

Upper bounds on 1D and 2D local stabilizer codes [Bravyi & Terhal 0810.1783]

Setup:  $C$  an  $[[N, k, d]]$  stabilizer code

stabilizers  $S_1, S_2, \dots, S_m$  (possibly an overcomplete set  
eg, for toric code, product of all plaquettes is  $\mathbb{1}$ )

geometric locality:

qubits can be placed on vertices of a  $D$ -dimensional square lattice,  $N = n^D$ ,  
s.t. each  $S_j$  is supported on a hypercube of side-length  $r$  (allowing  
periodic boundary conditions)

Theorem!: Under the above conditions,  $d \leq r \cdot n^{D-1}$  (if  $n \geq 2(r-1)^2$ )

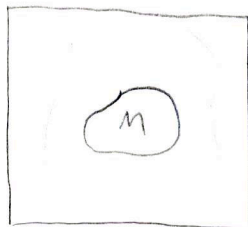
$\Rightarrow$  1D codes cannot have macroscopic distance

only  $n^{D-1}$

bound saturated by the 2D toric code, no known matching codes for  $D \geq 3$

Remark: Same bound holds for the distance protecting any one logical qubit,  
ie. also for subsystem codes with local stabilizers. (Note, though, that the  
 $n \times n$  Bacon-Shor code does not have local stabilizers, only local gauge operators.)

Lemma (Cleaning): Let  $M$  be an arbitrary subset of qubits. Then either  
there is a logical operator supported inside  $M$ , or every logical operator  
has an equivalent representative supported entirely off of  $M$ .



Proof: Let  $S(M) = \{ \text{stabilizers supported within } M \}$ ,  
an abelian group.

Let  $S_M = \{ \text{Paulis } M \text{ that can be extended to an} \}$   
element of the stabilizer, a (nonabelian) group.  $S(M) \subseteq S_M$ .

For a logical operator  $L \in N(S)$ , we can write  $L = L_M \otimes L_{\bar{M}}$ , and  
the claim is that either we can cancel  $L_M$  using elements of  $S_M$ , or  
that there is an  $L'_M \otimes I_{\bar{M}} \in N(S)$ .

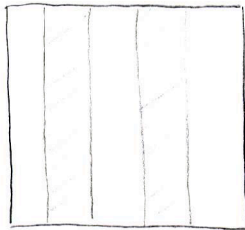
Consider the stabilizer code stabilized by  $S(M)$  on  $M$ . Since every  
element of  $S_M$  commutes with every element of  $S(M)$ ,  $S_M \setminus S(M)$  are logical  
operators for the code.

1. If  $S_M \setminus S(M)$  are all the logical operators, ie.  $N(S(M)) = S_M \setminus M$ ,  
then  $L_M \in N(S(M))$  can indeed be canceled by elements of  $S_M$ .

2. Otherwise, there is another logical operator, i.e. something that commutes with the elements of  $S(M)$  and  $S_M$ , that is supported on  $M$ , and is not in  $S_M$ . But this is also a logical operator for the original code.  $\square$

Proof of Theorem 1: Let  $L$  be a <sup>nontrivial</sup> logical operator.

Divide the lattice into vertical strips of width  $r$ .



Clean out every other strip.

$\Rightarrow$  Either  $L$  has an equivalent representative  $L'$  supported only on the odd strips, or for one of the strips there is (another) logical op. supported inside that strip. If  $d > r \cdot n^{D-1}$ , then the latter possibility isn't.

Now  $L'$  restricted to any one of the strips is also a logical operator, i.e. commutes with every stabilizer. This is because no (local) stabilizer can cross over a cleaned strip. One of these restrictions must be nontrivial, giving  $d \leq r \cdot n^{D-1}$ .  $\square$

Note: This argument required that  $n$  be a multiple of  $2r$ , so there are an even number of strips of width  $r$  (even to deal with periodic boundary conditions). But the same argument works for strips of width  $r$  or  $r-1$  and all large enough  $n$  can be divided into an even number of strips of width  $r$  or  $r-1$ .  $\checkmark$

If  $\mathcal{H} = -\sum_{j=1}^m S_j$ , then the energy cost of an error  $E$  is  
 $\mathcal{E}(E) = \#\{S_j \text{ that } E \text{ anticommutes with}\}$ .

A walk of Pauli operators is a sequence that changes one qubit at a time. Let

$$\mathcal{E}_{\max}(\gamma) = \max_{P \in \gamma} \mathcal{E}(P)$$

be the maximum energy reached by  $\gamma$ .

Theorem 2: For a code satisfying the above conditions on a  $D=2$  dimensional lattice, such that every qubit is touched by  $\mathcal{O}(1)$  generators, there must be a logical operator that can be walked to with only  $\mathcal{O}(1)$  maximum energy cost.

Proof: Consider a logical operator supported in a vertical strip of width  $r$ . It can be generated at constant cost by walking downward.  $\square$

A self-correcting memory should satisfy:

1. Its groundspace is the codespace of a Hamiltonian with macroscopic distance, i.e.  $\Omega(n)$ .
2. A macroscopic energy barrier has to be traversed by any sequence of one-qubit Paulis that generate a logical error.

Theorem 1 makes this impossible for stabilizer codes in 1D, Theorem 2 for 2D. It may still be possible for a subsystem code with the Hamiltonian the sum of gauge operators. Only Theorem 1 is known to generalize to the case of local gauge operators.