{2}}{4} & -\frac{\sqrt{2}}{4} \\
-\frac{\sqrt{2}}{4} & \frac{2+\sqrt{2}}{4}
\end{array}\right), \\
\nu_{1}(0)=\left(\begin{array}{cc}
\frac{2+\sqrt{2}}{4} & -\frac{\sqrt{2}}{4} \\
-\frac{\sqrt{2}}{4} & \frac{2-\sqrt{2}}{4}
\end{array}\right), & \nu_{1}(1)=\left(\begin{array}{cc}
\frac{2-\sqrt{2}}{4} & \frac{\sqrt{2}}{4} \\
\frac{\sqrt{2}}{4} & \frac{2+\sqrt{2}}{4}
\end{array}\right) .
\end{array}
$$

Figure 6.1 Matrix representations of the measurement operators described in Example 6.59.
for each $a \in \Sigma_{\mathrm{A}}, b \in \Sigma_{\mathrm{B}}, c \in \Gamma_{\mathrm{A}}$, and $d \in \Gamma_{\mathrm{B}}$. A calculation reveals that the quantum correlation operator defined by (6.369) is given by

$$
C=\left(\begin{array}{llll}
\frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8}  \tag{6.374}\\
\frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} \\
\frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} \\
\frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8}
\end{array}\right) .
$$

It will be demonstrated shortly that the operator $C$ is not a probabilistic correlation operator.

Example 6.60 Let $\Sigma_{\mathrm{A}}, \Sigma_{\mathrm{B}}, \Gamma_{\mathrm{A}}$, and $\Gamma_{\mathrm{B}}$ all be equal to the binary alphabet $\Sigma=\{0,1\}$. There are 16 deterministic correlation operators, which are in correspondence with the 16 possible pairs of functions $(f, g)$ having the form $f: \Sigma_{\mathrm{A}} \rightarrow \Gamma_{\mathrm{A}}$ and $g: \Sigma_{\mathrm{B}} \rightarrow \Gamma_{\mathrm{B}}$. As matrices, these operators are as described in Figure 6.2.

## Bell inequalities

By its definition, the set of all probabilistic correlation operators of the form

$$
\begin{equation*}
C \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.375}
\end{equation*}
$$

is convex. Indeed, this set is given by the convex hull of a finite set, as there are finitely many deterministic correlation operators of the same form. From this fact it follows that the set of all probabilistic correlation operators of the form (6.375) is compact. Therefore, by the separating hyperplane theorem

$$
\begin{aligned}
& \left(\begin{array}{llll}
1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 \\
1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
1 & 0 & 0 & 1 \\
0 & 0 & 0 & 0 \\
1 & 0 & 0 & 1 \\
0 & 0 & 0 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 1 & 1 & 0 \\
0 & 0 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 1 & 0 & 1 \\
0 & 0 & 0 & 0 \\
0 & 1 & 0 & 1 \\
0 & 0 & 0 & 0
\end{array}\right), \\
& \left(\begin{array}{llll}
1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
1 & 0 & 1 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
1 & 0 & 0 & 1 \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
1 & 0 & 0 & 1
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 1 & 1 & 0 \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
0 & 1 & 1 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 1 & 0 & 1 \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
0 & 1 & 0 & 1
\end{array}\right), \\
& \left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
1 & 0 & 1 & 0 \\
1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
1 & 0 & 0 & 1 \\
1 & 0 & 0 & 1 \\
0 & 0 & 0 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
0 & 1 & 0 & 1 \\
0 & 1 & 0 & 1 \\
0 & 0 & 0 & 0
\end{array}\right), \\
& \left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 \\
1 & 0 & 1 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
1 & 0 & 0 & 1 \\
0 & 0 & 0 & 0 \\
1 & 0 & 0 & 1
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 0 & 0 \\
0 & 1 & 1 & 0
\end{array}\right), \quad\left(\begin{array}{llll}
0 & 0 & 0 & 0 \\
0 & 1 & 0 & 1 \\
0 & 0 & 0 & 0 \\
0 & 1 & 0 & 1
\end{array}\right) .
\end{aligned}
$$

Figure 6.2 Matrix representations of the correlation operators described in Example 6.60.
(Theorem 1.11), if an operator

$$
\begin{equation*}
D \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.376}
\end{equation*}
$$

is not a probabilistic correlation operator, there must exist an operator

$$
\begin{equation*}
K \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.377}
\end{equation*}
$$

and a real number $\alpha$ such that

$$
\begin{equation*}
\langle K, D\rangle>\alpha \quad \text { and } \quad\langle K, C\rangle \leq \alpha \tag{6.378}
\end{equation*}
$$

for all probabilistic correlation operators $C$ of the form (6.375).
For a fixed choice of an operator $K$ and a real number $\alpha$, the inequality $\langle K, C\rangle \leq \alpha$ is traditionally called a Bell inequality, assuming it is satisfied for every probabilistic correlation operator $C$ of the form (6.375). When this is the case, the inequality $\langle K, D\rangle>\alpha$ is called a Bell inequality violation if it holds for some choice of a quantum correlation operator $D$.

The illustration of a Bell inequality violation can provide a convenient way to demonstrate that certain correlation operators are not probabilistic, as the following example illustrates.

Example 6.61 (Clauser-Horn-Shimony-Holt inequality) Let $\Sigma_{\mathrm{A}}, \Sigma_{\mathrm{B}}, \Gamma_{\mathrm{A}}$, and $\Gamma_{\mathrm{B}}$ all be equal to the binary alphabet $\Sigma=\{0,1\}$, and define

$$
\begin{equation*}
K \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.379}
\end{equation*}
$$

as

$$
K=\left(\begin{array}{cccc}
1 & -1 & 1 & -1  \tag{6.380}\\
-1 & 1 & -1 & 1 \\
1 & -1 & -1 & 1 \\
-1 & 1 & 1 & -1
\end{array}\right)
$$

For every deterministic correlation operator

$$
\begin{equation*}
C \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.381}
\end{equation*}
$$

it holds that

$$
\begin{equation*}
\langle K, C\rangle \leq 2 \tag{6.382}
\end{equation*}
$$

which may be verified by an inspection of the 16 deterministic correlation operators in Example 6.60. It follows by convexity that the same inequality holds for $C$ being any probabilistic correlation operator. On the other hand, the quantum correlation operator

$$
D=\left(\begin{array}{cccc}
\frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8}  \tag{6.383}\\
\frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} \\
\frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} \\
\frac{2-\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2+\sqrt{2}}{8} & \frac{2-\sqrt{2}}{8}
\end{array}\right)
$$

described in Example 6.59 satisfies

$$
\begin{equation*}
\langle K, D\rangle=2 \sqrt{2} \tag{6.384}
\end{equation*}
$$

This demonstrates that $D$ is not a probabilistic correlation operator.
Correlations among binary-valued measurements
For a given choice of alphabets $\Sigma_{\mathrm{A}}, \Sigma_{\mathrm{B}}, \Gamma_{\mathrm{A}}$, and $\Gamma_{\mathrm{B}}$, and an operator

$$
\begin{equation*}
K \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.385}
\end{equation*}
$$

it may be quite difficult in some cases to determine the supremum value of $\langle K, C\rangle$, optimized over all quantum correlation operators of the form

$$
\begin{equation*}
C \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.386}
\end{equation*}
$$

There is, however, an interesting class of operators $K$ for which this problem is solvable. This is the class for which the output alphabets $\Gamma_{A}$ and $\Gamma_{B}$ are both equal to the binary alphabet $\Sigma=\{0,1\}$, and furthermore the operator $K$ takes the form

$$
K=M \otimes\left(\begin{array}{cc}
1 & -1  \tag{6.387}\\
-1 & 1
\end{array}\right)
$$

for some choice of an operator

$$
\begin{equation*}
M \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}}}\right) \tag{6.388}
\end{equation*}
$$

Operators of the form (6.387) have a simple interpretation when considered in the context of Bell inequalities and violations-they effectively assign the value $M(a, b)$ to the event that Alice and Bob output equal binary-valued answers, and the value $-M(a, b)$ to the event that their outputs differ, for each possible input pair $(a, b)$.

The following theorem, known as Tsirelson's theorem, provides the basis for a solution to the problem under consideration.

Theorem 6.62 (Tsirelson's theorem) Let $\Sigma_{\mathrm{A}}$ and $\Sigma_{\mathrm{B}}$ be alphabets and let $X \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}}}\right)$ be an operator. The following statements are equivalent:

1. There exist complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, a state $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$, and two collections

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma_{\mathrm{A}}\right\} \subset \operatorname{Herm}(\mathcal{X}) \quad \text { and } \quad\left\{B_{b}: b \in \Sigma_{\mathrm{B}}\right\} \subset \operatorname{Herm}(\mathcal{Y}) \tag{6.389}
\end{equation*}
$$

of operators satisfying $\left\|A_{a}\right\| \leq 1,\left\|B_{b}\right\| \leq 1$, and

$$
\begin{equation*}
X(a, b)=\left\langle A_{a} \otimes B_{b}, \rho\right\rangle \tag{6.390}
\end{equation*}
$$

for every $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$.
2. Statement 1 holds under the additional requirement that, for some choice of an alphabet $\Gamma$, one has $\mathcal{X}=\mathbb{C}^{\Gamma}, \mathcal{Y}=\mathbb{C}^{\Gamma}$, and

$$
\begin{equation*}
\rho=\frac{1}{|\Gamma|} \sum_{c, d \in \Gamma} E_{c, d} \otimes E_{c, d} \tag{6.391}
\end{equation*}
$$

and furthermore that the operators in the collections

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma_{\mathrm{A}}\right\} \quad \text { and } \quad\left\{B_{b}: b \in \Sigma_{\mathrm{B}}\right\} \tag{6.392}
\end{equation*}
$$

are unitary (in addition to being Hermitian).
3. There exist operators

$$
\begin{equation*}
P \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{A}}}\right) \quad \text { and } \quad Q \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{B}}}\right) \tag{6.393}
\end{equation*}
$$

with $P(a, a)=1$ and $Q(b, b)=1$ for every $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$, such that

$$
\left(\begin{array}{cc}
P & X  \tag{6.394}\\
X^{*} & Q
\end{array}\right) \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{A}}} \oplus \mathbb{C}^{\Sigma_{\mathrm{B}}}\right)
$$

4. There exist two collections

$$
\begin{equation*}
\left\{u_{a}: a \in \Sigma_{\mathrm{A}}\right\},\left\{v_{b}: b \in \Sigma_{\mathrm{B}}\right\} \subset \mathbb{R}^{\Sigma_{\mathrm{A}}} \oplus \mathbb{R}^{\Sigma_{\mathrm{B}}} \tag{6.395}
\end{equation*}
$$

of unit vectors such that

$$
\begin{equation*}
X(a, b)=\left\langle u_{a}, v_{b}\right\rangle \tag{6.396}
\end{equation*}
$$

for every $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$.
The proof of this theorem will make use of a collection of unitary and Hermitian operators known as Weyl-Brauer operators.

Definition 6.63 Let $m$ be a positive integer, let $\Gamma=\{0,1\}$, and let $\mathcal{Z}=\mathbb{C}^{\Gamma}$. The Weyl-Brauer operators

$$
\begin{equation*}
V_{0}, \ldots, V_{2 m} \in \mathrm{~L}\left(\mathcal{Z}^{\otimes m}\right) \tag{6.397}
\end{equation*}
$$

of order $m$ are defined as follows: $V_{0}=\sigma_{z}^{\otimes m}$ and

$$
\begin{align*}
V_{2 k-1} & =\sigma_{z}^{\otimes(k-1)} \otimes \sigma_{x} \otimes \mathbb{1}^{\otimes(m-k)} \\
V_{2 k} & =\sigma_{z}^{\otimes(k-1)} \otimes \sigma_{y} \otimes \mathbb{1}^{\otimes(m-k)} \tag{6.398}
\end{align*}
$$

for $k=1, \ldots, m$, where $\mathbb{1}$ denotes the identity operator on $\mathcal{Z}$ and $\sigma_{x}, \sigma_{y}$, and $\sigma_{z}$ are given by the Pauli operators. In matrix form, these operators are as follows:

$$
\mathbb{1}=\left(\begin{array}{ll}
1 & 0  \tag{6.399}\\
0 & 1
\end{array}\right), \quad \sigma_{x}=\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right), \quad \sigma_{y}=\left(\begin{array}{cc}
0 & -i \\
i & 0
\end{array}\right), \quad \sigma_{z}=\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right) .
$$

Example 6.64 In the case $m=3$, the Weyl-Brauer operators $V_{0}, \ldots, V_{6}$ are

$$
\begin{align*}
& V_{0}=\sigma_{z} \otimes \sigma_{z} \otimes \sigma_{z} \\
& V_{1}=\sigma_{x} \otimes \mathbb{1} \otimes \mathbb{1} \\
& V_{2}=\sigma_{y} \otimes \mathbb{1} \otimes \mathbb{1} \\
& V_{3}=\sigma_{z} \otimes \sigma_{x} \otimes \mathbb{1}  \tag{6.400}\\
& V_{4}=\sigma_{z} \otimes \sigma_{y} \otimes \mathbb{1} \\
& V_{5}=\sigma_{z} \otimes \sigma_{z} \otimes \sigma_{x} \\
& V_{6}=\sigma_{z} \otimes \sigma_{z} \otimes \sigma_{y} .
\end{align*}
$$

A proposition summarizing the properties of the Weyl-Brauer operators that are relevant to the proof of Tsirelson's theorem follows.

Proposition 6.65 Let $m$ be a positive integer, let $V_{0}, \ldots, V_{2 m}$ denote the Weyl-Brauer operators of order $m$, and let

$$
\begin{equation*}
\left(\alpha_{0}, \ldots, \alpha_{2 m}\right),\left(\beta_{0}, \ldots, \beta_{2 m}\right) \in \mathbb{R}^{2 m+1} \tag{6.401}
\end{equation*}
$$

be vectors of real numbers. It holds that

$$
\begin{equation*}
\left(\sum_{k=0}^{2 m} \alpha_{k} V_{k}\right)^{2}=\left(\sum_{k=0}^{2 m} \alpha_{k}^{2}\right) \mathbb{1}^{\otimes m} \tag{6.402}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{2^{m}}\left\langle\sum_{j=0}^{2 m} \alpha_{j} V_{j}, \sum_{k=0}^{2 m} \beta_{k} V_{k}\right\rangle=\sum_{k=0}^{2 m} \alpha_{k} \beta_{k} \tag{6.403}
\end{equation*}
$$

Proof The Pauli operators anti-commute in pairs:

$$
\begin{equation*}
\sigma_{x} \sigma_{y}=-\sigma_{y} \sigma_{x}, \quad \sigma_{x} \sigma_{z}=-\sigma_{z} \sigma_{x}, \quad \text { and } \quad \sigma_{y} \sigma_{z}=-\sigma_{z} \sigma_{y} \tag{6.404}
\end{equation*}
$$

By an inspection of the definition of the Weyl-Brauer operators, it follows that $V_{0}, \ldots, V_{2 m}$ also anti-commute in pairs:

$$
\begin{equation*}
V_{j} V_{k}=-V_{k} V_{j} \tag{6.405}
\end{equation*}
$$

for distinct choices of $j, k \in\{0, \ldots, 2 m\}$. Moreover, each $V_{k}$ is both unitary and Hermitian, and therefore $V_{k}^{2}=\mathbb{1}^{\otimes m}$. It follows that

$$
\begin{align*}
\left(\sum_{k=0}^{2 m} \alpha_{k} V_{k}\right)^{2} & =\sum_{k=0}^{2 m} \alpha_{k}^{2} V_{k}^{2}+\sum_{0 \leq j<k \leq 2 m} \alpha_{j} \alpha_{k}\left(V_{j} V_{k}+V_{k} V_{j}\right)  \tag{6.406}\\
& =\left(\sum_{k=0}^{2 m} \alpha_{k}^{2}\right) \mathbb{1}^{\otimes m}
\end{align*}
$$

Moreover,

$$
\left\langle V_{j}, V_{k}\right\rangle= \begin{cases}2^{m} & \text { if } j=k  \tag{6.407}\\ 0 & \text { if } j \neq k\end{cases}
$$

and therefore

$$
\begin{equation*}
\frac{1}{2^{m}}\left\langle\sum_{j=0}^{2 m} \alpha_{j} V_{j}, \sum_{k=0}^{2 m} \beta_{k} V_{k}\right\rangle=\frac{1}{2^{m}} \sum_{j=0}^{2 m} \sum_{k=0}^{2 m} \alpha_{j} \beta_{k}\left\langle V_{j}, V_{k}\right\rangle=\sum_{k=0}^{2 m} \alpha_{k} \beta_{k} \tag{6.408}
\end{equation*}
$$

as required.

Proof of Theorem 6.62 The following implications among the statements will suffice to prove the theorem:

$$
\begin{equation*}
(2) \Rightarrow(1) \Rightarrow(3) \Rightarrow(4) \Rightarrow(2) \tag{6.409}
\end{equation*}
$$

The first implication, that statement 2 implies statement 1 , is trivial.
Assume statement 1 holds, define an operator

$$
\begin{equation*}
K=\sum_{a \in \Sigma_{\mathrm{A}}} e_{a} \operatorname{vec}\left(\left(A_{a} \otimes \mathbb{1}\right) \sqrt{\rho}\right)^{*}+\sum_{b \in \Sigma_{\mathrm{B}}} e_{b} \operatorname{vec}\left(\left(\mathbb{1} \otimes B_{b}\right) \sqrt{\rho}\right)^{*} \tag{6.410}
\end{equation*}
$$

and consider the operator $K K^{*} \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{A}} \sqcup \Sigma_{\mathrm{B}}}\right)$, which may be written in a block form as

$$
K K^{*}=\left(\begin{array}{cc}
P & Y  \tag{6.411}\\
Y^{*} & Q
\end{array}\right)
$$

for $P \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{A}}}\right), Q \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{B}}}\right)$, and $Y \in \mathrm{~L}\left(\mathbb{C}^{\Sigma_{\mathrm{B}}}, \mathbb{C}^{\Sigma_{\mathrm{A}}}\right)$. It holds that

$$
\begin{equation*}
Y(a, b)=\left\langle\left(A_{a} \otimes \mathbb{1}\right) \sqrt{\rho},\left(\mathbb{1} \otimes B_{b}\right) \sqrt{\rho}\right\rangle=\left\langle A_{a} \otimes B_{b}, \rho\right\rangle=X(a, b) \tag{6.412}
\end{equation*}
$$

for every $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$, and therefore $Y=X$. Moreover, for each $a \in \Sigma_{\mathrm{A}}$ one has

$$
\begin{equation*}
P(a, a)=\left\langle\left(A_{a} \otimes \mathbb{1}\right) \sqrt{\rho},\left(A_{a} \otimes \mathbb{1}\right) \sqrt{\rho}\right\rangle=\left\langle A_{a}^{2} \otimes \mathbb{1}, \rho\right\rangle, \tag{6.413}
\end{equation*}
$$

which is necessarily a nonnegative real number in the interval $[0,1]$; and through a similar calculation, one finds that $Q(b, b)$ is also a nonnegative integer in the interval $[0,1]$ for each $b \in \Sigma_{\mathrm{B}}$. A nonnegative real number may be added to each diagonal entry of this operator to yield another positive semidefinite operator, so one has that statement 3 holds. It has therefore been proved that statement 1 implies statement 3 .

Next, assume statement 3 holds, and observe that

$$
\frac{1}{2}\left(\begin{array}{cc}
P & X  \tag{6.414}\\
X^{*} & Q
\end{array}\right)+\frac{1}{2}\left(\begin{array}{cc}
P & X \\
X^{*} & Q
\end{array}\right)^{\top}=\left(\begin{array}{cc}
\frac{P+\bar{P}}{2} & X \\
X^{*} & \frac{Q+\bar{Q}}{2}
\end{array}\right)
$$

is a positive semidefinite operator having real number entries, and all of its diagonal entries are equal to 1 . Define

$$
u_{a}=\left(\begin{array}{cc}
\frac{P+\bar{P}}{2} & X  \tag{6.415}\\
X^{*} & \frac{Q+\bar{Q}}{2}
\end{array}\right)^{\frac{1}{2}}\binom{e_{a}}{0} \quad \text { and } \quad v_{b}=\left(\begin{array}{cc}
\frac{P+\bar{P}}{2} & X \\
X^{*} & \frac{Q+\bar{Q}}{2}
\end{array}\right)^{\frac{1}{2}}\binom{0}{e_{b}}
$$

for each $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$. As the square root of a positive semidefinite operator having real number entries also has real number entries, one has
that $u_{a}$ and $v_{b}$ are unit vectors with real number entries, and moreover it holds that

$$
\begin{equation*}
\left\langle u_{a}, v_{b}\right\rangle=X(a, b) \tag{6.416}
\end{equation*}
$$

for all $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$. It has therefore been proved that statement 3 implies statement 4.

Finally, assume statement 4 holds. Let

$$
\begin{equation*}
m=\left\lceil\frac{\left|\Sigma_{\mathrm{A}}\right|+\left|\Sigma_{\mathrm{B}}\right|-1}{2}\right\rceil, \tag{6.417}
\end{equation*}
$$

so that $2 m+1 \geq\left|\Sigma_{\mathrm{A}}\right|+\left|\Sigma_{\mathrm{B}}\right|$, and let $f: \Sigma_{\mathrm{A}} \sqcup \Sigma_{\mathrm{B}} \rightarrow\{0, \ldots, 2 m\}$ be a fixed but otherwise arbitrarily chosen injective function. Let $\Gamma=\{0,1\}$, let $\mathcal{Z}=\mathbb{C}^{\Gamma}$, and define

$$
\begin{equation*}
A_{a}=\sum_{c \in \Sigma_{\mathrm{A}} \sqcup \Sigma_{\mathrm{B}}} u_{a}(c) V_{f(c)} \quad \text { and } \quad B_{b}=\sum_{c \in \Sigma_{\mathrm{A}} \sqcup \Sigma_{\mathrm{B}}} v_{b}(c) V_{f(c)}^{\top} \tag{6.418}
\end{equation*}
$$

for each $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$, for $V_{0}, \ldots, V_{2 m}$ being the Weyl-Brauer operators of order $m$, regarded as elements of $\mathrm{L}\left(\mathcal{Z}^{\otimes m}\right)$. As the vectors $\left\{u_{a}: a \in \Sigma_{\mathrm{A}}\right\}$ and $\left\{v_{b}: b \in \Sigma_{\mathrm{B}}\right\}$ are unit vectors having real number entries, it follows from Proposition 6.65 that the operators $\left\{A_{a}: a \in \Sigma_{\mathrm{A}}\right\}$ and $\left\{B_{b}: b \in \Sigma_{\mathrm{B}}\right\}$ are unitary, and it is evident that they are Hermitian as well. Define

$$
\begin{equation*}
\tau=\frac{1}{2^{m}} \operatorname{vec}\left(\mathbb{1}_{\mathcal{Z}}^{\otimes m}\right) \operatorname{vec}\left(\mathbb{1}_{\mathcal{Z}}^{\otimes m}\right)^{*} \tag{6.419}
\end{equation*}
$$

For each choice of $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$ it holds that

$$
\begin{align*}
& \left\langle A_{a} \otimes B_{b}, \tau\right\rangle=\frac{1}{2^{m}} \operatorname{Tr}\left(A_{a} B_{b}^{\top}\right) \\
& \quad=\frac{1}{2^{m}} \sum_{c, d \in \Sigma_{\mathrm{A}} \sqcup \Sigma_{\mathrm{B}}}\left\langle u_{a}(c) V_{f(c)}, v_{b}(d) V_{f(d)}\right\rangle=\left\langle u_{a}, v_{b}\right\rangle, \tag{6.420}
\end{align*}
$$

again by Proposition 6.65. This is equivalent to statement 2 (taking $\Gamma^{m}$ in place of $\Gamma$ ). It has therefore been proved that statement 4 implies statement 2 , which completes the proof.

As a consequence of Tsirelson's theorem (Theorem 6.62), there exists a semidefinite program for the supremum value of the inner product $\langle K, C\rangle$, for $K$ taking the form (6.387) and for $C$ ranging over all quantum correlation operators of the form

$$
\begin{equation*}
C \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}} \times \Gamma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}} \times \Gamma_{\mathrm{A}}}\right) \tag{6.421}
\end{equation*}
$$

for $\Sigma_{\mathrm{A}}$ and $\Sigma_{\mathrm{B}}$ being arbitrary alphabets and $\Gamma_{\mathrm{A}}$ and $\Gamma_{\mathrm{B}}$ both being equal to the binary alphabet $\Gamma=\{0,1\}$.

To understand why this is so, consider an arbitrary quantum correlation operator $C$, which must be given by

$$
\begin{equation*}
C((a, c),(b, d))=\left\langle\mu_{a}(c) \otimes \nu_{b}(d), \rho\right\rangle \tag{6.422}
\end{equation*}
$$

for every $a \in \Sigma_{\mathrm{A}}, b \in \Sigma_{\mathrm{B}}$, and $c, d \in \Gamma$, for some choice of complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$, a state $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$, and two collections of measurements $\left\{\mu_{a}: a \in \Sigma_{\mathrm{A}}\right\}$ and $\left\{\nu_{b}: b \in \Sigma_{\mathrm{B}}\right\}$ whose elements take the form

$$
\begin{equation*}
\mu_{a}: \Gamma \rightarrow \operatorname{Pos}(\mathcal{X}) \quad \text { and } \quad \nu_{b}: \Gamma \rightarrow \operatorname{Pos}(\mathcal{Y}) . \tag{6.423}
\end{equation*}
$$

For an operator $K$ of the form (6.387) for some choice of $M \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}}}\right)$, one has that the value of the inner product $\langle K, C\rangle$ is given by

$$
\begin{equation*}
\sum_{(a, b) \in \Sigma_{\mathrm{A}} \times \Sigma_{\mathrm{B}}} M(a, b)\left\langle\left(\mu_{a}(0)-\mu_{a}(1)\right) \otimes\left(\nu_{b}(0)-\nu_{b}(1)\right), \rho\right\rangle . \tag{6.424}
\end{equation*}
$$

Now, an operator $H$, acting on an arbitrary complex Euclidean space, may be written as

$$
\begin{equation*}
H=\mu(0)-\mu(1) \tag{6.425}
\end{equation*}
$$

for some binary-valued measurement $\mu$ if and only if $H$ is Hermitian and satisfies $\|H\| \leq 1$. Thus, an optimization of the expression (6.424) over all choices of the measurements $\left\{\mu_{a}: a \in \Sigma_{\mathrm{A}}\right\}$ and $\left\{\nu_{b}: b \in \Sigma_{\mathrm{B}}\right\}$ is equivalent to an optimization of the expression

$$
\begin{equation*}
\sum_{(a, b) \in \Sigma_{\mathrm{A}} \times \Sigma_{\mathrm{B}}} M(a, b)\left\langle A_{a} \otimes B_{b}, \rho\right\rangle \tag{6.426}
\end{equation*}
$$

over all collections

$$
\begin{equation*}
\left\{A_{a}: a \in \Sigma_{\mathrm{A}}\right\} \subset \operatorname{Herm}(\mathcal{X}) \quad \text { and } \quad\left\{B_{b}: b \in \Sigma_{\mathrm{B}}\right\} \subset \operatorname{Herm}(\mathcal{Y}) \tag{6.427}
\end{equation*}
$$

of Hermitian operators satisfying $\left\|A_{a}\right\| \leq 1$ and $\left\|B_{b}\right\| \leq 1$, for every $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$, respectively.

By optimizing over all complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$ and density operators $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$, one finds (by Theorem 6.62) that the supremum value of $\langle K, C\rangle$ over all quantum correlation operators $C$ is equal to the supremum value of the inner product $\langle M, X\rangle$ over all choices of operators $X \in \mathrm{~L}\left(\mathbb{R}^{\Sigma_{\mathrm{B}}}, \mathbb{R}^{\Sigma_{\mathrm{A}}}\right)$ for which it holds that

$$
\left(\begin{array}{cc}
P & X  \tag{6.428}\\
X^{*} & Q
\end{array}\right) \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{A}}} \oplus \mathbb{C}^{\Sigma_{\mathrm{B}}}\right)
$$

for $P \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{A}}}\right)$ and $Q \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{B}}}\right)$ satisfying $P(a, a)=1$ and $Q(b, b)=1$ for every $a \in \Sigma_{\mathrm{A}}$ and $b \in \Sigma_{\mathrm{B}}$. Such an optimization corresponds directly to the following primal problem of a semidefinite program:

$$
\begin{array}{ll} 
& \frac{\text { Primal problem }}{\text { maximize: }} \\
\frac{1}{2}\langle M, X\rangle+\frac{1}{2}\left\langle M^{*}, X^{*}\right\rangle \\
\text { subject to: } & \left(\begin{array}{cc}
P & X \\
X^{*} & Q
\end{array}\right) \geq 0 \\
& \Delta(P)=\mathbb{1}, \Delta(Q)=\mathbb{1} \\
& P \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{A}}}\right), Q \in \operatorname{Pos}\left(\mathbb{C}^{\Sigma_{\mathrm{B}}}\right) \\
& X \in \mathrm{~L}\left(\mathbb{C}^{\Sigma_{\mathrm{B}}}, \mathbb{C}^{\Sigma_{\mathrm{A}}}\right)
\end{array}
$$

In this problem, $\Delta$ refers to the completely dephasing channel, defined with respect to either $\mathbb{C}^{\Sigma_{\mathrm{A}}}$ or $\mathbb{C}^{\Sigma_{\mathrm{B}}}$, and $\mathbb{1}$ denotes the identity operator on either of these spaces, as the context dictates without ambiguity.

The dual problem of this semidefinite program is as follows:

$$
\begin{array}{cc}
\frac{\text { Dual problem }}{} \\
\text { minimize: } & \frac{1}{2} \operatorname{Tr}(Y)+\frac{1}{2} \operatorname{Tr}(Z) \\
\text { subject to: } & \left(\begin{array}{ll}
\Delta(Y) & -M \\
-M^{*} & \Delta(Z)
\end{array}\right) \geq 0 \\
& Y \in \operatorname{Herm}\left(\mathbb{C}^{\Sigma_{\mathrm{A}}}\right) \\
& Z \in \operatorname{Herm}\left(\mathbb{C}^{\Sigma_{\mathrm{B}}}\right)
\end{array}
$$

It follows from Slater's theorem (Theorem 1.18) that strong duality holds for this semidefinite program - strict feasibility holds for both the primal and dual problems.

Example 6.66 (Tsirelson's bound) Consider the operator

$$
K=\left(\begin{array}{cccc}
1 & -1 & 1 & -1  \tag{6.429}\\
-1 & 1 & -1 & 1 \\
1 & -1 & -1 & 1 \\
-1 & 1 & 1 & -1
\end{array}\right)=M \otimes\left(\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right)
$$

for

$$
M=\left(\begin{array}{cc}
1 & 1  \tag{6.430}\\
1 & -1
\end{array}\right)
$$

which was examined in Example 6.61. One has $\|M\|=\sqrt{2}$, so that

$$
\left(\begin{array}{cc}
\sqrt{2} \mathbb{1} & -M  \tag{6.431}\\
-M^{*} & \sqrt{2} \mathbb{1}
\end{array}\right) \geq 0
$$

By taking $Y=\sqrt{2} \mathbb{1}$ and $Z=\sqrt{2} \mathbb{1}$ in the dual problem above, a feasible dual solution achieving the objective value $2 \sqrt{2}$ is obtained. Therefore,

$$
\begin{equation*}
\langle K, C\rangle \leq 2 \sqrt{2} \tag{6.432}
\end{equation*}
$$

for every quantum correlation operator $C$. The Bell inequality violation exhibited in Example 6.61 is therefore optimal for this choice of $K$.

### 6.4 Exercises

Exercise 6.1 Let $\Phi \in \mathrm{C}(\mathcal{X}, \mathcal{Y})$ be a channel, for complex Euclidean spaces $\mathcal{X}$ and $\mathcal{Y}$. Prove that the following three statements are equivalent:

1. For every complex Euclidean space $\mathcal{Z}$ and every state $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Z})$, it holds that

$$
\begin{equation*}
\left(\Phi \otimes \mathbb{1}_{\mathrm{L}(\mathcal{Z})}\right)(\rho) \in \operatorname{SepD}(\mathcal{Y}: \mathcal{Z}) \tag{6.433}
\end{equation*}
$$

2. $J(\Phi) \in \operatorname{Sep}(\mathcal{Y}: \mathcal{X})$.
3. There exists an alphabet $\Sigma$, a measurement $\mu: \Sigma \rightarrow \operatorname{Pos}(\mathcal{X})$, and a collection of states $\left\{\sigma_{a}: a \in \Sigma\right\} \subseteq \mathrm{D}(\mathcal{Y})$ such that

$$
\begin{equation*}
\Phi(X)=\sum_{a \in \Sigma}\langle\mu(a), X\rangle \sigma_{a} \tag{6.434}
\end{equation*}
$$

for all $X \in \mathrm{~L}(\mathcal{X})$.
Channels for which these statements hold are called entanglement-breaking channels.

Exercise 6.2 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, let $n=\operatorname{dim}(\mathcal{Y})$, and assume $n \leq \operatorname{dim}(\mathcal{X})$. Also let $\left\{U_{1}, \ldots, U_{m}\right\} \in \mathrm{U}(\mathcal{Y}, \mathcal{X})$ be an orthogonal collection of isometries, and let $u_{k} \in \mathcal{X} \otimes \mathcal{Y}$ be the unit vector defined as

$$
\begin{equation*}
u_{k}=\frac{1}{\sqrt{n}} \operatorname{vec}\left(U_{k}\right) \tag{6.435}
\end{equation*}
$$

for each $k \in\{1, \ldots, m\}$. Prove that if $\mu:\{1, \ldots, m\} \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ is a measurement satisfying $\mu(k) \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y})$ for every $k \in\{1, \ldots, m\}$, then

$$
\begin{equation*}
\sum_{k=1}^{m}\left\langle\mu(k), u_{k} u_{k}^{*}\right\rangle \leq \operatorname{dim}(\mathcal{X}) \tag{6.436}
\end{equation*}
$$

(Observe that a correct solution to this exercise generalizes Theorem 6.30.)

Exercise 6.3 Let $X$ and $Y$ be registers and let $\rho \in \mathrm{D}(\mathcal{X} \otimes \mathcal{Y})$ be a state of the pair $(\mathrm{X}, \mathrm{Y})$. With respect to $\rho$, one defines the entanglement of formation between $X$ and $Y$ as

$$
\begin{equation*}
\mathrm{E}_{\mathrm{F}}(\mathrm{X}: \mathrm{Y})=\inf \left\{\sum_{a \in \Sigma} p(a) \mathrm{H}\left(\operatorname{Tr}_{\mathcal{Y}}\left(u_{a} u_{a}^{*}\right)\right): \sum_{a \in \Sigma} p(a) u_{a} u_{a}^{*}=\rho\right\} \tag{6.437}
\end{equation*}
$$

where the infimum is over all choices of an alphabet $\Sigma$, a probability vector $p \in \mathcal{P}(\Sigma)$, and a collection of unit vectors $\left\{u_{a}: a \in \Sigma\right\} \subset \mathcal{X} \otimes \mathcal{Y}$ for which it holds that

$$
\begin{equation*}
\sum_{a \in \Sigma} p(a) u_{a} u_{a}^{*}=\rho \tag{6.438}
\end{equation*}
$$

(a) Prove that the infimum in (6.437) is achieved for some choice of $\Sigma, p$, and $\left\{u_{a}: a \in \Sigma\right\}$ for which $|\Sigma| \leq \operatorname{dim}(\mathcal{X} \otimes \mathcal{Y})^{2}$.
(b) Suppose that Z and W are registers and $\Phi \in \operatorname{LOCC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$ is an LOCC channel. Prove that

$$
\begin{equation*}
\mathrm{E}_{\mathrm{F}}(\mathrm{Z}: \mathrm{W})_{\sigma} \leq \mathrm{E}_{\mathrm{F}}(\mathrm{X}: \mathrm{Y})_{\rho} \tag{6.439}
\end{equation*}
$$

where $\sigma=\Phi(\rho)$ and $\mathrm{E}_{\mathrm{F}}(\mathrm{X}: \mathrm{Y})_{\rho}$ and $\mathrm{E}_{\mathrm{F}}(\mathrm{Z}: \mathrm{W})_{\sigma}$ denote the entanglement of formation of the pairs $(X, Y)$ and $(Z, W)$ with respect to the states $\rho$ and $\sigma$, respectively.
(c) Prove a more general statement than the one required of a solution to part (b), holding not only for all LOCC channels, but for all separable channels of the form $\Phi \in \operatorname{SepC}(\mathcal{X}, \mathcal{Z}: \mathcal{Y}, \mathcal{W})$.

Exercise 6.4 Let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces, and assume that both spaces have dimension at least 2 . Prove that there exist entanglementbreaking channels $\Phi_{0}, \Phi_{1} \in \mathrm{C}(\mathcal{X}, \mathcal{Y})$, as defined in Exercise 6.1, such that

$$
\begin{equation*}
\left\|\left\|\Phi_{0}-\Phi_{1}\right\|\right\|_{1}>\left\|\Phi_{0}(\rho)-\Phi_{1}(\rho)\right\|_{1} \tag{6.440}
\end{equation*}
$$

for every $\rho \in \mathrm{D}(\mathcal{X})$. Such channels have the seemingly strange property that they destroy entanglement, and yet evaluating them on an entangled state helps to discriminate between them.

Exercise 6.5 Let $\Sigma$ be an alphabet, let $\mathcal{X}$ and $\mathcal{Y}$ be complex Euclidean spaces of the form $\mathcal{X}=\mathbb{C}^{\Sigma}$ and $\mathcal{Y}=\mathbb{C}^{\Sigma}$, let $n=|\Sigma|$, and consider the projections $\Delta_{0}, \Delta_{1}, \Pi_{0}$, and $\Pi_{1}$ defined in Example 6.10. Also define

$$
\begin{equation*}
\rho_{0}=\frac{\Pi_{0}}{\binom{n+1}{2}}, \quad \rho_{1}=\frac{\Pi_{1}}{\binom{n}{2}}, \quad \sigma_{0}=\Delta_{0}, \quad \text { and } \quad \sigma_{1}=\frac{\Delta_{1}}{n^{2}-1} \tag{6.441}
\end{equation*}
$$

The states $\rho_{0}$ and $\rho_{1}$ are therefore Werner states, while $\sigma_{0}$ and $\sigma_{1}$ are isotropic states.
(a) Prove that if $\mu:\{0,1\} \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ is a measurement satisfying $\mu(0), \mu(1) \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y})$, then

$$
\begin{equation*}
\frac{1}{2}\left\langle\mu(0), \rho_{0}\right\rangle+\frac{1}{2}\left\langle\mu(1), \rho_{1}\right\rangle \leq \frac{1}{2}+\frac{1}{n+1} . \tag{6.442}
\end{equation*}
$$

Prove that there exists an LOCC measurement $\mu$ for which (6.442) holds with equality.
(b) Prove that if $\nu:\{0,1\} \rightarrow \operatorname{Pos}(\mathcal{X} \otimes \mathcal{Y})$ is a measurement satisfying $\nu(0), \nu(1) \in \operatorname{PPT}(\mathcal{X}: \mathcal{Y})$, then

$$
\begin{equation*}
\frac{1}{2}\left\langle\nu(0), \sigma_{0}\right\rangle+\frac{1}{2}\left\langle\nu(1), \sigma_{1}\right\rangle \leq 1-\frac{1}{2 n+2} . \tag{6.443}
\end{equation*}
$$

Prove that there exists an LOCC measurement $\nu$ for which (6.443) holds with equality.

Exercise 6.6 Let $N$ and $m$ be positive integers, and assume that there exist unitary and Hermitian operators $U_{0}, \ldots, U_{2 m} \in \mathrm{~L}\left(\mathbb{C}^{N}\right)$ that anticommute in pairs: $U_{j} U_{k}=-U_{k} U_{j}$ for distinct choices of $j, k \in\{0, \ldots, 2 m\}$. Prove that the collection

$$
\begin{equation*}
\left\{U_{0}^{a_{0}} \cdots U_{2 m}^{a_{2 m}}: a_{0}, \ldots, a_{2 m} \in\{0,1\}, a_{0}+\cdots+a_{2 m} \text { is even }\right\} \tag{6.444}
\end{equation*}
$$

is an orthogonal collection, and conclude that $N \geq 2^{m}$. (Observe that a correct solution to this exercise implies that the Weyl-Brauer operators have the minimum possible dimension required to possess the properties mentioned above.)

### 6.5 Bibliographic remarks

Although it was not formally defined or called entanglement therein, the phenomenon of entanglement was first recognized by Einstein, Podolsky, and Rosen (1935). Einstein, Podolsky, and Rosen's work inspired Schrödinger to investigate the phenomenon of entanglement, and to give it its name; he published a three-part paper in German (Schrödinger, 1935a,b,c), as well as two related English-language papers (Schrödinger, 1935d, 1936) discussing entanglement and other issues, as they pertained to the nature of quantum physics at that time. (An English translation of Schrödinger's three-part paper in German was published later (Trimmer, 1980).) The identification of entanglement with a lack of separability is due to Werner (1989), who used
the terms classically correlated and EPR correlated rather than separable and entangled.

The equivalence of the first two statements in Theorem 6.9 was proved by M. Horodecki, P. Horodecki, and R. Horodecki (1996), and Proposition 6.6 was proved by P. Horodecki (1997). Several elementary analytic facts about the set of separable states that have been discussed in Section 6.1.1 were also observed in the papers proving these facts. The equivalence of the third statement in Theorem 6.9 to the first two was proved a few years later by P. Horodecki (2001). In general, it is likely to be a computationally difficult task to test a bipartite density operator for separability, as suggested by the computational hardness result proved by Gurvits (2003).

The fact that any operator sufficiently close to the identity operator in a bipartite tensor product space is separable was first proved by Życzkowski, P. Horodecki, Sanpera, and Lewenstein (1998). Theorem 6.13 is due to Gurvits and Barnum (2002).

The local operations and classical communication paradigm, also called the distant labs paradigm, arose naturally in quantum information theory as various quantum information processing tasks were considered. Among the first researchers to consider this paradigm were Peres and Wootters (1991), who compared the capabilities of LOCC measurements to general measurements in a setting in which information is encoded into bipartite product states. The teleportation procedure of Bennett, Brassard, Crepéau, Jozsa, Peres, and Wootters (1993) followed shortly after.

There are natural extensions of the definition of LOCC channels that have not been discussed in this chapter. In particular, the definition of LOCC channels in the present chapter requires an LOCC channel to be a finite composition of one-way LOCC channels, corresponding to a fixed number of classical message transmissions between two individuals implementing the channel, but one may also consider channels implemented by a potentially unbounded number of message transmissions. It is known that the set of LOCC channels, as they have been defined in this chapter, is generally not closed for a fixed choice of spaces; this was proved (for bipartite channels) by Chitambar, Leung, Mančinska, Ozols, and Winter (2014). The definition of LOCC channels presented in this chapter is based on one of the definitions considered by these authors.

The class of separable channels was identified by Vedral, Plenio, Rippin, and Knight (1997), although they did not raise the possibility (first suggested by Rains (1997)) that some separable channels might not be LOCC channels. The existence of separable measurements that are not LOCC measurements (and, in fact, not even approached by a sequence of LOCC measurements
in the limit) was proved by Bennett, DiVincenzo, Fuchs, Mor, Rains, Shor, Smolin, and Wootters (1999b). Childs, Leung, Mančinska, and Ozols (2013) give a simplified proof of this fact, along with some generalizations of it.

The distillable entanglement and entanglement cost measures were defined by Bennett, Bernstein, Popescu, and Schumacher (1996a). They used the term entanglement of formation rather than entanglement cost-but that terminology has since come to refer to the measure of entanglement described in Exercise 6.3. Theorem 6.38 was proved in the same paper through the design and analysis of LOCC channels for entanglement distillation and its reverse for pure states.

Entanglement distillation for general quantum states was considered by Bennett, Brassard, Popescu, Schumacher, Smolin, and Wootters (1996c) and Bennett, DiVincenzo, Smolin, and Wootters (1996b) around the same time. It is known that the entanglement cost of every bipartite entangled state is nonzero (Yang et al., 2005).

The entanglement rank was first defined by Terhal and P. Horodecki (2000), who referred to it as the Schmidt number of a density operator (as it generalizes the number of nonzero terms in a Schmidt decomposition of the vector representation of a given pure state). They also proved that the entanglement rank of a state cannot increase under the action of an LOCC channel, based on related observations by Lo and Popescu (2001) regarding pure states, and that it is generally not multiplicative with respect to tensor products.

Theorem 6.30 was proved by Nathanson (2005), and Theorem 6.32 was proved by Walgate, Short, Hardy, and Vedral (2000).

The equivalence of statements 1,2 , and 3 in Theorem 6.33 , as well as statement 4 for LOCC channels rather than separable channels, was proved by Nielsen (1999). Nielsen's proof used the fact that every bipartite pure state transformation induced by an LOCC channel is also induced by a oneway LOCC channel, which was proved earlier by Lo and Popescu (2001). The proof of Theorem 6.38 concerning entanglement distillation and cost for pure states also appears in the same paper of Nielsen. The equivalence between statement 4 of Nielsen's theorem and the first three was proved by Gheorghiu and Griffiths (2008).

Peres (1996) proposed the computationally efficient partial transpose test for separability of bipartite density operators; he observed that separable states are necessarily PPT, and that interesting families of entangled states were revealed to be entangled through this test. By the Horodecki criterion (Theorem 6.9) proved shortly after, it follows that the partial transpose test correctly identifies all entangled state in a tensor product of two complex

Euclidean spaces, both of dimension 2 or one of dimension 2 and one of dimension 3, based on work of Størmer (1963) and Woronowicz (1976), but that entangled PPT states in higher dimensions must exist (Horodecki et al., 1996). The first explicit examples of entangled PPT states were given by P. Horodecki (1997); the unextendable product set construction of such states is due to Bennett, DiVincenzo, Mor, Shor, Smolin, and Terhal (1999c), who introduced the notion of an unextendable product set as well as the specific example given in this chapter. Proposition 6.44 and Theorem 6.45 were proved by M. Horodecki, P. Horodecki, and R. Horodecki (1998).

As was already mentioned above, the teleportation procedure described in Example 6.50 is due to Bennett, Brassard, Crepéau, Jozsa, Peres, and Wootters (1993). The dense coding procedure described in Example 6.55 is due to Bennett and Wiesner (1992). These procedures have been generalized in various ways. The general presentation of teleportation and dense coding in this chapter is based on work of Werner (2001).

The fact that entangled states may induce non-classical correlations was discovered by Bell in a highly influential 1964 paper (Bell, 1964). The Bell inequality described in Example 6.61 is due to Clauser, Horn, Shimony, and Holt (1969). Some entangled states fail to induce non-classical correlationsthis was demonstrated for the special case in which projective measurements are made on the two parts of a bipartite state by Werner (1989), and for the general case (allowing arbitrary measurements) by Barrett (2002). The entangled states constructed by Werner that have this property are among those described in Example 6.10. Theorem 6.62 is due to Tsirelson (1987).

This chapter has presented just a small part of an extensive body of work on entanglement. Readers interested in learning more about this topic are referred to the survey of R. Horodecki, P. Horodecki, M. Horodecki, and K. Horodecki (2009).


[^0]:    ${ }^{3}$ As one may expect, the notation $\operatorname{SepC}(\mathcal{X}: \mathcal{Y})$ is a shorthand for $\operatorname{SepC}(\mathcal{X}, \mathcal{X}: \mathcal{Y}, \mathcal{Y})$.

